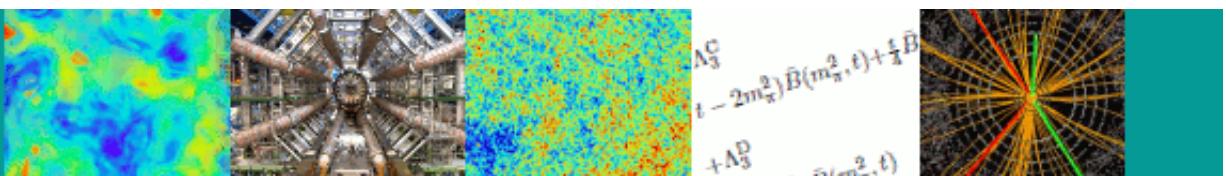


# TeV-Scale Sterile Neutrinos

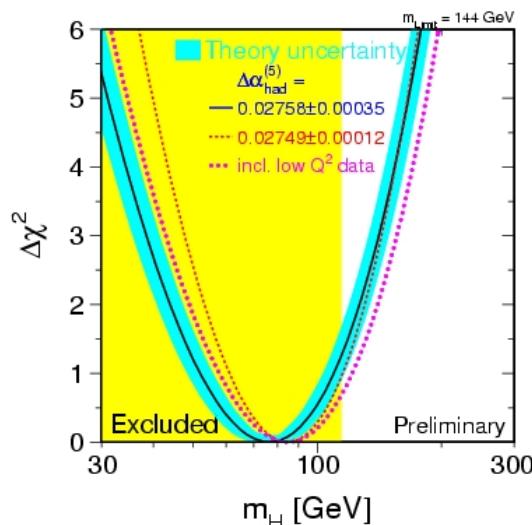
Manfred Lindner



MAX-PLANCK-INSTITUT  
FÜR KERNPHYSIK  
HEIDELBERG

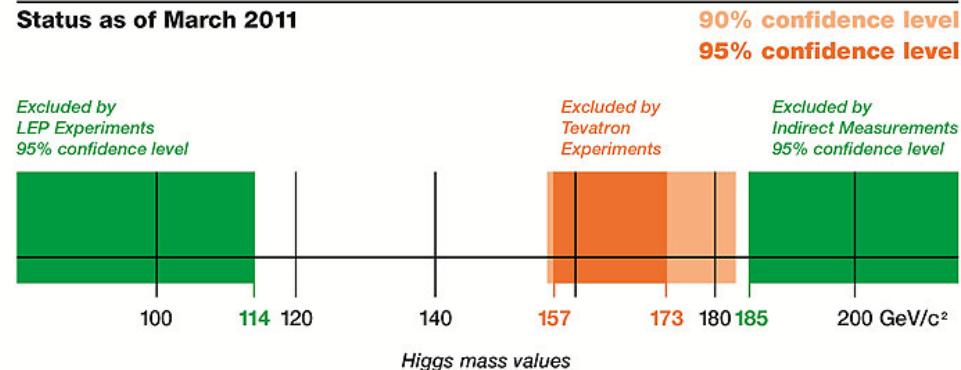


# SM works perfectly & Higgs seems to be there

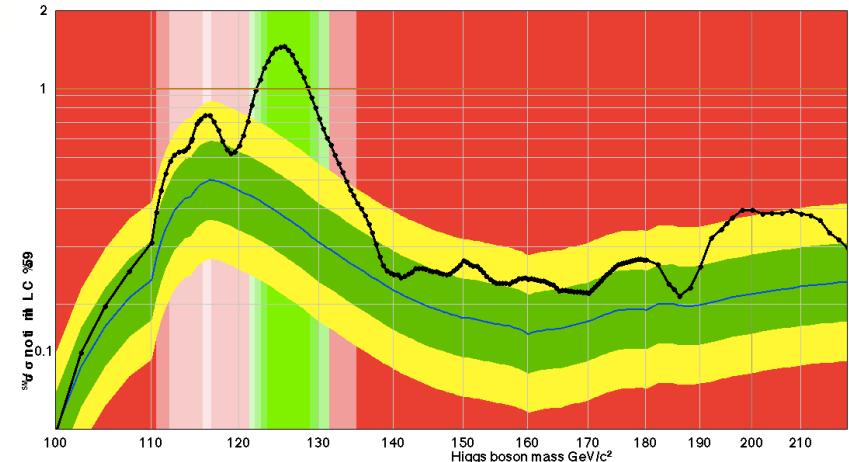


## Search for the Higgs Particle

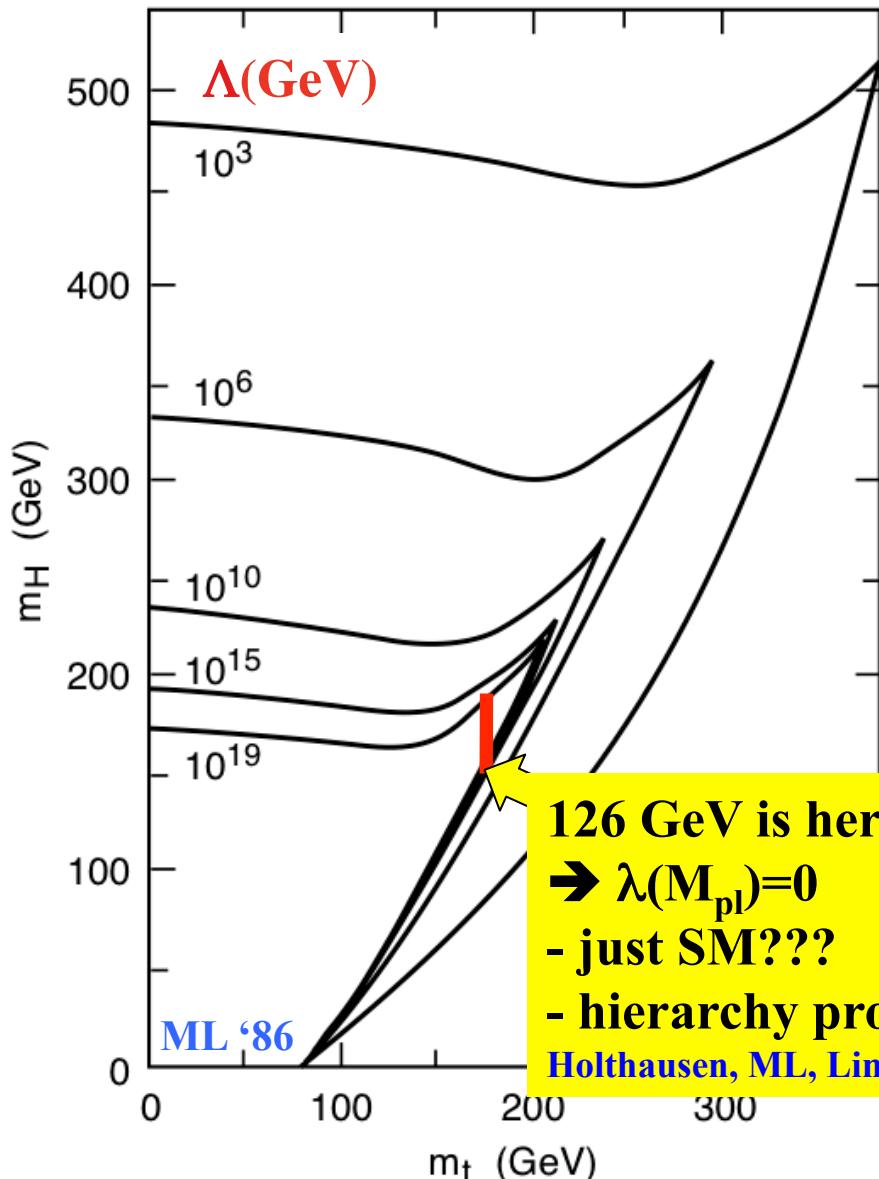
Status as of March 2011



- mass range was shrinking...
- ... and is now rather precisely known
- no signs for anything else
- just the SM?**
  
- It's a well defined QFT
- just like QED, QCD, ...
- 18+1 parameters → experimental input
- tests by over-constraining (S,T,U, running α's, ...)
- per se no hierarchy problem! ←→ embedding QFT self-consistency →

