

Non-Thermal Dark Matter

From the EW Phase Transition:

Heavy WIMP's and *Baby-Zillas*

Jose Miguel No (Sussex U.)

with Adam Falkowski

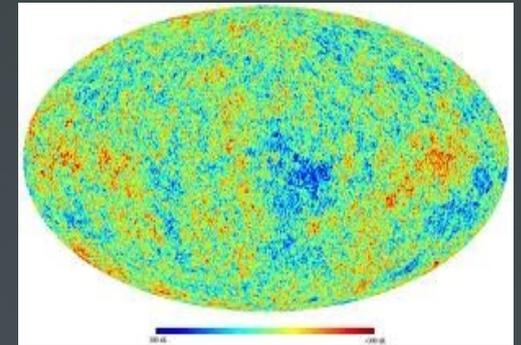
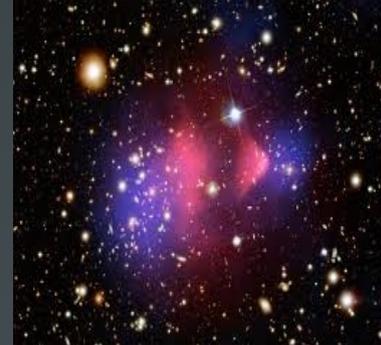
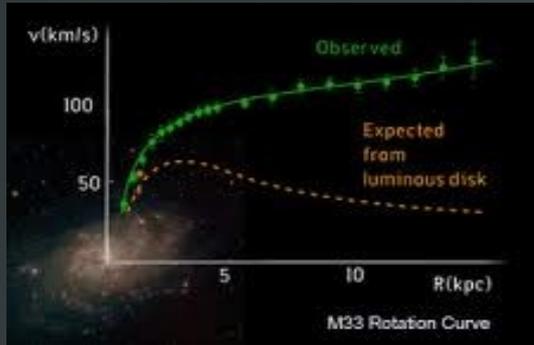
JHEP **1302** (2013) 034 (arXiv:1211.5615)

February 22nd 2013



Dark Matter in the Universe

Galaxy Rotation Curves, Gravitational Lensing (Bullet Cluster), CMB...

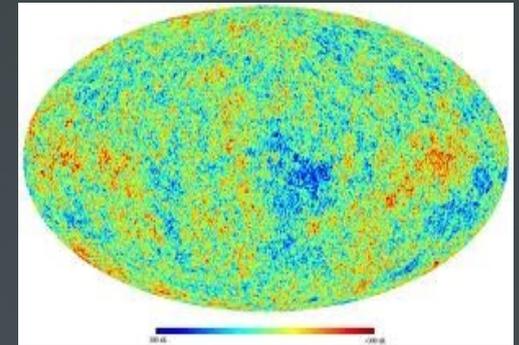
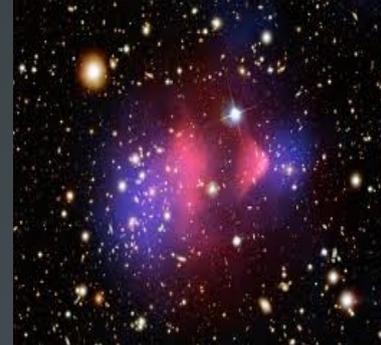
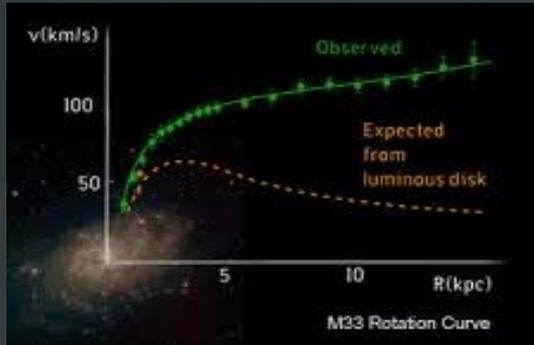


Particle Physics \rightarrow DM Candidates:  WIMP
Non-WIMP: KeV Sterile Neutrino, Axion, Gravitino, Wimpzillas...



