

Precise electroweak calculation of the charged current Drell-Yan process

Carlo M. Carloni Calame

INFN, Sezione di Pavia

Southampton, November 3, 2006

thanks to Stefano and the SHEP group for the invitation!

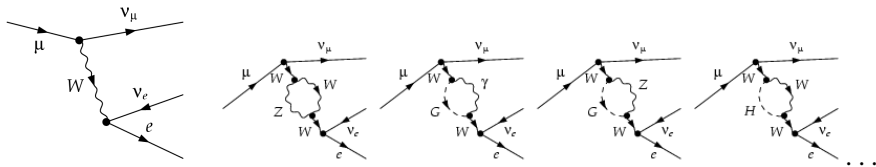
- Motivations for precise predictions of Drell-Yan processes
 - ★ precise measurement of M_W (and Γ_W)
 - ★ luminosity monitor
 - ★ PDFs constraint
 - ★ background to New Physics searches
- Overview of the literature
 - ★ QCD calculations
 - ★ EW calculations
- The event generator [HORACE](#)
 - first version (with a QED Parton Shower)
 - inclusion of exact $\mathcal{O}(\alpha)$ electro-weak corrections
 - technicalities
 - results
- Conclusions & outlook

M_W in the Standard Model

$$\mathcal{L}_{SM} = \mathcal{L}_{SM}(\alpha, M_W, M_Z; M_H; m_f; ckm)$$

$$\frac{G_\mu}{\sqrt{2}} = \frac{g^2}{8M_W^2}(1 + \Delta r) \quad \Delta r = \Delta r(m_{top}, M_W, M_Z, M_H, \dots)$$

$$M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right) = \frac{\pi\alpha}{\sqrt{2}G_\mu(1 - \Delta r)}$$

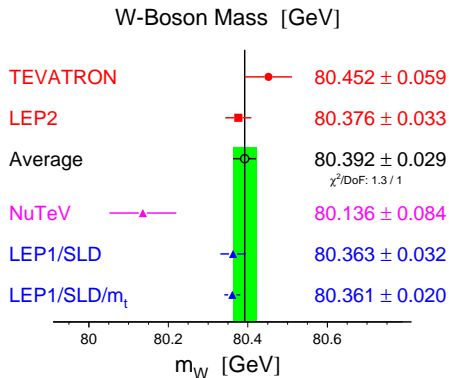


- the W mass can be predicted

$$M_W^2 = \frac{M_Z^2}{2} \left(1 + \sqrt{1 - \frac{4\pi\alpha(1 + \Delta r)}{G_\mu\sqrt{2}M_Z^2}} \right) \rightarrow M_W = 80.363 \pm 0.032 \text{ GeV}$$

Direct measurement of M_W

- at LEP2, from $e^+e^- \rightarrow WW$ (at threshold and higher energies)
- at hadron colliders, **from the M_T distribution**



Future goals for ΔM_W

- ★ Tevatron Run II $\Rightarrow 27$ MeV
- ★ LHC $\Rightarrow 15$ MeV

Future goals for $\Delta \Gamma_W$

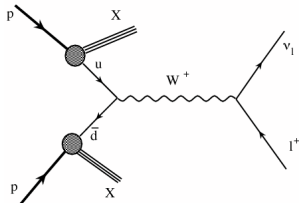
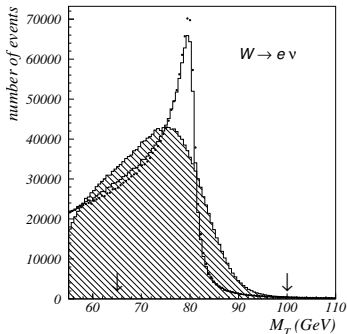
- ★ Tevatron Run II $\Rightarrow 30$ MeV
- ★ LHC $\Rightarrow \leq 30$ MeV

- A small ΔM_W (and Δm_{top}) will constraint the indirect limit on M_H

$$\Delta M_W = 27 [15] \text{ MeV and } \Delta m_{top} = 2.7 [1] \text{ GeV} \rightarrow \Delta M_H / M_H \simeq 35 [18]\%$$

M_W at Hadron Colliders

- M_W is extracted from the p_{\perp}^{ℓ} distribution, showing a (Jacobian) peak at $M_W/2$
- more reliable is $M_T^W = \sqrt{2p_{\perp}^{\ell}p_{\perp}^{\nu}(1 - \cos \phi_{\ell\nu})}$
 - ★ less sensitive to QCD corrections (e.g. p_{\perp}^W)



- The th. description of M_T spectrum has to match the aimed exp. accuracy

- the ratio $\frac{d\sigma/dM_T^W}{d\sigma/dM_T^Z}$ can be also used to extract M_W . Competitive at high luminosities

- it is possible to monitor the collider luminosity, the parton luminosities or to measure the PDFs

Frixione & Mangano '04 and refs. therein

- ★ relevant observables: total cross section, W (and Z) rapidity distribution, lepton rapidity distribution

$$\sigma^{\text{exp}} \equiv \frac{1}{\text{BR}(W \rightarrow \ell\nu)} \frac{1}{\int \mathcal{L} dt} \frac{N^{\text{obs}}}{A} = \sigma^{\text{theory}} \equiv \sum_{ab} \mathcal{P}_{ab} \otimes \hat{\sigma}_{ab}$$

- ★ an accuracy of $\mathcal{O}(1\%)$ is required/achievable
- DY processes are background to New Physics searches
 - ★ invariant mass (transverse mass) distribution tails have to be precisely simulated

Theoretical cross section at hadron colliders

$$\sigma^{\text{theory}} = \sum_{i,j} \int_0^1 dx_1 dx_2 f_{i,A}(x_1, \mu^2) f_{j,B}(x_2, \mu^2) \int d\sigma_{i,j}^H(x_1 x_2 s, \mu^2)$$

- it relies on *factorization theorems*
- $f_{k,C}$ ($k = u, \bar{u}, d, \dots, g, \dots$) are the PDFs of hadron C
- $d\sigma_{k,l}^H$ describes the hard parton-parton process, as accurately as possible, including
 - QCD Parton Shower evolution
 - QCD fixed order corrections
 - EW corrections
 - ...
- μ^2 is the factorization scale. Evolution up to μ^2 is driven by DGLAP equations