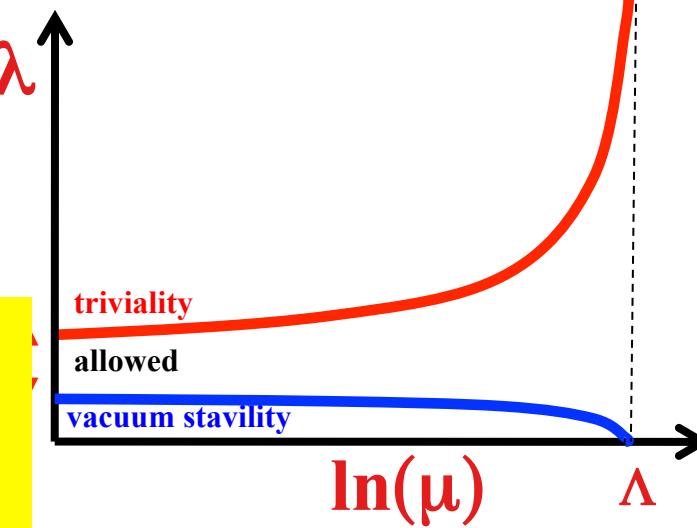


# Triviality and Vacuum Stability



126 GeV <  $m_H$  < 174 GeV

SM does not exist w/o embedding  
- U(1) coupling , Higgs self-coupling



→ RGE arguments seem to work  
→ we need some embedding

# The SM works perfectly but it must be extended...

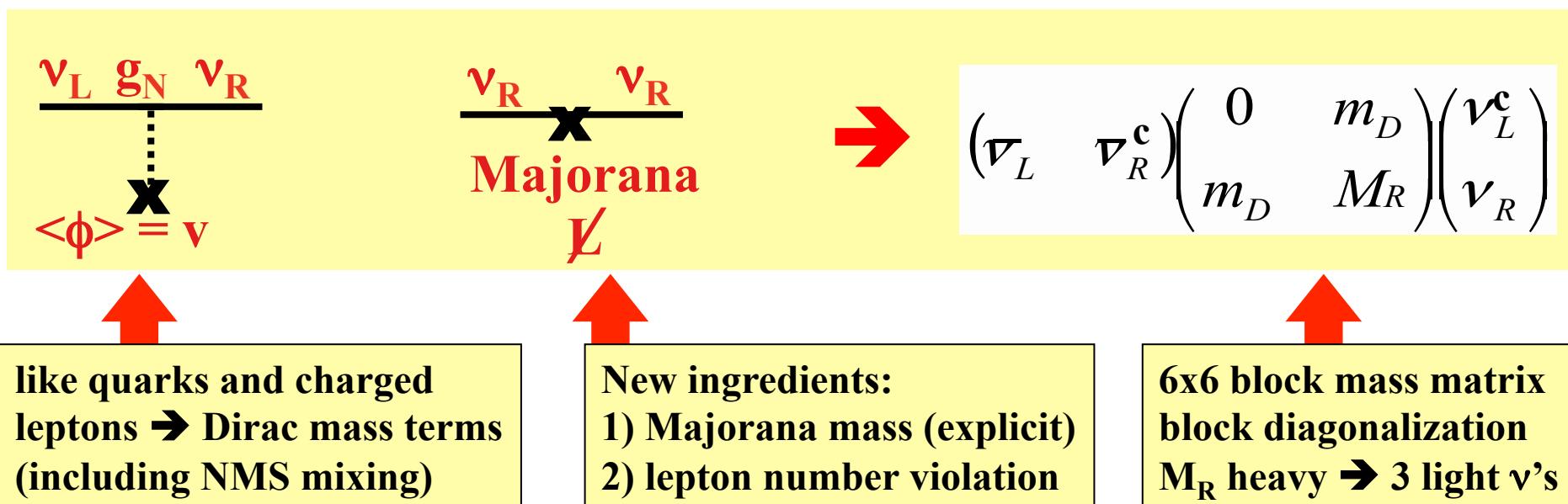
- **Many theoretical reasons for BSM physics:**
  - Flavour problem ...  $\leftrightarrow$  maybe insights via neutrinos
  - Hierarchy problem
    - only a problem once the SM is embedded...  $\rightarrow$  GUTs, ...
    - separation of two scalar scales: SM Higgs and some other scalar at  $\Lambda$
    - Planck scale physics  $\rightarrow$  new concepts, spin 2, no new scalars?
      - $\rightarrow$  no 2<sup>nd</sup> scalar for direct embeddings at the Planck scale?
      - $\rightarrow$  SM boundary conditions that point to  $\Lambda_{\text{Planck}}$   $\rightarrow$  vacuum stability line!
- **Experimental facts:**
  - SM cannot explain Baryon Asymmetry of the Universe
    - BUT: neutrino masses require SM extension  $\rightarrow$  SM+ $\rightarrow$
    - leptogenesis = one of the best BAU explanations
  - Dark Matter
    - some extra particle is required which is DM
    - particles connected to the hierarchy problem, strong CP, ...  $\rightarrow$  neutrinos
  - Neutrino Masses and maybe evidences for sterile neutrinos

# New Physics: Neutrino Mass Terms

$\text{SM} \sim m\phi \bar{L}R = (2,1)$   
 $\rightarrow$  new fields

$\rightarrow$  Simplest possibility:  
add 3 right handed neutrino fields

Field	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$
$L_Q = \begin{pmatrix} l_u \\ l_d \end{pmatrix}$	3	2	1/3
$r_u$	3	1	4/3
$r_d$	3	1	-2/3
$L_L = \begin{pmatrix} l_\nu \\ l_e \end{pmatrix}$	1	2	-1
$r_\nu$ ???	1	1	0
$r_e$	1	1	-2



NEW ingredients, 9 parameters  $\rightarrow$  SM+ and sea-saw

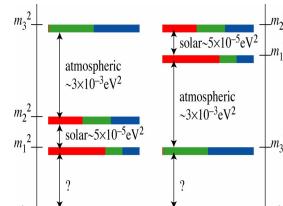
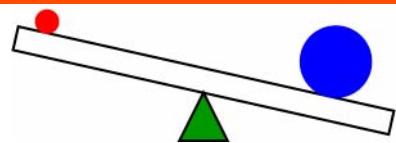
# Suggestive Seesaw Features

QFT: natural value of mass operators  $\leftrightarrow$  scale of symmetry

$m_D \sim$  electro-weak scale

$M_R \sim$  L violation scale  $\leftarrow ? \rightarrow$  embedding (GUTs, ...)

See-saw (type I)



?  $\rightarrow$  EW scale

$$m_\nu = m_D M_R^{-1} m_D^T$$

$$m_h = M_R$$

Numerical hints:

For  $m_3 \sim (\Delta m_{\text{atm}}^2)^{1/2}$ ,  $m_D \sim$  leptons  $\rightarrow M_R \sim 10^{11} - 10^{16} \text{ GeV}$   
 $\rightarrow \nu$ 's are Majorana particles,  $m_\nu$  probes  $\sim$  GUT scale physics!  
 $\rightarrow$  smallness of  $m_\nu \leftrightarrow$  high scale of L, symmetries of  $m_D, M_R$

# The Neutrino Spectrum

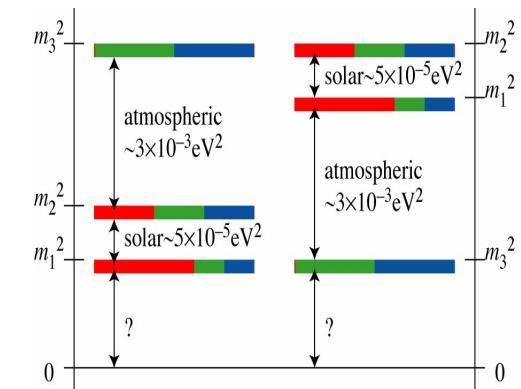
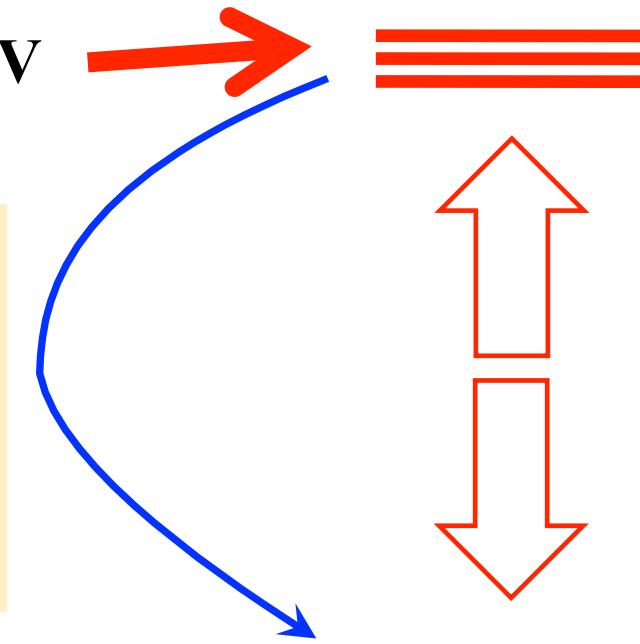
The standard picture:

3 heavy sterile neutrinos typ.  $\geq 10^{13}$  GeV

→ leptogenesis, role in GUTs, ...

Some mechanism which makes  
1, 2, ... heavy states light?  
→ light sterile neutrino(s)  
→ tiny heavy-light mixing expected  
 $\theta^2 < \mathcal{O}(m_\nu/m_s)$

3 light active neutrinos  
→ this could easily be wrong  
- more than 3  $N_R$  states, ...  
-  $M_R$  may have special eigenvalues, ...  
→ light sterile neutrinos ?!



# Evidences for Light Sterile Neutrinos

## Particle Physics:

Reactor anomaly, LSND, MiniBooNE, MINOS, Gallex...

- New and better data / experiments are needed to clarify the situation
- maybe something exciting around the corner?
- would hint at eV scale and sizable percentage type mixings

CMB: Extra eV-ish neutrinos possible J. Hamann et al. , ...