Dark Matter in the Universe

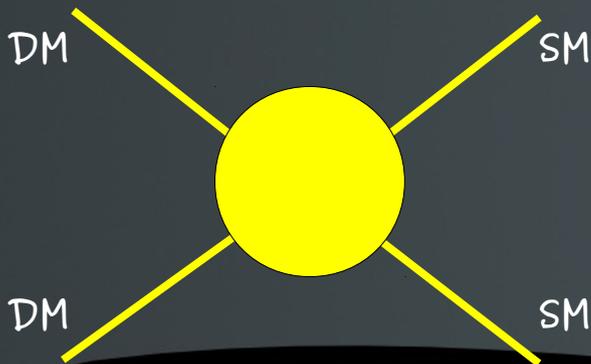
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Particle Physics \rightarrow DM Candidates: $\left\{ \begin{array}{l} \text{WIMP} \\ \text{Non-WIMP: KeV sterile Neutrino, Axion, Gravitino, Wimpzillas...} \end{array} \right.$

“WIMP Paradigm”

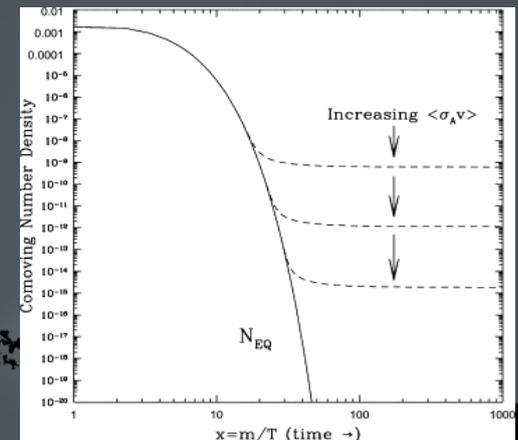
SUSY Neutralino, KK Dark Matter, Little Higgs (w. T-Parity), Singlet Scalar (Hidden Sector DM), Inert Doublet Model...



Dark Matter Annihilation



Thermal Freeze-out



Dark Matter in the Universe

Other Appealing Dark Matter Production Mechanisms...

Non-Thermal DM Production

⇒ Non-thermal DM Disconnected from EW Physics?

Axions, WIMPZILLAS, Modulus decay products...

DM Direct Detection? DM Indirect Detection?



Dark Matter in the Universe

Other Appealing Dark Matter Production Mechanisms...

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DM Direct Detection? DM Indirect Detection?

⇒ EW Phase Transition May Be Out-of-Equilibrium Process

+

⇒ Most DM candidates at the EW Scale Couple to Higgs

Non-Thermal DM Production at the EW Phase Transition



Outline

- Dynamics of the Electroweak Phase Transition
 - ⇒ Bubble 1st Order Phase Transitions: Nucleation & Expansion
 - ⇒ Bubble Expansion Velocities: Runaway Bubbles
 - ⇒ Bubble Collisions
- Particle Production in Bubble Collisions
- Non-Thermal Dark Matter from the EW Phase Transition
 - ⇒ Multi-TeV WIMP's
 - ⇒ Baby-Zillas



Electroweak Phase Transition

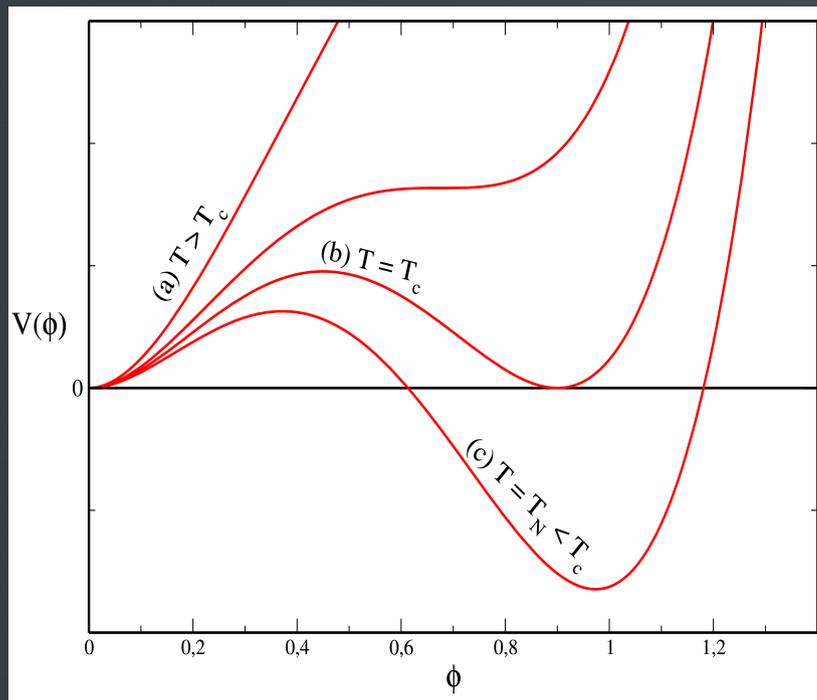
Universe Expands Adiabatically \Rightarrow Equilibrium Thermal Field Theory **OK**