- NLO/NNLO corrections to W/Z total production rate

G. Altarelli, R.K. Ellis, M. Greco and G. Martinelli, Nucl. Phys. **B246** (1984) 12

R. Hamberg, W.L. van Neerven, T. Matsuura, Nucl. Phys. **B359** (1991) 343

R.V. Harlander and W.B. Kilgore, Phys. Rev. Lett. **88** (2002) 201801

- resummation of LL/NLL p_T^W / M_W logs (**RESBOS**)

C. Balazs and C.P. Yuan, Phys. Rev. **D56** (1997) 5558

- NLO ME merged with HERWIG Parton Shower [PS] (**MC@NLO**)

S. Frixione and B.R. Webber, JHEP **0206** (2002) 029

- Matrix elements Monte Carlos (**ALPGEN**, **SHERPA**,...) matched with PS

M.L. Mangano et al., JHEP **0307**, 001 (2003)

F. Krauss et al., JHEP **0507**, 018 (2005)

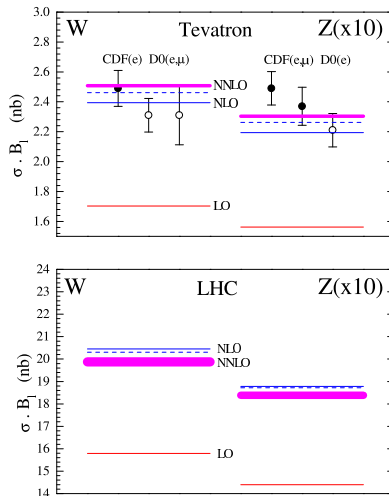
- NNLO corrections to W/Z rapidity distribution (**VRAP**)

C. Anastasiou et al., Phys. Rev. **D69** (2004) 094008

K. Melnikov and F. Petriello, hep-ph/0603182

QCD predictions for W/Z total rates

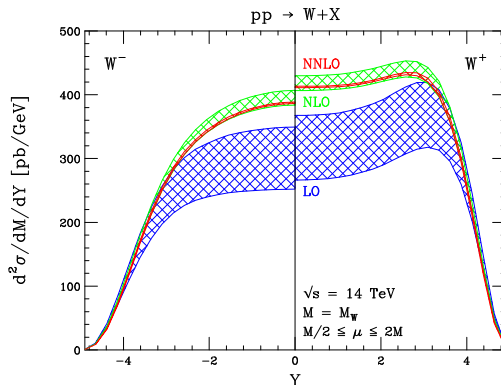
A.D. Martin *et al.*, Eur. Phys. J. **C18** (2000) 117



- Good convergence of α_s expansion. NLO-NNLO difference $\sim 2\%$ at the LHC

High-precision QCD: W/Z rapidity @ NNLO

C. Anastasiou et al., Phys. Rev. **D69** (2004) 094008



- First calculation of a differential distribution at NNLO in α_s . NNLO corrections at $\sim 2\%$ at the LHC and residual scale dependence below 1%.
- $\mathcal{O}(\alpha_s^2) \approx \mathcal{O}(\alpha_{em}) \rightarrow$ need to worry about electroweak corrections!

- Electroweak corrections to W production

- ★ Pole approximation ($\sqrt{\hat{s}} = M_W$)

- D. Wackerth and W. Hollik, PRD **55** (1997) 6788

- U. Baur et al., PRD **59** (1999) 013002

- ★ Complete $\mathcal{O}(\alpha)$ corrections

- V.A. Zykunov et al., EPJC **3** 9 (2001)

- S. Dittmaier and M. Krämer, PRD **65** (2002) 073007

- U. Baur and D. Wackerth, PRD **70** (2004) 073015

- A. Arbuzov, et al., EPJC **46**, 407 (2006)

- C.M.C.C. et al., hep-ph/0609170

DK
WGRAD2
SANC
HORACE

- Multi-photon radiation

- C.M.C.C. et al., PRD **69**, 037301 (2004), JHEP 0505:019 (2005),
hep-ph/0609170

- S. Jadach, W. Płaczek, EPJC **29** 325 (2003)

HORACE
WINHAC

- Electroweak corrections to W production

- ★ Pole approximation ($\sqrt{\hat{s}} = M_W$)

- D. Wackerth and W. Hollik, PRD **55** (1997) 6788

- U. Baur et al., PRD **59** (1999) 013002

- ★ Complete $\mathcal{O}(\alpha)$ corrections

- V.A. Zykunov et al., EPJC **3** 9 (2001)

- S. Dittmaier and M. Krämer, PRD **65** (2002) 073007

- U. Baur and D. Wackerth, PRD **70** (2004) 073015

- A. Arbuzov, et al., EPJC **46**, 407 (2006)

- C.M.C.C. et al., hep-ph/0609170

DK
WGRAD2
SANC
HORACE

- Multi-photon radiation

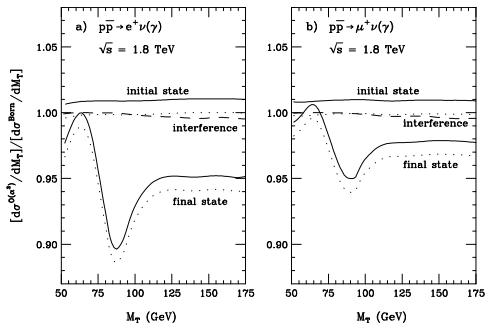
- C.M.C.C. et al., PRD **69**, 037301 (2004), JHEP 0505:019 (2005),
hep-ph/0609170

- S. Jadach, W. Płaczek, EPJC **29** 325 (2003)

HORACE
WINHAC

$\mathcal{O}(\alpha)$ EW correction effects

U. Baur, S. Keller, D. Wackerroth, Phys. Rev. **D59** (1999) 013002



- around the W peak, EW RC dominated by **final-state [FS] (photonic) corrections**
- **FS radiation** modifies the shape of the M_T distribution
- at TeVatron, $\mathcal{O}(\alpha)$ EW RC shift M_W by $\mathcal{O}(100)$ MeV

CDF coll., PRD 64 052001 (2001)

- ★ **are QED higher orders (h.o.) important?** Quoted as systematic uncertainty of $\mathcal{O}(10-20)$ MeV by TeVatron colls.