BBN: Extra ν's possible:  $N_\nu \simeq 3.7 \pm 1$

E. Aver, K. Olive, E. Skillman (2010), Y. Izotov, T. Thuan(2010)

Astrophysics: keV-ish sterile neutrinos could explain pulsar kicks

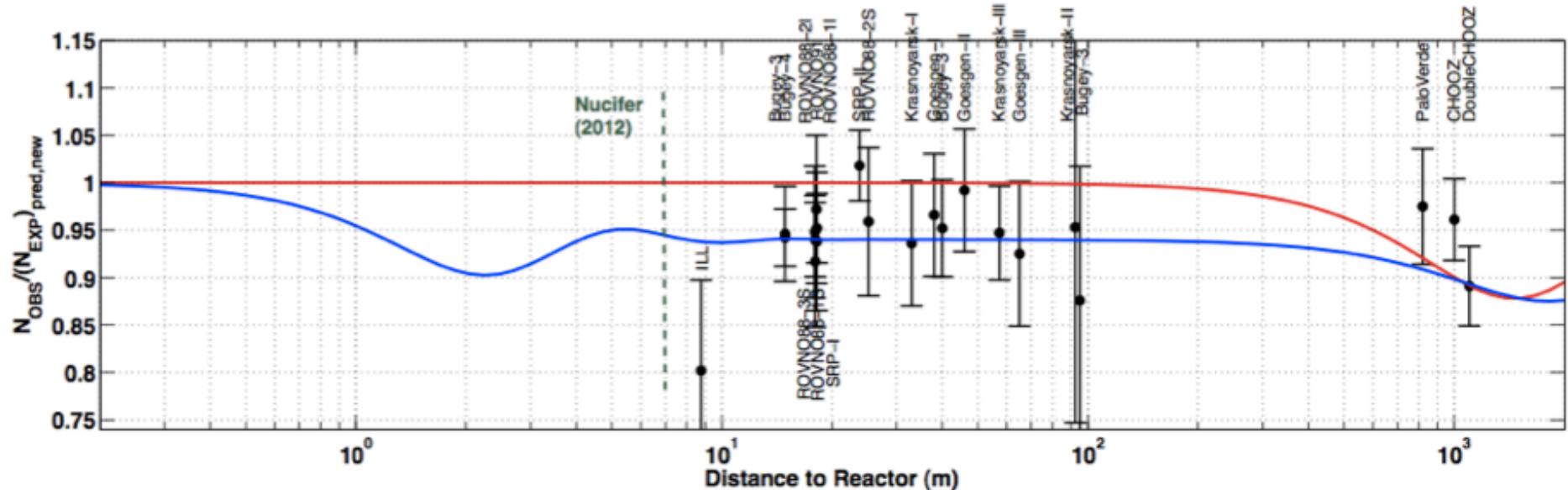
Kusenko, Segre, Mocioiu, Pascoli, Fuller et al., Biermann & Kusenko, Stasielak et al., Loewenstein et al., Dodelson, Widrow, Dolgov, ...

My believe: Most likely not all of them are correct

- but any one has far reaching consequences!

# The Reactor Anomaly

# New reactor fluxes and global reactor data: Mention et al.



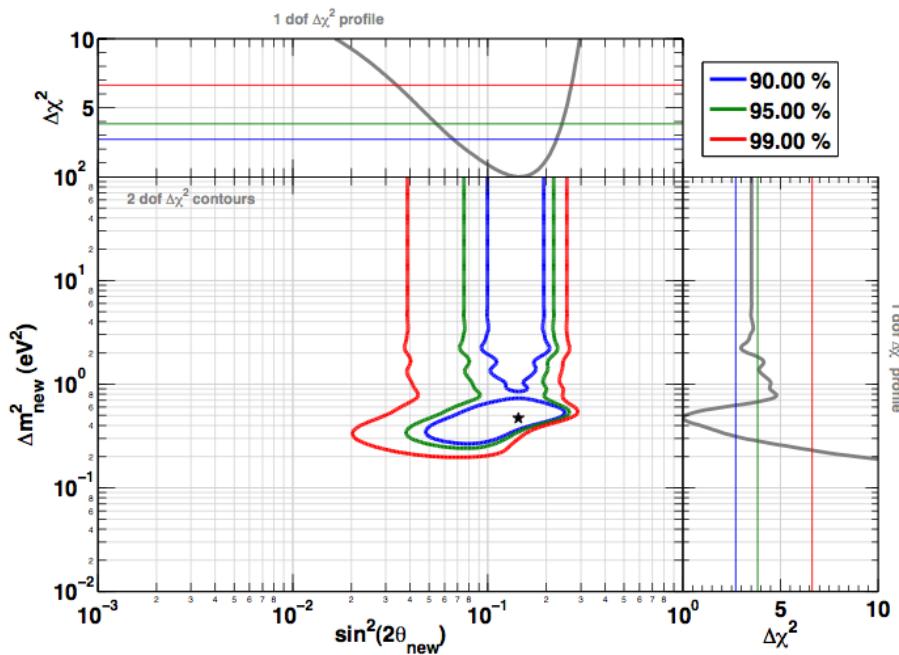
$\theta_{13}$  can reduce flux at  $L \geq 1$  km, but not at shorter baselines  
 Sterile neutrino with  $\Delta m^2 \sim \text{eV}$  can nicely account for reduction  
 $\rightarrow$  3+1 fits ; all evidences for eV scale  $\rightarrow$  3+2, ...

See e.g. T. Schwetz at NEUTRINO 2012

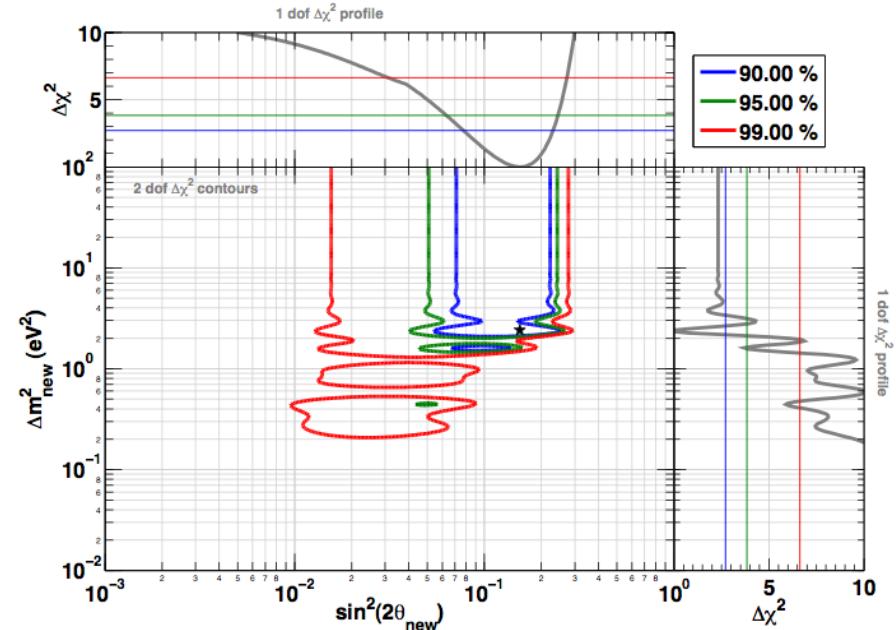
→ will be tested by new experiments (e.g. NUCIFER @ few meters)

# Reactor Anomaly

Rate only analysis



Rate + Shape (Bugey 3) analysis



- Best fit:  $\Delta m^2 = 2.4 \text{ eV}^2$  ;  $\sin^2(2\theta_{\text{new}}) = 0.14$
- 2.9 sigma significance