Finite-T Effective Potential $V(\phi, T)$ for the Higgs

$$V(\phi, T) \approx (a T^2 - \mu^2) \phi^2 - b T \phi^3 + \lambda \phi^4$$

1st Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$ Discontinuous

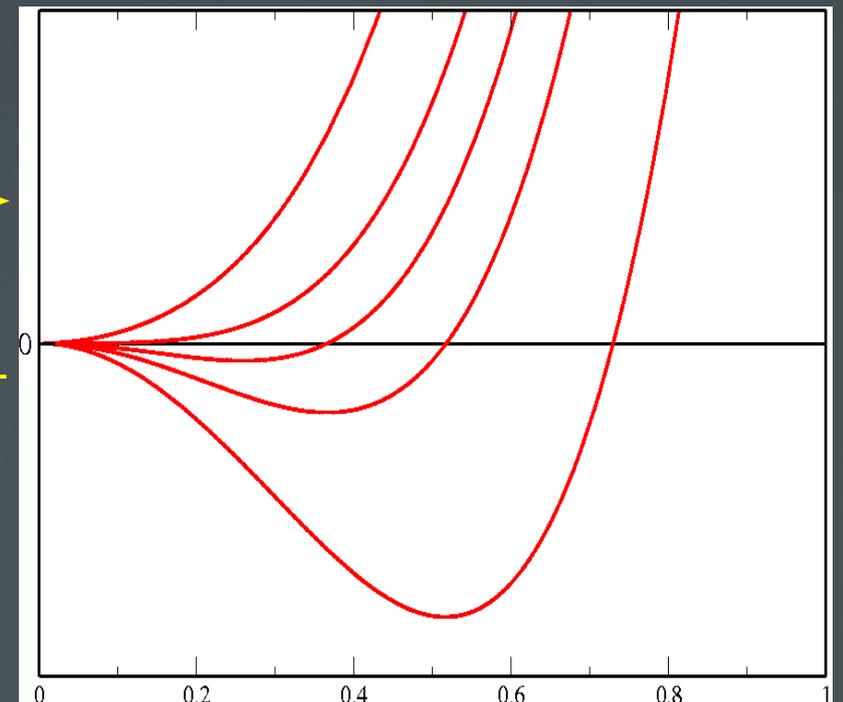


Larger m_h

New Bosons

2nd Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$ Continuous



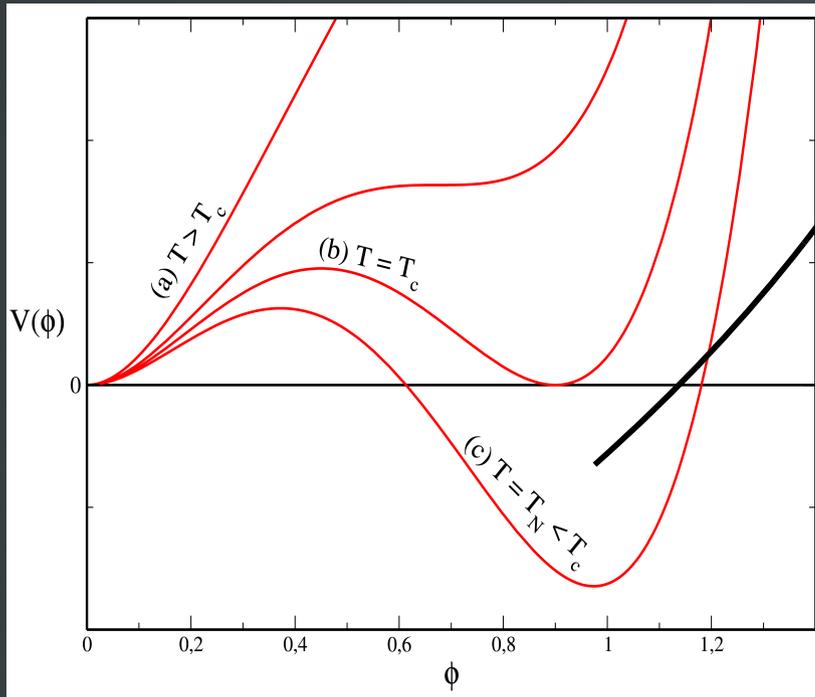
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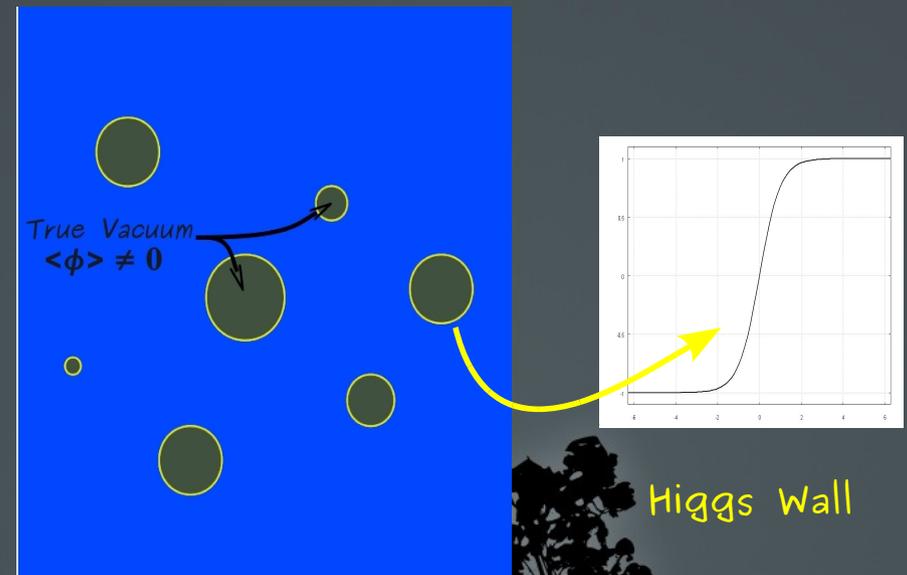


Nucleation of True Vacuum Bubbles
(in False Vacuum Sea)

*J. S. Langer, Ann. Phys. **54** (1969) 258*

*S. R. Coleman, Phys. Rev. D **15** (1977) 2929*

*A. D. Linde, Nucl. Phys. B **216** (1983) 421*



How do Higgs Bubbles Expand?

Pressure $\sim \Delta\mathcal{F}$

$$\partial_\mu \partial^\mu \phi + \frac{\partial \mathcal{F}}{\partial \phi} - \mathcal{K}(\phi) = 0$$

$$\mathcal{K}(\phi) = - \sum_i \frac{dm_i^2}{d\phi} \int \frac{d^3 p}{(2\pi)^3 2E_i} \delta f_i(p)$$

Friction

Particles Gain Mass
When Crossing Wall

Particle Distributions $f_a(p)$
Away from Equilibrium
Close to Wall

$$F_\eta(v_w) = \int_{-\infty}^{\infty} dz \frac{d\phi}{dz} \sum_i \frac{dm_i^2}{d\phi} \int \frac{d^3 p}{(2\pi)^3 2E_i} \delta f_i(p) \simeq \eta v_w$$

Friction Reaches a Terminal Value! (F_{\max})

D. Bodeker and G. Moore, JCAP 0905 (2009) 009



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→ If Pressure Exceeds F_{\max} ⇒ Accelerated Bubble Expansion
"Runaway Bubbles"

D. Bodeker and G. Moore, JCAP 0905 (2009) 009

→ Energy of the EW Phase Transition Stored on Higgs Bubble Walls

J. R. Espinosa, T. Konstandin, J. M. N. and G. Servant, JCAP 1004 (2010) 028



How do Higgs Bubbles Expand?