HORACE: first version

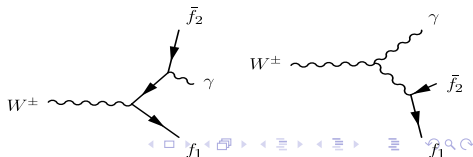
- The Monte Carlo event generator **HORACE** was originally developed to simulate **QED multi-photon radiation** in DY (W and Z) processes in **LL accuracy**, by means of a **QED Parton Shower [PS]**. Only final state radiation was accounted for

C.M.C.C. et al., PRD **69** 037301 (2004)

C.M.C.C et al., JHEP 0505:019 (2005)

- as in QCD, the QED PS solves the QED DGLAP equation, allowing for
 - ★ inclusion of QED LL corrections up to all orders (resummation)
 - ★ fully exclusive event generation (up to ∞ photons)
- by comparing the resummed PS and its $\mathcal{O}(\alpha)$ truncation, the effects purely due to h.o. can be disentangled
- photons' angular generation in W decay

$$\cos \vartheta_\gamma \propto E_\gamma^2 \left(\frac{p_\ell}{p_\ell \cdot k_\gamma} - \frac{p_W}{p_W \cdot k_\gamma} \right)^2$$



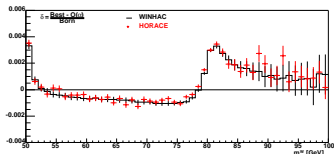
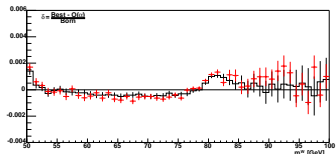
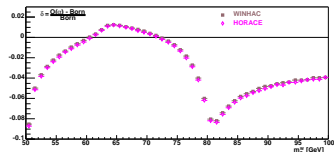
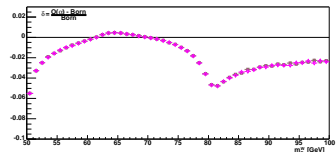
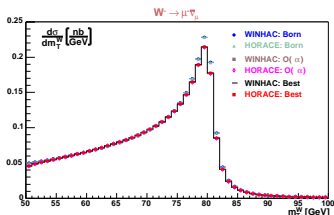
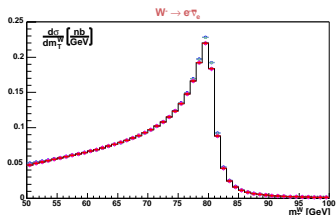
First HORACE vs WINHAC

C.M.C.C et al., Acta Phys. Pol. B35 1643 (2004), hep-ph/0402235

- during the 2003 MC4LHC CERN workshop, HORACE and WINHAC (Jadach & Płaczek) where compared
- WINHAC exploits the YFS approach to exponentiate exact $\mathcal{O}(\alpha)$ EW corrections to W decay
- e.g., cross sections for W^- production (in nb) at parton (p) and hadron (h) level, with (c) and without (nc) cuts at LHC:

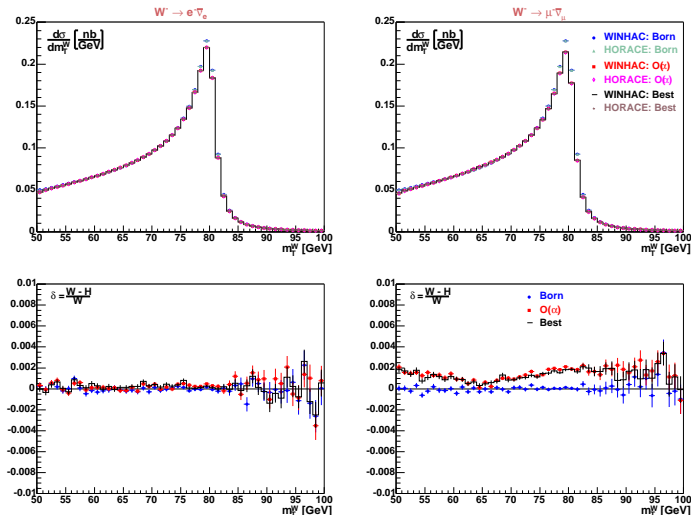
	Born		$\mathcal{O}(\alpha)$		with h.o.	
	HORACE	WINHAC	HORACE	WINHAC	HORACE	WINHAC
p e^- (nc)	8.8872	8.8871(2)	8.8872	8.8855(1)	8.8872	8.8840
p μ^- (nc)	8.8872	8.8871(1)	8.8869	8.8853(1)	8.8869	8.8844
h e^- (nc)	7.7331(4)	7.7331	7.7331(4)	7.7317(1)	7.7325(4)	7.7304
h μ^- (nc)	7.7332(4)	7.7332(1)	7.7332(4)	7.7316	7.7328(4)	7.7307
h e^- (c)	3.2363(1)	3.2363(1)	3.1871(1)	3.1878(1)	3.1869(1)	3.1876(1)
h μ^- (c)	3.2363(1)	3.2363(1)	3.1599(1)	3.1642(1)	3.1601(1)	3.1641(1)

First HORACE vs WINHAC: M_T distribution



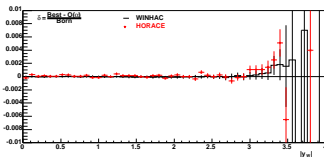
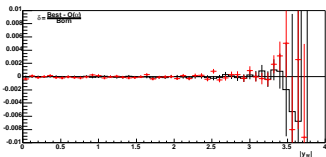
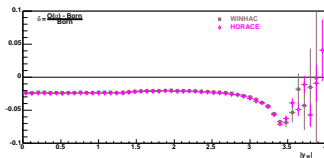
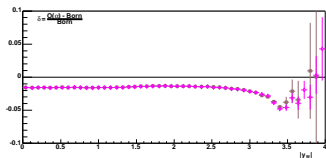
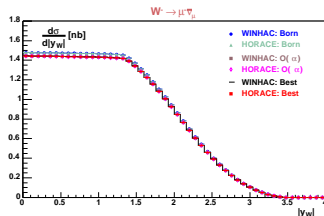
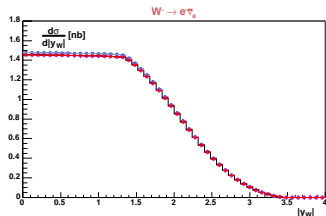
- for e^- , a calorimetric event selection [ES] criterium is used