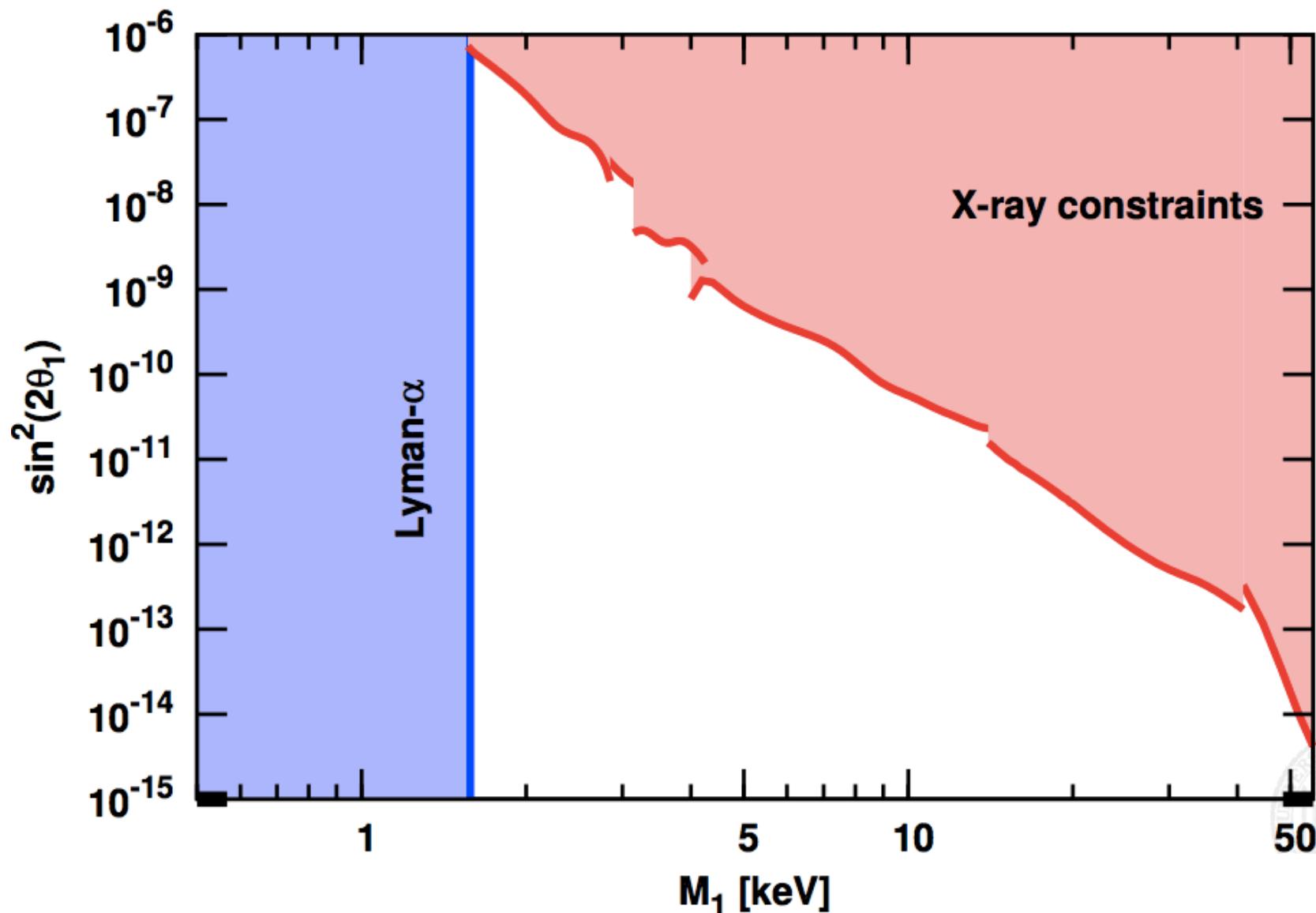
# Could Neutrinos be Dark Matter?

- Active neutrinos would be perfect Hot Dark Matter → ruled out:
  - destroys small scale structures in cosmological evolution
  - measured neutrino masses too small → maybe HDM component
- keV sterile neutrinos: Warm Dark Matter → works very well:
  - relativistic at decoupling
  - non-relativistic at radiation to matter dominance transition
  - OK for  $M_X \simeq$  few keV with very tiny mixing
  - reduced small scale structure → smoother profile, less dwarf satellites
  - scenario where one sterile neutrino is keV-ish, the others heavy
  - tiny active – sterile mixings  $O(m_\nu/M_R)$

**Note: Right-handed neutrinos exist probably anyway – just make one light!**

Asaka, Blanchet, Shaposhnikov, Asaka, Shaposhnikov; Kusenko, Segre, Mocioiu, Pascoli, Fuller et al., Biermann & Kusenko, Stasielak et al., Loewenstein et al., Dodelson, Widrow, Dolgov, ... Bezrukov, Hettmannsperger, ML

# Allowed Range for keV Sterile $\nu$ Dark Matter



# Could TeV-Scale Sterile Neutrinos Exist?

E. Akhmedov, A. Kartavtsev, ML, L. Michaels and J. Smirnov  
arXiv 1302.1872

Let's assume that TeV-scale sterile neutrinos exist.

- what are the limits?
- do they improve things or make them worse?

→ improvements of over-all fits!

# Simplest Extension of the SM

- Add 3 right-handed neutrinos (Majorana fields)  $N_i = N_i^c$

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \bar{N}_i (i\cancel{\partial} - M_i) N_i - h_{\alpha i} \bar{\ell}_\alpha \tilde{\phi} N_i - h_{i\alpha}^\dagger \bar{N}_i \tilde{\phi}^\dagger \ell_\alpha$$

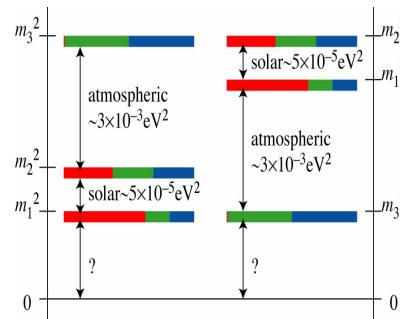
- This can solve experimentally required facts :
  - small neutrino masses via type I see-saw
  - baryon asymmetry of the Universe via leptogenesis
  - eV-ish sterile neutrinos (reactor anomaly)
  - keV sterile neutrinos (dark matter)
  - ...
- Explaining the sterile spectrum ...???
- Study effects of sterile neutrinos with any mass → TeV

# N=3 sterile Neutrinos

sterile neutrinos  $\nu_{4,5,6}$   
spectrum ???



active neutrinos  $\nu_{4,5,6}$



Global fits of active neutrino data  
usually assume 3x3 unitarity!  
→ unconstraint fits s up to 4%  
changes Antusch et al.

6x6 mixing matrix:

$$U = \begin{pmatrix} \mathcal{U} & \mathcal{R} \\ \mathcal{W} & \mathcal{V} \end{pmatrix}$$

(3 × 3) PMNS matrix  $\mathcal{U}$   
is not exactly unitary  
→ deviations

$$\epsilon_\alpha \equiv \sum_{i \geq 4} |U_{\alpha i}|^2$$

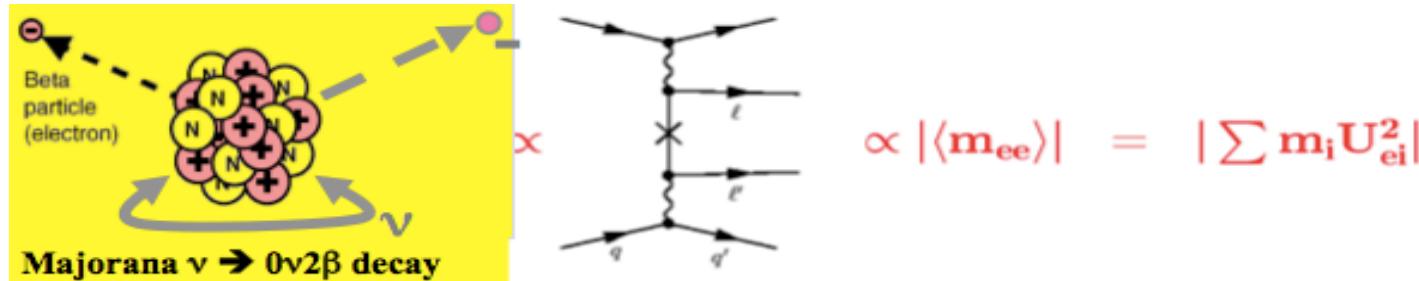
$$\epsilon_e - \epsilon_\mu = 0.0022 \pm 0.0025$$

$$\epsilon_\mu - \epsilon_\tau = 0.0017 \pm 0.0038$$

$$\epsilon_e - \epsilon_\tau = 0.0039 \pm 0.0040$$

# Consequences of sterile Neutrinos

- Shifts in active neutrino parameters in global analyses
- L-violating admixture in light  $\nu$ 's  $\rightarrow 0\nu\beta\beta$  Beta Decay



- Effective mass including all states:

$$|\langle m_{ee} \rangle| \approx \left| \sum_{i=1}^3 U_{ei}^2 m_i - \sum_{i=4}^{3+n} F(A, M_i) U_{ei}^2 m_i \right|$$

$$F(A, m_i) \approx (m_a/m_i)^2 f(A)$$

$f(A)$  depends on the decaying isotope

today  $|\langle m_{ee} \rangle| < 0.4$  eV  $\rightarrow$  will improve in next years!

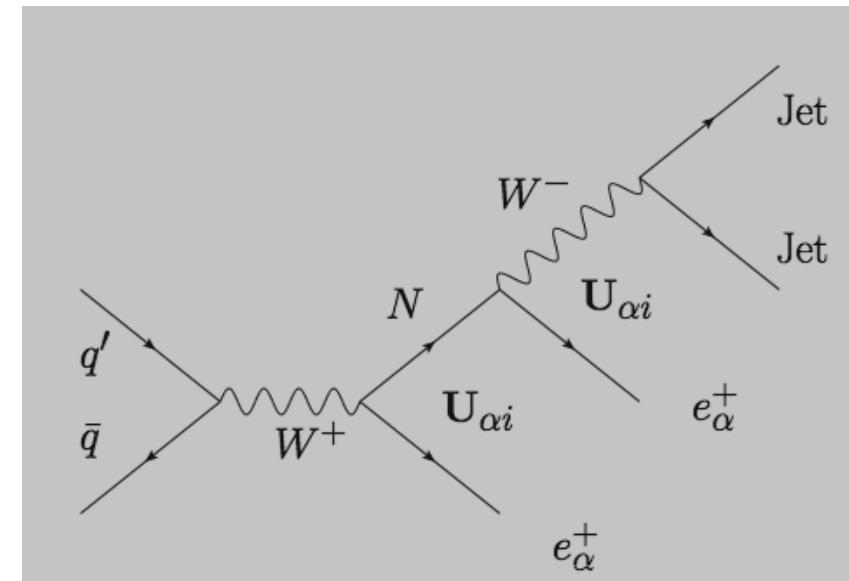
- Additional interactions of W and Z:

$$\begin{aligned}\mathcal{L}_{\text{int}} = & -\frac{e}{2c_w s_w} Z_\mu \sum_{i,j=1}^{3+n} \sum_{\alpha=e,\mu,\tau} \bar{\nu}_i \mathbf{U}_{i\alpha}^\dagger \gamma^\mu P_L \mathbf{U}_{\alpha j} \nu_j \\ & - \frac{e}{\sqrt{2} s_w} W_\mu \sum_{i=1}^{3+n} \sum_{\alpha=e,\mu,\tau} \bar{\nu}_i \mathbf{U}_{i\alpha}^\dagger \gamma^\mu P_L e_\alpha + \text{h.c.}\end{aligned}$$