First HORACE vs WINHAC: M_T distribution differences

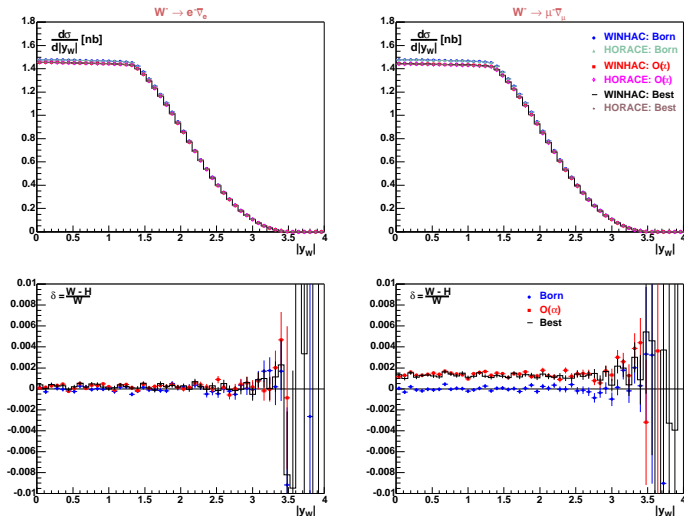


- expected (0.2%) differences at $\mathcal{O}(\alpha)$. The relative effect of exponentiation is the same

First HORACE vs WINHAC: y_W distribution



First HORACE vs WINHAC: y_W distribution differences



- expected (0.2%) differences at $\mathcal{O}(\alpha)$. The relative effect of exponentiation is the same

M_W shift induced by h.o. QED RC

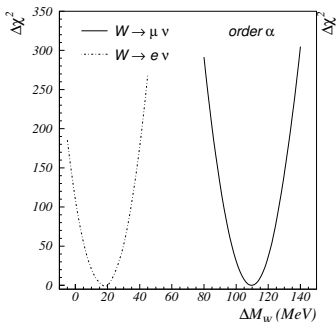
- by means of the first HORACE, we estimated the impact of multi-photon radiation on the extraction of M_W from M_T distr.
 - a “pseudo-experiment” was performed (χ^2 analysis):
- 1 generate a sample of pseudo-data at Born level for M_W^{ref}
 - 2 consider the M_T spectrum and bin it into 100 bins within the fit region 65 - 100 GeV
 - 3 consider N different M_W around M_W^{ref} and generate N radiatively corrected M_T spectra
 - 4 for each mass, calculate the χ^2 between corrected and Born spectra

$$\chi^2(M_W) = \sum_{i=bins} \left(\frac{d\sigma_{i,QED}}{\sigma_{QED}} - \frac{d\sigma_{i,Born}}{\sigma_{Born}} \right)^2 / \left[\left(\Delta \frac{d\sigma_{i,QED}}{\sigma_{QED}} \right)^2 + \left(\Delta \frac{d\sigma_{i,Born}}{\sigma_{Born}} \right)^2 \right]$$

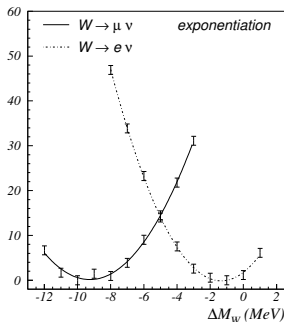
- 5 at the minimum of the χ^2 distribution, read the M_W shift

M_W shift induced by h.o. QED RC

- the exercise was performed by including (naive) detector effects, which are important



$$\Delta M_W^{\alpha,e} \sim 20 \text{ MeV}$$
$$\Delta M_W^{\alpha,\mu} \sim 110 \text{ MeV}$$

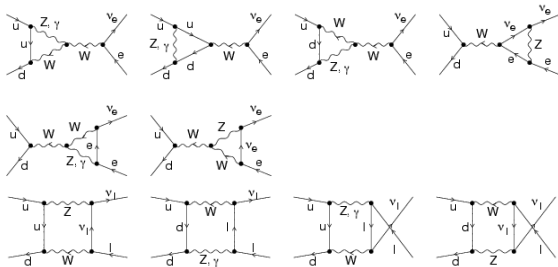


$$\Delta M_W^{\infty,e} \sim 2 \text{ MeV}$$
$$\Delta M_W^{\infty,\mu} \sim 10 \text{ MeV}$$

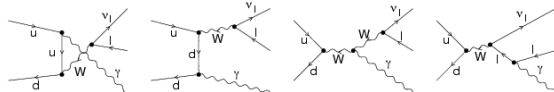
- for the electron, a recombination criterium was adopted \rightarrow smaller effect
- W -mass shift due to multiphoton radiation is about **10%** of that caused by one photon emission \rightarrow **non negligible for precise W mass!**

- <http://www.pv.infn.it/hepcomplex/horace.html>
- HORACE now includes **exact $\mathcal{O}(\alpha)$ EW corrections**, in order to go **beyond the LL QED accuracy and include weak corrections** (e.g. important at high M_T)

★ virtual RC



★ real RC



NLO EW calculation

As usual, the (partonic) NLO EW calculation is split into two parts

- $u\bar{d} \rightarrow \ell^+ \nu_\ell$

$$\mathcal{M}_{2 \rightarrow 2} = \mathcal{M}_0 + \mathcal{M}_\alpha^{virt}(\lambda)$$

- $u\bar{d} \rightarrow \ell^+ \nu_\ell \gamma$

$$\mathcal{M}_{2 \rightarrow 3} = \mathcal{M}_\alpha^{soft}(\lambda, \Delta E) + \mathcal{M}_\alpha^{hard}(\Delta E)$$

$$|\mathcal{M}_\alpha^{soft}(\lambda, \Delta E)|^2 = \delta^{soft}(\lambda, \Delta E) |\mathcal{M}_0|^2$$

- ★ $2 \rightarrow 2$ cross section

$$d\sigma_{2 \rightarrow 2} = d\sigma_{SV} \propto |\mathcal{M}_0|^2 + 2\Re[\mathcal{M}_0^* \mathcal{M}_\alpha^{virt}(\lambda)] + \delta^{soft}(\lambda, \Delta E) |\mathcal{M}_0|^2$$

- ★ $2 \rightarrow 3$ cross section

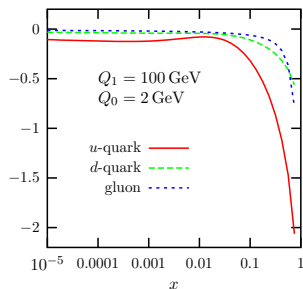
$$d\sigma_{2 \rightarrow 3} = d\sigma_H \propto |\mathcal{M}_\alpha^{hard}(\Delta E)|^2$$

- the phase space integration is performed with MC techniques
- infrared singularities are regularized with a small photon mass λ ;
collinear ones with a finite (unphysical) quark mass
- ★ IS collinear singularities must be subtracted

Subtraction of initial state collinear singularities

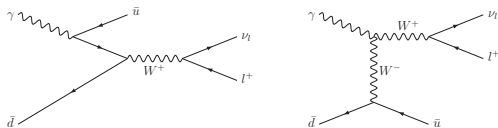
- IS quark masses regularize the collinear QED divergencies
- the QED IS singularities **have to be subtracted from the hard cross section [in analogy with NLO QCD]**, since they are already accounted in the (QED) evolution of PDFs
- the set **MRSTQED (2004)** includes the QED evolution

δf [%]



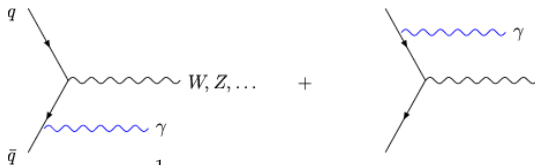
e.g. M. Roth, S. Weinzierl, PLB 590 190 (2004)

- ★ QED evolution modifies PDFs at 0.1% level for $x < 0.1$
- ★ dynamic generation of photon distr. function. **Need to include photon induced processes in DY**



Subtraction of IS collinear singularities

- QED initial-state collinear singularities are universal \rightarrow **are absorbed into PDFs**, as in QCD
- the singular partonic cross section is convoluted with a modified PDF, subtracting IS singularities



$$f(x) \rightarrow f(x, \mu_F^2) - \int_x^1 \frac{dz}{z} f\left(\frac{x}{z}, \mu_F^2\right) \frac{\alpha}{2\pi} Q_q^2 \times$$

$$\times \left\{ \ln\left(\frac{\mu_F^2}{m_q^2}\right) [P_{ff}(z)]_+ - [P_{ff}(z) (2 \ln(1-z) + 1)]_+ + C(z) \right\}$$

$$C(z) = \left\{ \begin{array}{l} 0 \\ [P_{ff}(z) (\ln(\frac{1-z}{z}) - \frac{3}{4}) + \frac{9+5z}{4}]_+ \end{array} \right. \overline{\text{MS}} \text{ DIS}$$

- during the Les Houches 2005 workshop, a tuned comparison among EW $\mathcal{O}(\alpha)$ calculations was done

- ★ DK (Dittmaier & Krämer), HORACE, SANC, WGRAD2

- Setup for comparison (cuts & lepton ID)

- ★ LHC, W^+ production

$$p_{\perp}^{\ell} > 25 \text{ GeV} \quad p_{\perp}^{\text{miss}} \equiv p_{\perp}^{\nu} > 25 \text{ GeV} \quad |\eta_{\ell}| < 2.5$$

$\ell = \text{bare } \mu^+, \text{bare } e^+, \text{recombined } e^+$

- ★ recombination criteria (kill FS collinear $\alpha \log \frac{s}{m_e^2}$, due to KLN theorem):

$$\text{if } |\eta_{\gamma}| > 2.5 \text{ and } R_{e+\gamma} = \sqrt{(\eta_{e^+} - \eta_{\gamma})^2 + \phi_{e+\gamma}^2} \leq 0.1 \text{ photon and electron momenta are summed (i.e. } p_e = p_e + p_{\gamma}\text{)}$$

- similar comparison for the TeV4LHC workshop still going on. Naive-detector effects are also investigated here.

Les Houches comparisons, varying p_{\perp}^{ℓ} cut

C. Buttar et al., hep-ph/0604120

pp $\rightarrow \nu_l l^+ (+\gamma)$ @ $\sqrt{s} = 14$ TeV (with MRSTQED04)

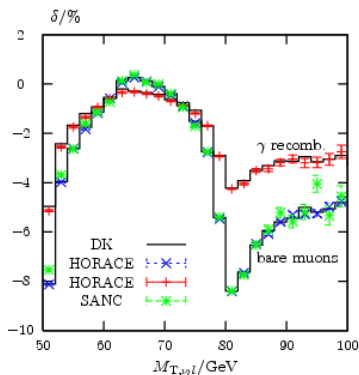
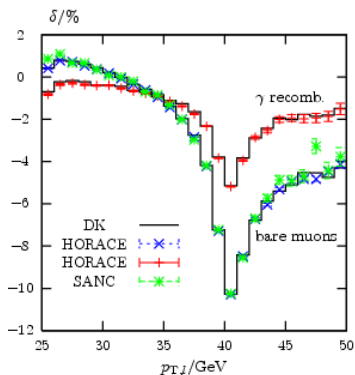
$p_{T,l}/\text{GeV}$	25- ∞	50- ∞	100- ∞	200- ∞	500- ∞	1000- ∞
σ_0/pb						
DK	2112.2(1)	13.152(2)	0.9452(1)	0.11511(2)	0.0054816(3)	0.00026212(1)
HORACE	2112.21(4)	13.151(6)	0.9451(1)	0.11511(1)	0.0054812(4)	0.00026211(2)
SANC	2112.22(2)	13.1507(2)	0.94506(1)	0.115106(1)	0.00548132(6)	0.000262108(3)
WGRAD	2112.3(1)	13.149(1)	0.94510(5)	0.115097(5)	0.0054818(2)	0.00026209(2)
$\delta_{e+\nu_e}/\%$						
DK	-5.19(1)	-8.92(3)	-11.47(2)	-16.01(2)	-26.35(1)	-37.92(1)
HORACE	-5.23(1)	-8.98(1)	-11.49(1)	-16.03(1)	-26.36(1)	-37.92(2)
WGRAD	-5.10(1)	-8.55(5)	-11.32(1)	-15.91(2)	-26.1(1)	-38.2(2)
$\delta_{\mu+\nu_{\mu}}/\%$						
DK	-2.75(1)	-4.78(3)	-8.19(2)	-12.71(2)	-22.64(1)	-33.54(2)
HORACE	-2.79(1)	-4.84(1)	-8.21(1)	-12.73(1)	-22.65(1)	-33.57(1)
SANC	-2.80(1)	-4.82(2)	-8.17(2)	-12.67(2)	-22.63(2)	-33.50(2)
WGRAD	-2.69(1)	-4.53(1)	-8.12(1)	-12.68(1)	-22.62(2)	-33.6(2)
$\delta_{\text{recomb}}/\%$						
DK	-1.73(1)	-2.45(3)	-5.91(2)	-9.99(2)	-18.95(1)	-28.60(1)
HORACE	-1.77(1)	-2.51(1)	-5.94(1)	-10.02(1)	-18.96(1)	-28.65(1)
SANC	-1.89(1)	-2.56(1)	-5.97(1)	-10.02(1)	-18.96(1)	-28.61(1)
WGRAD	-1.71(1)	-2.32(1)	-5.94(1)	-10.11(2)	-19.08(3)	-28.73(6)
$\delta_{\gamma q}/\%$						
DK	+0.071(1)	+5.24(1)	+13.10(1)	+16.44(2)	+14.30(1)	+11.89(1)