→ L-violating processes

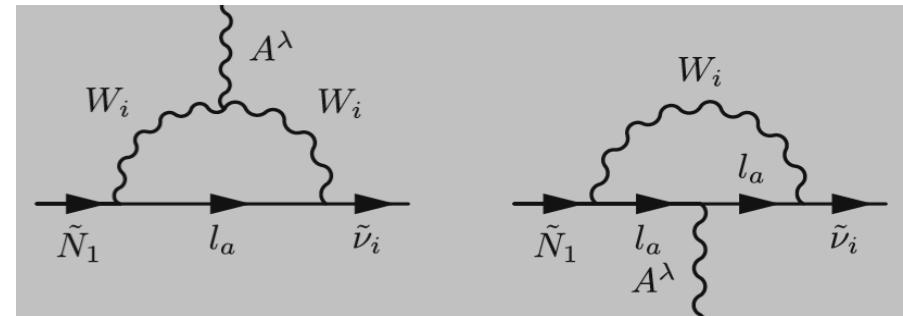
e.g. production of heavy neutrino at LHC as intermediate state

→ after upgrade to 14 TeV with 100 fb<sup>-1</sup> expected limit



$|\sum_i \mathbf{U}_{\alpha i}^2 m_i^{-1}| \simeq 6.5 \cdot 10^{-3} \text{ TeV}^{-1}$  sensitivity to 800 GeV sterile neutrinos!

- L-violating decays;  
e.g. :  $\mu \rightarrow e\gamma$



$$\text{BR}(\mu \rightarrow e\gamma) = \frac{\Gamma(\mu \rightarrow e\gamma)}{\Gamma(\mu \rightarrow e\nu\bar{\nu})} = \frac{3\alpha}{32\pi} |\delta_\nu|^2$$

:  $\delta_\nu = 2 \sum_i \mathbf{U}_{ei}^* \mathbf{U}_{\mu i} g(m_i^2/M_W^2)$

$$g(x) = \int_0^1 \frac{(1-\alpha)d\alpha}{(1-\alpha)+\alpha x} [2(1-\alpha)(2-\alpha) + \alpha(1+\alpha)x]$$

$\rightarrow$  
$$\delta_\nu = 2 \sum_{i=4}^{3+n} \mathbf{U}_{ei}^* \mathbf{U}_{\mu i} [g(m_i^2/M_W^2) - 5/3]$$

**MEG bound:**  $\text{BR}(\mu^+ \rightarrow e^+ \gamma) \leq 2.4 \cdot 10^{-12}$

- Non-unitarity in neutrino oscillations  
→ zero distance effec in addtion to oscillations

$$P_{\alpha\beta}(L = 0) = \frac{\delta_{\alpha\beta} (1 - 2\epsilon_\alpha) + \epsilon_\alpha \epsilon_\beta}{(1 - \epsilon_\alpha)(1 - \epsilon_\beta)}$$

- not easy to detect ↔ normalization
- note that this cannot explain the reactor anomaly!

- **Electro-weak precision observables**
  - tree level mixing and loop effects (S,T,U)
  - may explain some ‘deviations’
  - fix SM definition → predict other observables

EWPO	Theory (Standard Model)	Experiment
$\Gamma_{\text{lept}}$ (MeV)	$84.005 \pm 0.015$	$83.984 \pm 0.086$
$\Gamma_{\text{inv}}/\Gamma_{\text{lept}}$	$5.9721 \pm 0.0002$	$5.942 \pm 0.016$
$\sin^2 \theta_W$	$0.23150 \pm 0.0001$	$0.2324 \pm 0.0012$
$g_L^2$	$0.3040 \pm 0.0002$	$0.3026 \pm 0.0012$
$g_R^2$	$0.0300 \pm 0.0002$	$0.0303 \pm 0.0010$
$M_W$ (GeV)	$80.359 \pm 0.011$	$80.385 \pm 0.015$

- Z-width very precisely known  $\simeq 2\sigma$  too low
- NuTeV anomaly

- Sterile neutrinos lead to modifications  
→ may explain some ‘deviations’

$$\Gamma_{\text{inv}} / [\Gamma_{\text{inv}}]_{\text{SM}} = \frac{1}{3} \sum_{\alpha} (1 - \epsilon_{\alpha})^2$$

**CC and NC scattering ↔ mixing**

$$\sigma_{\alpha}^{\text{CC}} = \sigma_{\alpha, \text{SM}}^{\text{CC}} (1 - \epsilon_{\alpha}),$$

$$\sigma_{\alpha}^{\text{NC}} = \sigma_{\alpha, \text{SM}}^{\text{NC}} (1 - \epsilon_{\alpha})^2$$

**definition of  $G_{\mu}$ :**

$$G_{\mu}^2 = G_F^2 (1 - \epsilon_{\mu})(1 - \epsilon_e)$$

## Corrections due to heavy sterile ν's:

- Both tree level ( $\epsilon$ 's) and loop (S,T,U) effects
- U is tiny
- identical or almost identical parameter combinations

→ cancellations

→ some observables don't change

→ some do  $\leftrightarrow$  data

$$\frac{\Gamma_{\text{lept}}}{[\Gamma_{\text{lept}}]_{\text{SM}}} = 1 + 0.6 (\epsilon_e + \epsilon_\mu + 0.0145 T) - 0.0021 S,$$

$$\frac{\Gamma_{\text{inv}}/\Gamma_{\text{lept}}}{[\Gamma_{\text{inv}}/\Gamma_{\text{lept}}]_{\text{SM}}} = 1 - 0.67 (\epsilon_e + \epsilon_\mu + \epsilon_\tau) + 0.0021 S - 0.0015 T,$$

$$\frac{\sin^2 \theta_w^{\text{lept}}}{[\sin^2 \theta_w^{\text{lept}}]_{\text{SM}}} = 1 - 0.72 (\epsilon_e + \epsilon_\mu + 0.0145 T) + 0.0016 S,$$

$$\frac{g_L^2}{[g_L^2]_{\text{SM}}} = 1 + 0.41 \epsilon_e - 0.59 \epsilon_\mu - 0.0090 S + 0.0022 T,$$

$$\frac{g_R^2}{[g_R^2]_{\text{SM}}} = 1 - 1.4 \epsilon_e - 2.4 \epsilon_\mu + 0.031 S - 0.0067 T,$$

$$\frac{M_W}{[M_W]_{\text{SM}}} = 1 + 0.11 \epsilon_e + 0.11 \epsilon_\mu - 0.0036 S + 0.0056 T + 0.0042 U$$

- Details...

$$S_{\text{tot}} = S_N + S_{\text{SM}} = -\frac{1}{2\pi M_Z^2}$$

$$\begin{aligned} & \times [\sum_{i,j=1}^{3+n} \sum_{\alpha\beta} \mathbf{U}_{i\alpha}^\dagger \mathbf{U}_{\alpha j} \mathbf{U}_{j\beta}^\dagger \mathbf{U}_{\beta i} \Delta Q(M_Z^2, m_i^2, m_j^2) \\ & + \sum_{i,j=1}^{3+n} \sum_{\alpha\beta} \mathbf{U}_{i\alpha}^\dagger \mathbf{U}_{\alpha j} \mathbf{U}_{i\beta}^\dagger \mathbf{U}_{\beta j} m_i m_j \Delta B_0(M_Z^2, m_i^2, m_j^2) \\ & + \sum_\alpha m_\alpha^2 B_0(0, m_\alpha^2, m_\alpha^2) + \sum_\alpha Q(M_Z^2, m_\alpha^2, m_\alpha^2) \\ & - 2 \sum_\alpha m_\alpha^2 B_0(M_Z^2, m_\alpha^2, m_\alpha^2)] , \end{aligned} \quad (20)$$

$$\begin{aligned} Q(q^2, m_1^2, m_2^2) & \equiv (D-2)B_{22}(q^2, m_1^2, m_2^2) \\ & + q^2 [B_1(q^2, m_1^2, m_2^2) + B_{21}(q^2, m_1^2, m_2^2)] \end{aligned}$$

$B_0$ ,  $B_1$ ,  $B_{21}$  and  $B_{22}$   
are the usual loop functions

$$\begin{aligned} T_{\text{tot}} = T_N + T_{\text{SM}} & = -\frac{1}{8\pi s_w^2 M_W^2} \\ & \times [\sum_{i,j=1}^{3+n} \sum_{\alpha\beta} \mathbf{U}_{i\alpha}^\dagger \mathbf{U}_{\alpha j} \mathbf{U}_{j\beta}^\dagger \mathbf{U}_{\beta i} Q(0, m_i^2, m_j^2) \\ & + \sum_{i,j=1}^{3+n} \sum_{\alpha\beta} \mathbf{U}_{i\alpha}^\dagger \mathbf{U}_{\alpha j} \mathbf{U}_{i\beta}^\dagger \mathbf{U}_{\beta j} m_i m_j B_0(0, m_i^2, m_j^2) \\ & - 2 \sum_{i=1}^{3+n} \sum_\alpha \mathbf{U}_{i\alpha}^\dagger \mathbf{U}_{\alpha i} Q(0, m_i^2, m_\alpha^2) \\ & + \sum_\alpha m_\alpha^2 B_0(0, m_\alpha^2, m_\alpha^2)] , \end{aligned} \quad (1)$$

$$\begin{aligned} U_{\text{tot}} = U_N + U_{\text{SM}} & = \frac{1}{2\pi M_Z^2} \\ & \times [\sum_{i,j=1}^{3+n} \sum_{\alpha\beta} \mathbf{U}_{i\alpha}^\dagger \mathbf{U}_{\alpha j} \mathbf{U}_{j\beta}^\dagger \mathbf{U}_{\beta i} \Delta Q(M_Z^2, m_i^2, m_j^2) \\ & + \sum_{i,j=1}^{3+n} \sum_{\alpha\beta} \mathbf{U}_{i\alpha}^\dagger \mathbf{U}_{\alpha j} \mathbf{U}_{i\beta}^\dagger \mathbf{U}_{\beta j} m_i m_j \Delta B_0(M_Z^2, m_i^2, m_j^2) \\ & + \sum_\alpha m_\alpha^2 B_0(0, m_\alpha^2, m_\alpha^2) - \sum_\alpha Q(M_Z^2, m_\alpha^2, m_\alpha^2) \\ & - 2 \sum_\alpha m_\alpha^2 B_0(M_Z^2, m_\alpha^2, m_\alpha^2) \\ & - 2(M_Z/M_W)^2 \sum_\alpha \mathbf{U}_{i\alpha}^\dagger \mathbf{U}_{\alpha i} \Delta Q(M_W^2, m_i^2, m_\alpha^2)] . \end{aligned} \quad (21)$$

# Results of global Fits to all Data

- Assume Lagrangian including see-saw type I  
 $\leftrightarrow$  parameter relations
- All data: LFV, LHC, EWPO and active neutrinos and consider 3 typical mass spectra

$$\sin^2 \theta_{12} = 0.30 \pm 0.013,$$

$$\sin^2 \theta_{23} = 0.41^{+0.037}_{-0.025},$$

$$\sin^2 \theta_{13} = 0.023 \pm 0.0023,$$

$$\delta_{CP} = 300^{+66}_{-138},$$

	NH	IH	QD
$m_1$ (eV)	$\sim 0$	$4.85 \cdot 10^{-2}$	$\sim 0.1$
$m_2$ (eV)	$8.660 \cdot 10^{-3}$	$4.93 \cdot 10^{-2}$	$\sim 0.1$
$m_3$ (eV)	$4.97 \cdot 10^{-2}$	$\sim 0$	$\sim 0.1$

- Quality of fit:

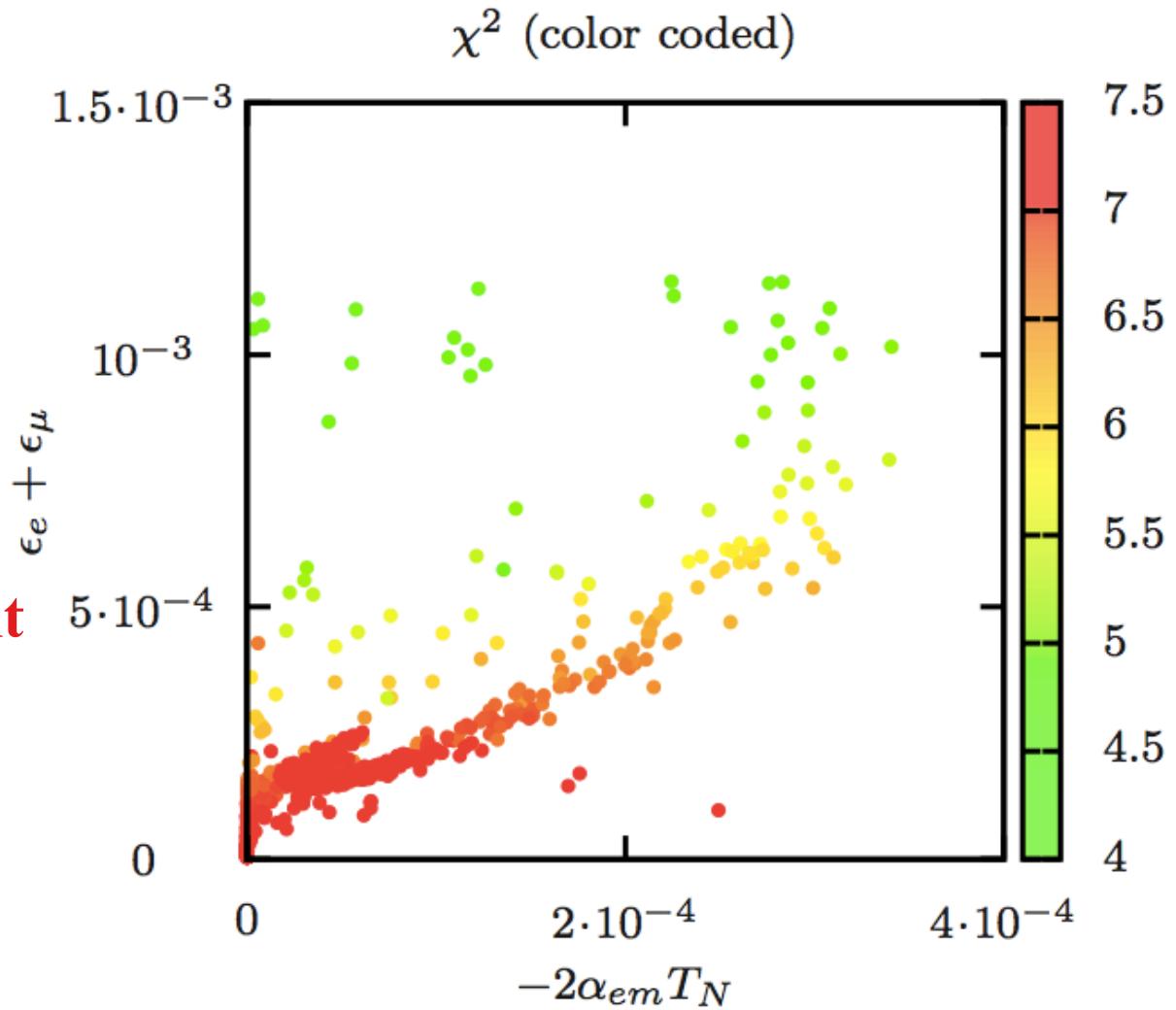
$$\chi^2_{\text{EWPO}} = \sum_i \frac{(O_i - O_{i,\text{SM}})^2}{(\delta O_i)^2 + (\delta O_{i,\text{SM}})^2}$$

# Some Examples

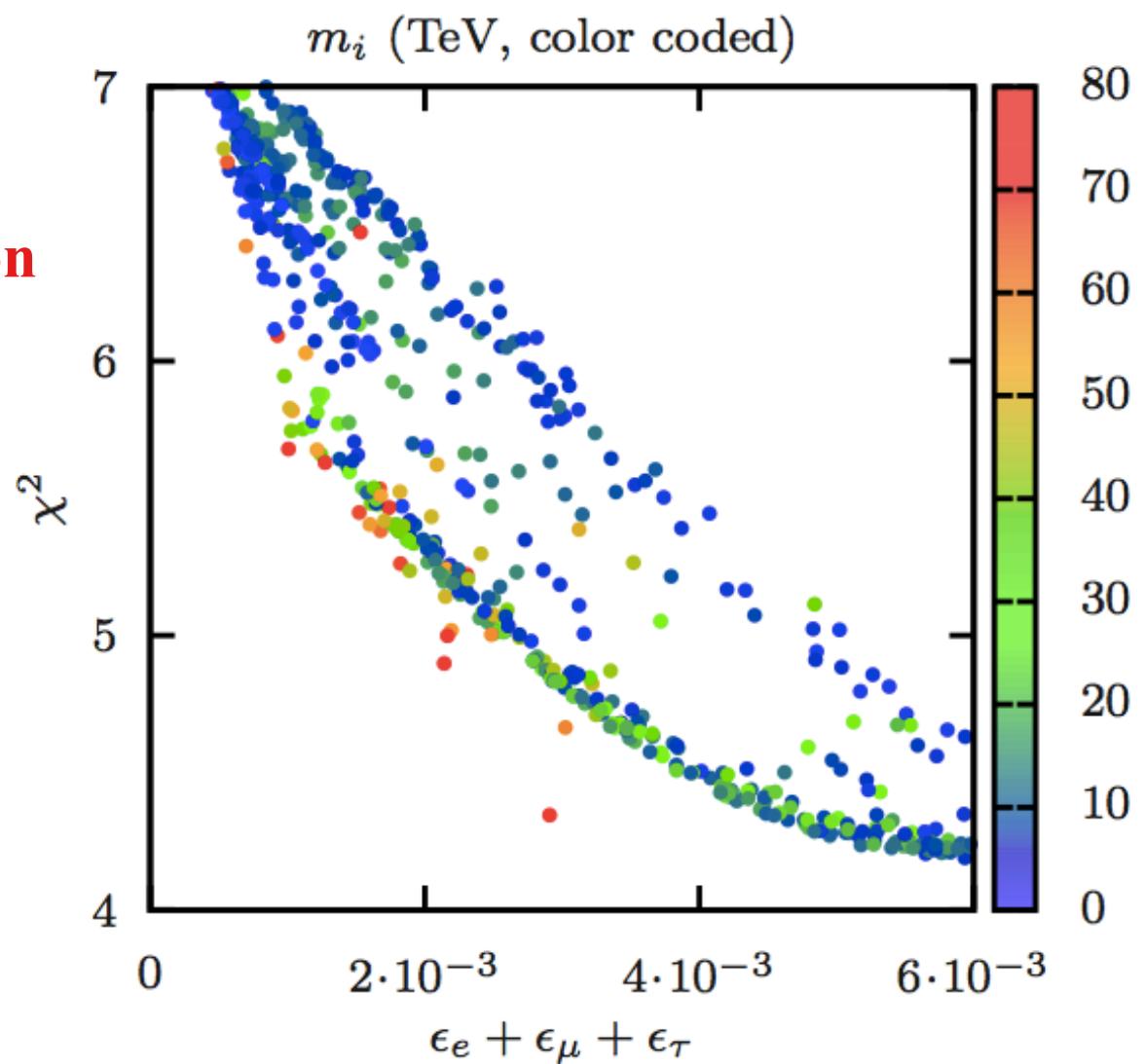
NH,  $\chi^2$  for 4 d.o.f.  
fit to all EWPOs