Les Houches comparisons: distributions

- EW corrections varying M_T cut

M_T/GeV	50- ∞	100- ∞	200- ∞	500- ∞	1000- ∞
$\delta_{\text{rec}}/\%$	-1.73(1)	-3.43(2)	-6.52(2)	-12.55(1)	-19.51(1)
$\delta_{\gamma q}/\%$	+0.0567(3)	+0.1347(1)	+0.2546(1)	+0.3333(1)	+0.3267(1)

- perfect agreement between independent calculations on p_{\perp}^{ℓ}, M_T distributions (δ due to EW $\mathcal{O}(\alpha)$)



The new HORACE: version 2.0

C.M.C.C., Montagna, Nicosini, Vicini, hep-ph/0609170, submitted to JHEP

- we would like to combine (merge, match) the old QED PS based formulation with the exact EW $\mathcal{O}(\alpha)$ calculation, in order to
 - ★ preserve PS advantages (multi-photon effects, exclusive event generation)
 - ★ go beyond its approximation (LL accuracy, missing contributions already at $\mathcal{O}(\alpha)$)
- the matching has to avoid the double counting of $\mathcal{O}(\alpha)$ LL, already accounted for by the PS, and to “produce” a formula well suited for Monte Carlo generation
- a solution to the problem has been found. . .

- the issue has a long story also in QCD (e.g. MC@NLO, POWHEG)

PS and exact $\mathcal{O}(\alpha)$ matrix elements (at parton level)

Consider the LL [$LL \equiv PS$] resummed, LL $\mathcal{O}(\alpha)$ and exact $\mathcal{O}(\alpha)$ cross sections

- $d\sigma_{LL}^{\infty} = \Pi(Q^2, \varepsilon) \sum_{n=0}^{\infty} \frac{1}{n!} |\mathcal{M}_{n,LL}|^2 d\Phi_n$
- $d\sigma_{LL}^{\alpha} = [1 + C_{\alpha,LL}] |\mathcal{M}_0|^2 d\Phi_0 + |\mathcal{M}_{1,LL}|^2 d\Phi_1 \equiv d\sigma_{SV}(\varepsilon) + d\sigma_H(\varepsilon)$
- $d\sigma_{exact}^{\alpha} = [1 + C_{\alpha}] |\mathcal{M}_0|^2 d\Phi_0 + |\mathcal{M}_1|^2 d\Phi_1$
- $F_{SV} = 1 + (C_{\alpha} - C_{\alpha,LL}) \quad F_H = 1 + \frac{|\mathcal{M}_1|^2 - |\mathcal{M}_{1,LL}|^2}{|\mathcal{M}_{1,LL}|^2}$
- $d\sigma_{exact}^{\alpha} \stackrel{\text{at } \mathcal{O}(\alpha)}{=} F_{SV} (1 + C_{\alpha,LL}) |\mathcal{M}_0|^2 d\Phi_0 + F_H |\mathcal{M}_{1,LL}|^2 d\Phi_1$

$$d\sigma_{matched}^{\infty} = F_{SV} \Pi(Q^2, \varepsilon) \sum_{n=0}^{\infty} \frac{1}{n!} \left(\prod_{i=0}^n F_{H,i} \right) |\mathcal{M}_{n,LL}|^2 d\Phi_n$$

Contents of the *matched* formula

- F_{SV} and $F_{H,i}$ are infrared safe and account for missing EW $\mathcal{O}(\alpha)$, avoiding double counting of QED LL
- $[\sigma_{matched}^\infty]_{\mathcal{O}(\alpha)} = \sigma_{exact}^\alpha$
- $\sigma_{matched}^\infty$ is “made of” exact $\mathcal{O}(\alpha)$ one-loop building blocks
- resummation of higher-order LL contributions preserved
- the cross section is still **fully differential** in the momenta of the final state particles (including the photons)
- Problem:
the $\mathcal{O}(\alpha)$ calculation presents IS collinear singularities ($\propto \alpha \log \frac{M_W^2}{m_q^2}$), which are exponentiated in the matched formula.
A subtraction procedure up to all orders needs to be devised!

Contents of the *matched* formula

- F_{SV} and $F_{H,i}$ are infrared safe and account for missing EW $\mathcal{O}(\alpha)$, avoiding double counting of QED LL
- $[\sigma_{matched}^\infty]_{\mathcal{O}(\alpha)} = \sigma_{exact}^\alpha$
- $\sigma_{matched}^\infty$ is “made of” exact $\mathcal{O}(\alpha)$ one-loop building blocks
- resummation of higher-order LL contributions preserved
- the cross section is still **fully differential** in the momenta of the final state particles (including the photons)
- **Problem:**
the $\mathcal{O}(\alpha)$ calculation presents IS collinear singularities ($\propto \alpha \log \frac{M_W^2}{m_q^2}$), which are exponentiated in the matched formula.
A subtraction procedure up to all orders needs to be devised!