→ overall  $\chi^2/\text{dof}$  goes  
from  $\simeq 2$  to  $\simeq 1$   
for  $\varepsilon_e + \varepsilon_\mu \sim 10^{-3}$

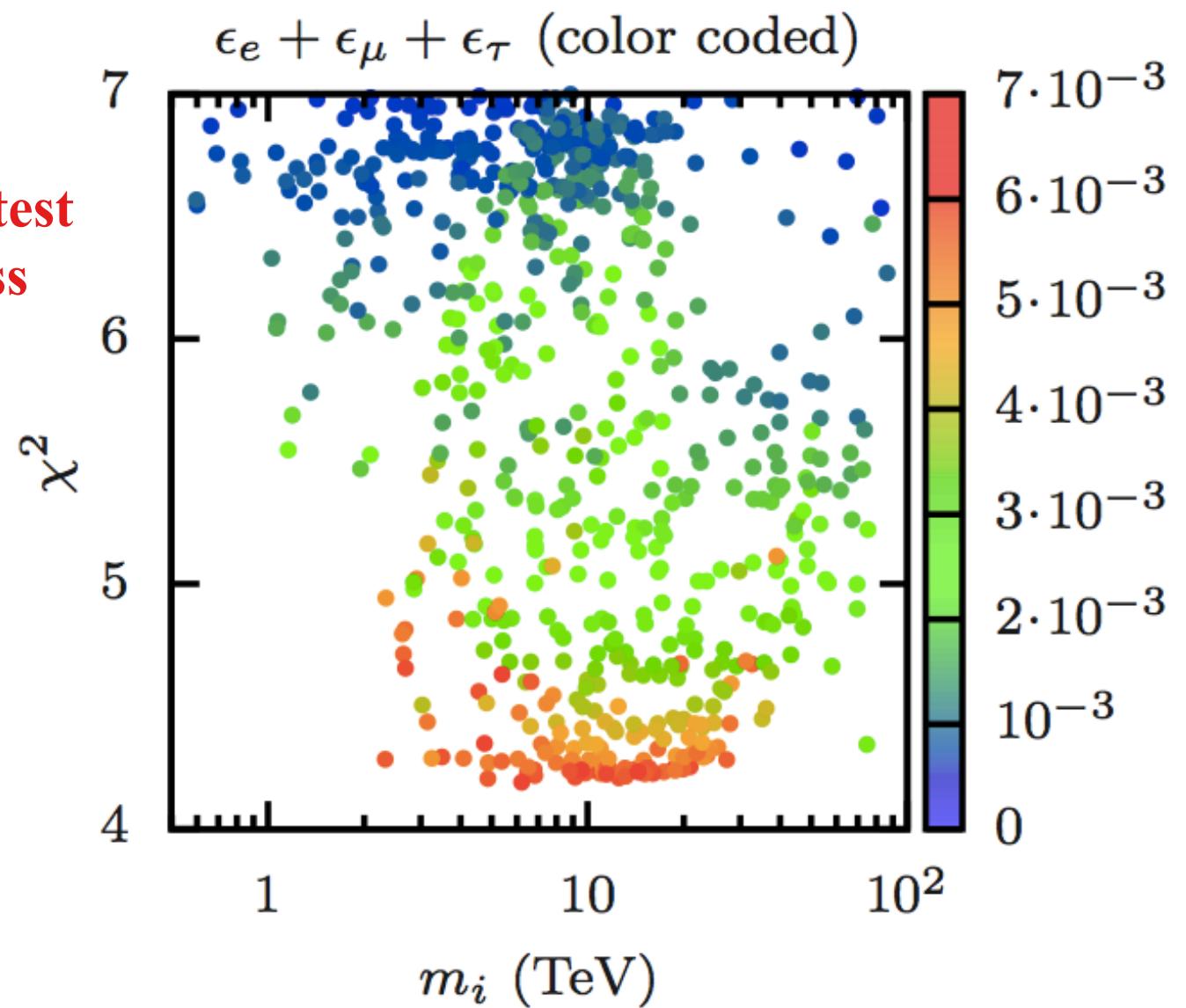
→ sterile ν's improve fit



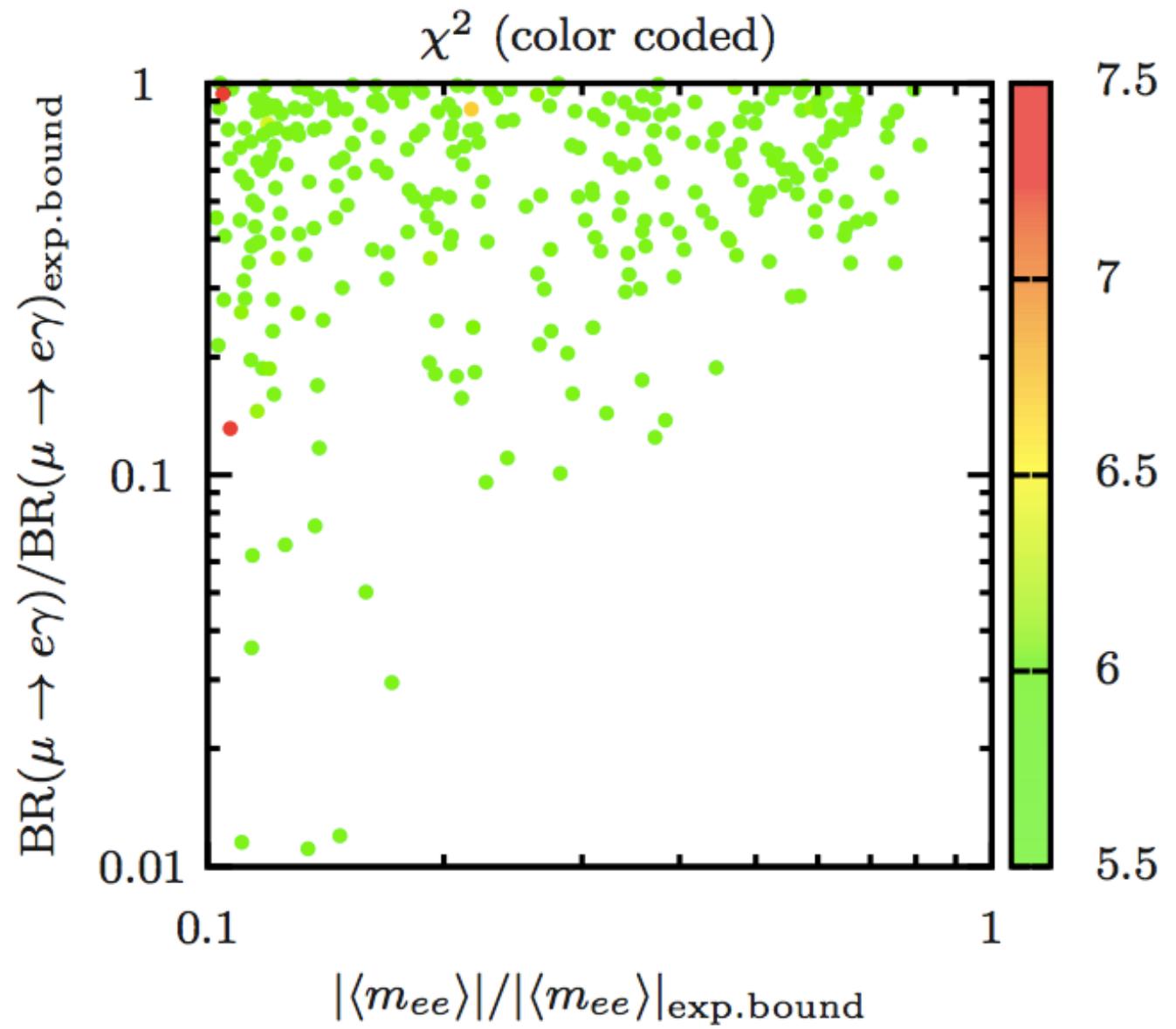
NH, the lightest heavy neutrino mass as function of  $\chi^2$  and  $\epsilon_e + \epsilon_\mu + \epsilon_\tau$  for 4dof