IS subtraction to all orders

- the idea is to start with a **subtracted $\mathcal{O}(\alpha)$ cross section**

$$d\sigma_{exact}^\alpha = d\sigma_{exact}^\alpha - d\sigma_{sub}^\alpha + d\sigma_{sub}^\alpha \equiv d\tilde{\sigma}_{exact}^\alpha + d\sigma_{sub}^\alpha$$

where $d\sigma_{sub}^\alpha$ contains the LL initial state collinear singularities, both for the S+V part and the real part

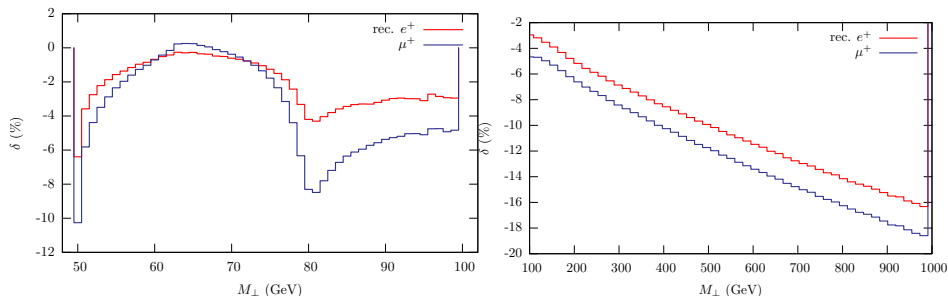
(see hep-ph/0609170 for explicit formulae)

- then the resummation and matching are performed by **using the subtracted $\mathcal{O}(\alpha)$ building blocks of $d\tilde{\sigma}_{exact}^\alpha$** , which is free of IS collinear singularities
- proof of the independence from quark masses**
e.g., W^+ cross section (nb) at LHC (within typical cuts)

	$\mathcal{O}(\alpha)$	matched
m_q	4410.98 ± 0.20	4412.14 ± 0.26
$m_q/10$	4410.92 ± 0.26	4411.89 ± 0.33
$m_q/100$	4410.99 ± 0.29	4411.92 ± 0.50

$\mathcal{O}(\alpha)$ EW results with HORACE

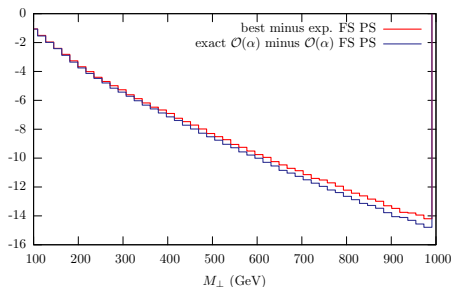
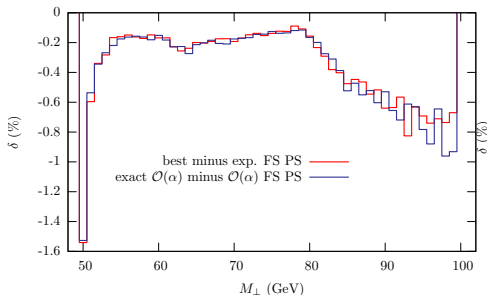
- LHC, $pp \rightarrow W^+ \rightarrow \ell^+ \nu_\ell$, $p_{\perp,\ell}$ and $p_{\perp,\nu} > 25$ GeV, $|\eta_\ell| < 2.5$
- $\mathcal{O}(\alpha)$ EW corrections to the M_T distribution



- $\mathcal{O}(\alpha)$ corrections at 5% - 10% level around the peak and increasingly large in the M_T tail due to the presence of the EW Sudakov $(\log s)^2$, $\alpha_W \log^2 \frac{s}{M_Z^2}$

Weak $\mathcal{O}(\alpha)$ and QED non-log corrections on M_T

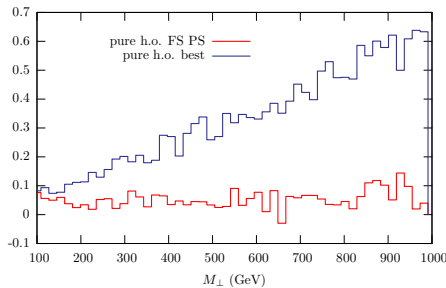
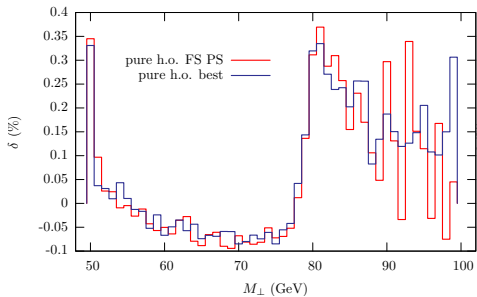
- differences between matched result and final-state [FS] QED PS



- blue = $(d\sigma_{exact}^\alpha - d\sigma_{PS}^\alpha)/d\sigma_0$ red = $(d\sigma_{matched}^\infty - d\sigma_{PS}^\infty)/d\sigma_0$
- Sum of weak $\mathcal{O}(\alpha)$, QED FS non-logs and QED IS remnant flat around the peak, increasingly large in the tail
- the FS QED PS calculation is improved consistently by missing $\mathcal{O}(\alpha)$ with the matching procedure

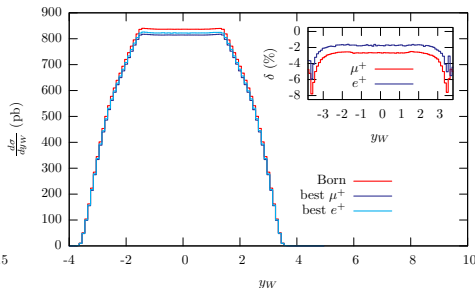
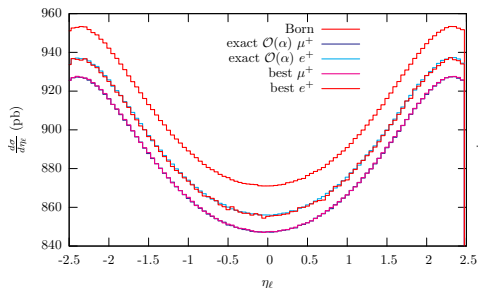
Effects of multi-photon radiation on M_T

- higher-order EW corrections to the M_T distribution



- blue = $(d\sigma_{matched}^{\infty} - d\sigma_{exact}^{\alpha})/d\sigma_0$ red = $(d\sigma_{PS}^{\infty} - d\sigma_{PS}^{\alpha})/d\sigma_0$
- QED h.o. around the peak distort the shape. In the tail, induced effects by EW Sudakov $\otimes \mathcal{O}(\alpha)$ QED LL
- the $\mathcal{O}(\alpha)$ calculation is improved consistently by h.o. with the matching procedure

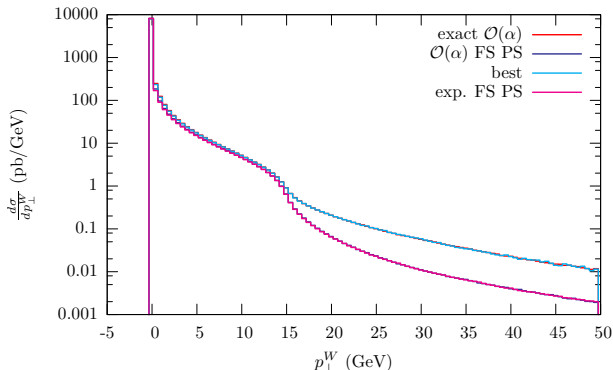
Effects on lepton and W rapidities



- $\mathcal{O}(\alpha)$ effects at the level of 2%, resummation negligible
- important for acceptances determination, luminosity and PDFs constraints