NH,  $\epsilon_e + \epsilon_\mu + \epsilon_\tau$  as a  
function of the lightest  
heavy neutrino mass  
for 4dof



IH,  
double beta decay  
and  $\mu \rightarrow e\gamma$  limits  
for 4dof



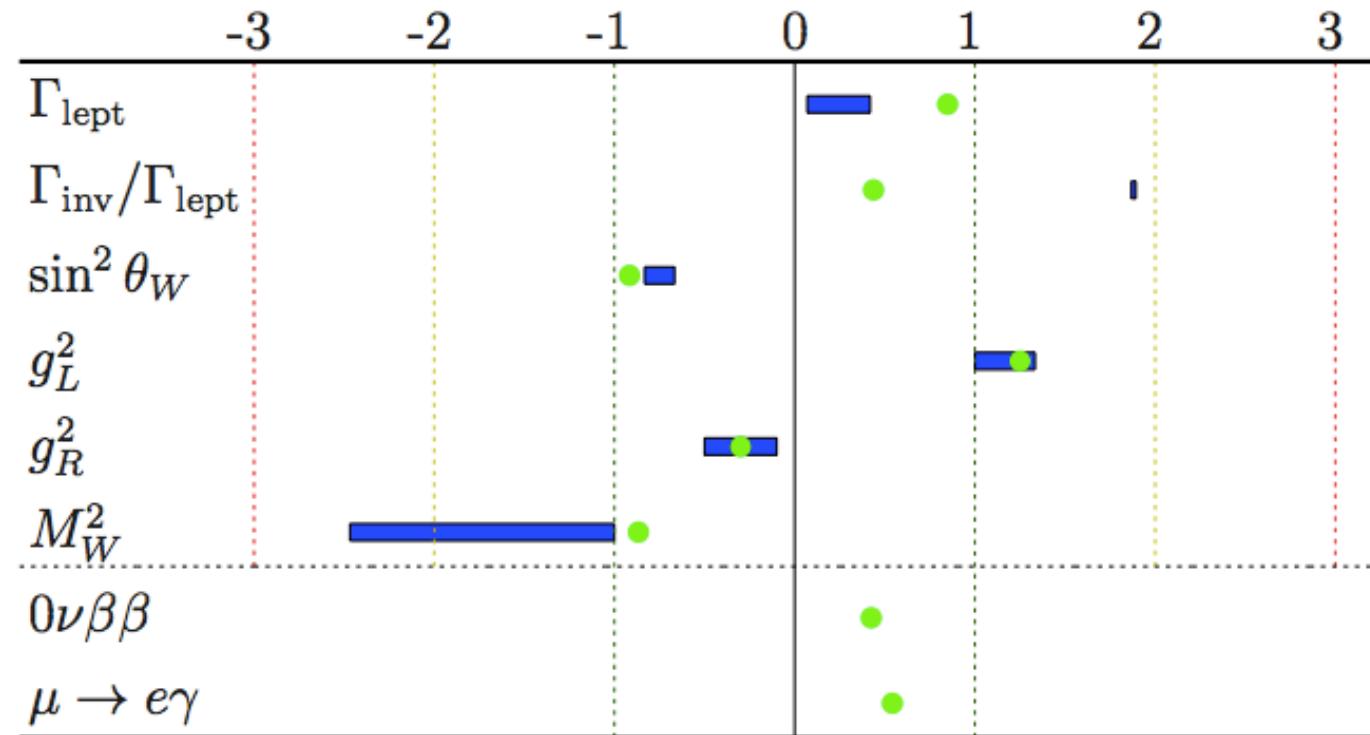
# Which Improvements

Agreement with data =  $0 + 1, 2, 3 \sigma$  ; assumed scenario: NH  
SM best fit  $\rightarrow$  blue bars

best fit with steril n's  $\rightarrow$  green points

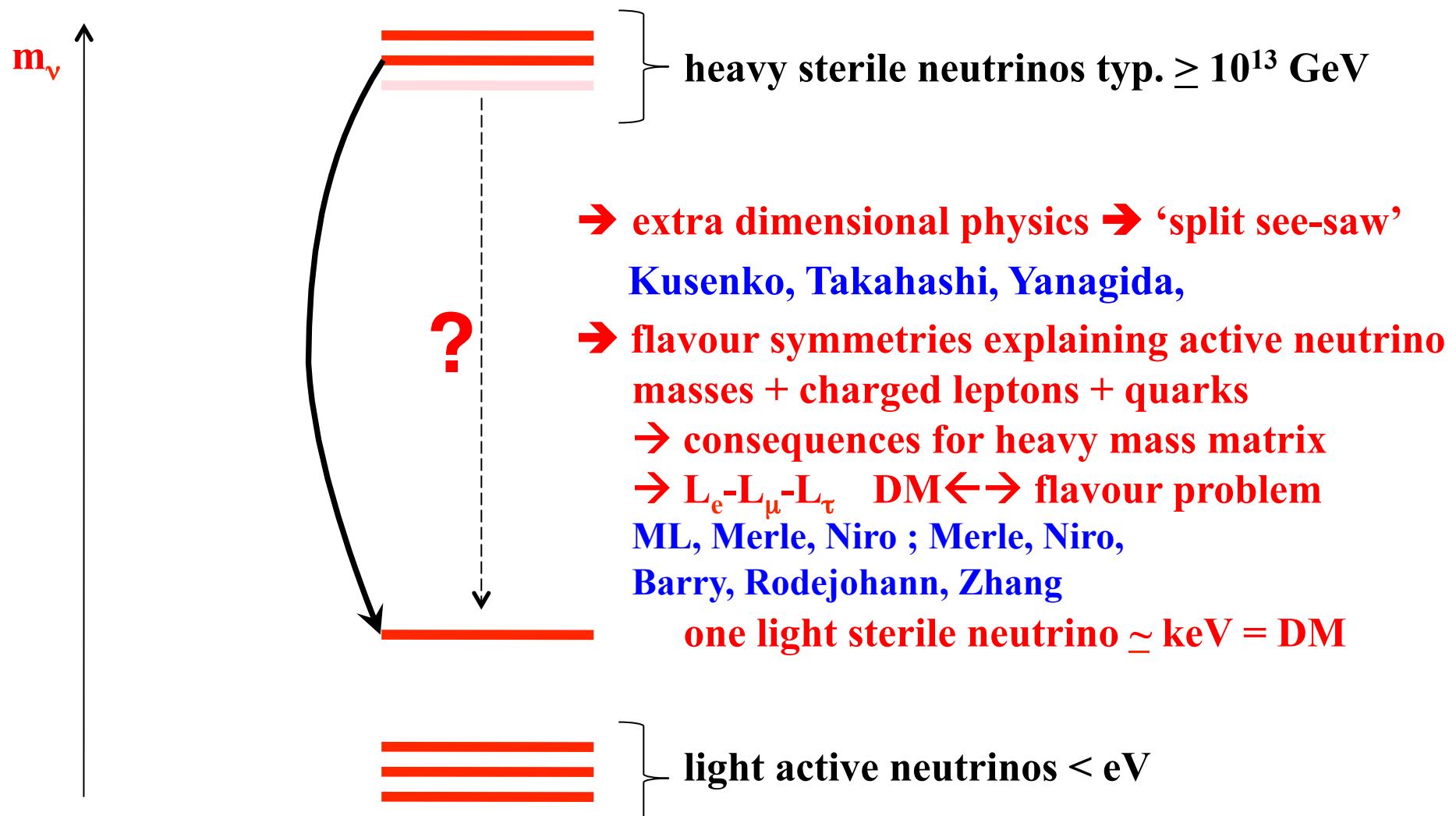
$M_1 = 20.3 \text{ TeV}$ ,  $M_2 = 14.1 \text{ TeV}$ ,  $M_3 = 21.0 \text{ TeV}$ ,

$\epsilon_e = 2.1 \cdot 10^{-3}$ ,  $\epsilon_\mu = 3.0 \cdot 10^{-6}$  and  $\epsilon_\tau = 4.5 \cdot 10^{-3}$



# Explaining keV-ish Sterile Neutrinos

Possible scenario: See-saw + a reason why 1 sterile  $\nu$  is light



# Conclusions

- Different indications / motivations for sterile neutrinos  
→ what about masses  $O(\text{TeV})$ ?
- Interesting cancellations between tree and loop effects of heavy sterile neutrinos
- Heavy sterile neutrinos improve global fits to all data and reduces tensions: deviation of Z-width, NuTeV, global
- Mass range from a few hundred GeV to 100 TeV
- $O(10^{-3})$  mixings
- Natural theoretical explanation of the spectrum required  
→ eV and keV states might also be explained  
→ interesting possibilities / combinations