Effects on p_{\perp}^W

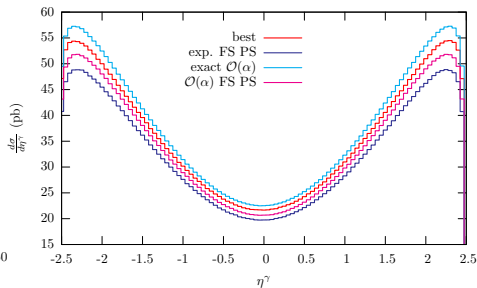
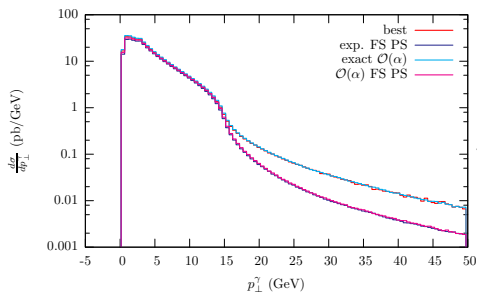
- no QCD effects
- p_{\perp}^W is defined as the p_{\perp} of the $\ell^+\nu$ pair, different from zero because of photon radiation



- at high p_{\perp}^W using the exact photon emission ME is crucial, both for $\mathcal{O}(\alpha)$ and resummed distribution

Photonic observables (radiative events)

- besides leptonic cuts, we require $|\eta_\gamma| < 2.5$ and $E_\gamma > 3$ GeV for the hardest photon
- this signature can be used e.g. to study the $WW\gamma$ trilinear vertex



- as expected, the exact real emission ME gives large corrections w.r.t. the LL approximation
- here radiative events (one more α) are selected

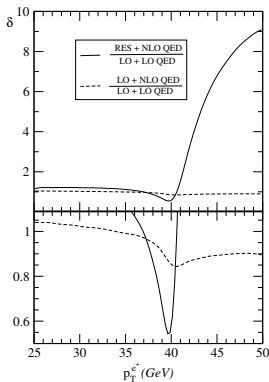
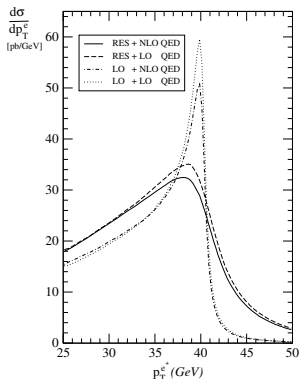
Combining EW & QCD corrections

A **unified tool** simulating “at best” EW and QCD corrections would be highly desirable for data analysis

- e.g., **RESBOS** combined with factorizable FS EW corrections of **WGRAD2** → **RESBOS-A**

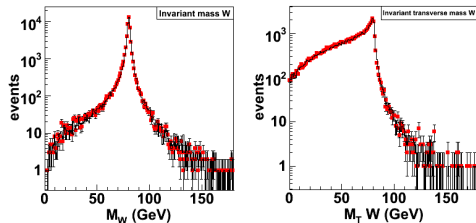
Q. Cao, C.P. Yuan, PRL 93 042001 (2004)

Q. Cao, C.P. Yuan, hep-ph/0401171

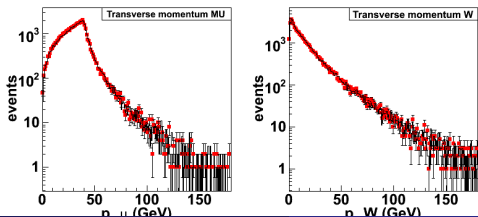


HORACE & QCD showering MCs

- HORACE is “Les Houches Accord” compliant. Its events can be passed through QCD Parton Shower & hadronization MCs (HERWIG or PYTHIA) to include at least LL QCD effects.
- e.g. HORACE (with QED PS)+PYTHIA vs. PYTHIA+PHOTOS

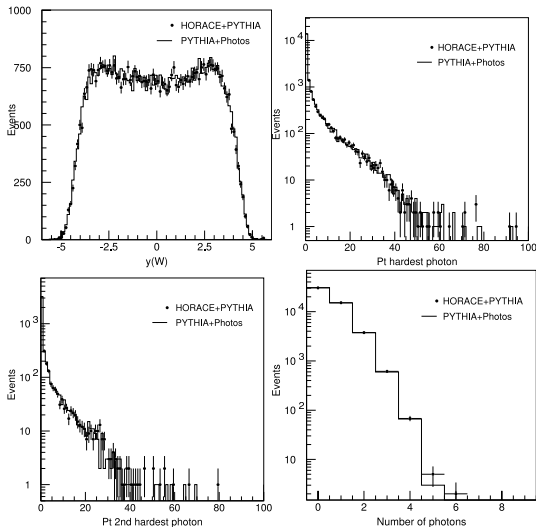


courtesy of G. Polesello and
M. Bellomo (ATLAS)



• HORACE+PYTHIA
— PYTHIA+PHOTOS

HORACE & QCD showering MCs (II)



★ interface to more refined QCD tools (MC@NLO, ALPGEN) is in progress...

Conclusions & outlook

- DY processes are a fertile ground for (challenging) precision physics at hadron colliders
 - ★ precise M_W measurement ($\Delta M_W = 15 \text{ MeV}$ at LHC)
 - ★ PDF constraints
 - ★ collider luminosity (with accuracy $\mathcal{O}(5\%)$)
 - ★ it is a background to New Physics searches
- Theoretical calculations are essential ingredients for the success of the physics program. Higher order QCD and EW corrections must be taken into account
- The Monte Carlo “EW event generator” HORACE has been developed, including
 - ★ exact $\mathcal{O}(\alpha)$ EW corrections matched with a
 - ★ QED Parton Shower to simulate multi-photon radiation
- Combining QCD and EW generators is needed for data analysis
- Work in progress to
 - ★ interface “at best” HORACE with more refined QCD tools
 - ★ extend the EW matching to $\gamma/Z \rightarrow \ell^+ \ell^-$ channel

thank you!