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Runaway Bubbles are Possible in Extensions of the SM

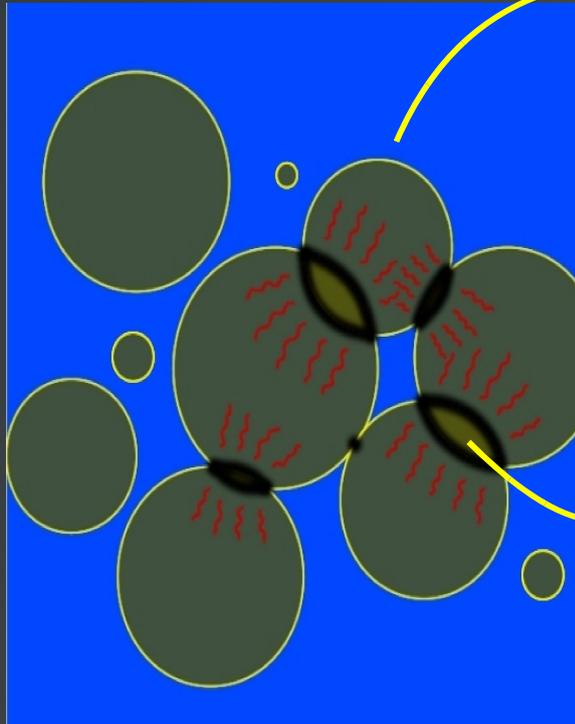
*D. Bodeker and G. Moore, JCAP **0905** (2009) 009*

*J. R. Espinosa, T. Konstandin, J. M. N. and G. Servant, JCAP **1004** (2010) 028*

S. Huber and M. Sopena, arXiv: 1302:1044



The End of the Phase Transition: Bubble Collisions



Runaway Bubbles May Reach $\gamma_{\max} \sim 10^{15}$

$$\gamma_w^{\max} \sim \beta^{-1} \frac{M_{\text{pl}}}{v}$$

Bubble Collisions Release Energy in Walls

- Gravitational Wave Production
- Particle Production



Particle Production from Bubble Collisions

Higgs Wall \Rightarrow Classical Background



Particle Production (in a space-time dependent Classical Background)

Background $\phi(\mathbf{r},t)$ (Bubble wall Profile) + Quantum fields coupled to it

- 1 Modelling the Background: a Bubble Collision
- 2 Particle Production in Given Background



① Modelling the Background: a Bubble Collision

⇒ Large Bubbles at time of Collision  Planar Limit

$$\boxed{(\partial_t^2 - \partial_z^2) \phi(z, t) = -\frac{\partial V(\phi)}{\partial \phi}}$$

+ Initial condition (2 planar bubble walls far away from each other and moving in opposite directions)

*S. Hawking, I. Moss and A. Steward, Phys. Rev. D **26** (1982) 2681*



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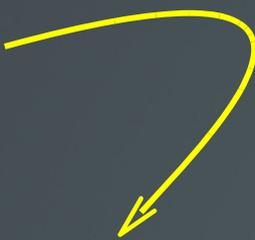
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*S. Hawking, I. Moss and A. Steward, Phys. Rev. D **26** (1982) 2681*

Each Bubble Wall interpolates between two minima of potential:

outside of wall → $\boxed{(\partial_t^2 - \partial_z^2) \phi(z, t) = 0}$ 

$$\boxed{\phi(z, t) = \frac{v}{2} \left[2 + \text{Tanh} \left(\gamma_w \frac{z+t}{l_w} \right) - \text{Tanh} \left(\gamma_w \frac{z-t}{l_w} \right) \right]}$$

Approximate solution Before the Collision ($t < 0$)

① Modelling the Background: a Bubble Collision

Elastic vs Inelastic Collisions:

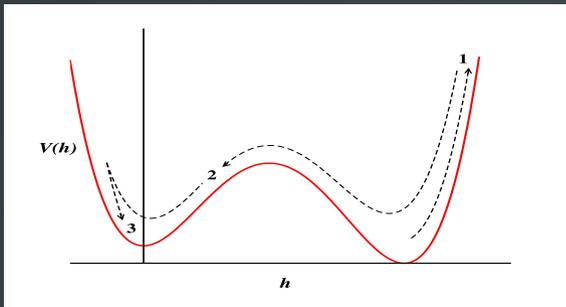
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T. Konstandin and G. Servant, JCAP **1107** (2011) 024

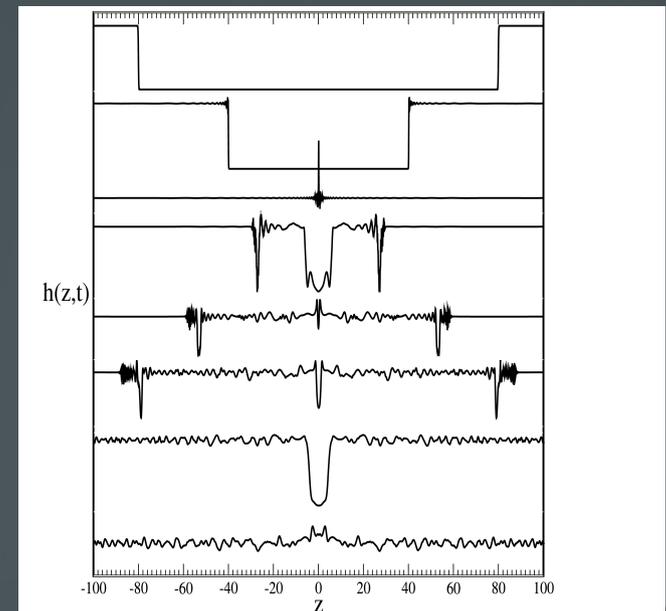
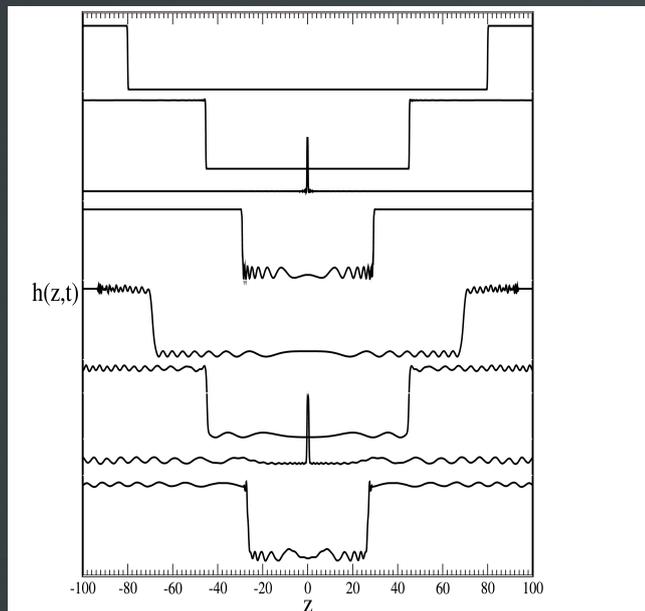
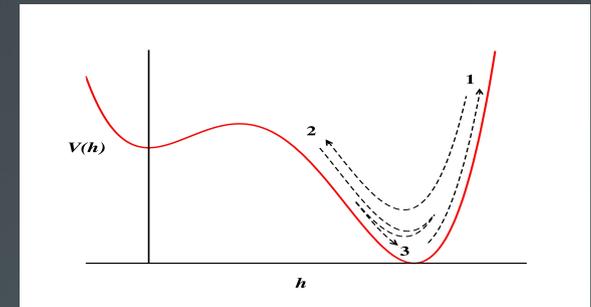
A. Falkowski and J. M. N. JHEP **1302** (2013) 034

$$(\partial_t^2 - \partial_z^2) \phi(z, t) = -\frac{\partial V(\phi)}{\partial \phi}$$

Elastic



Inelastic



① Modelling the Background: a Bubble Collision

Elastic vs Inelastic Collisions:

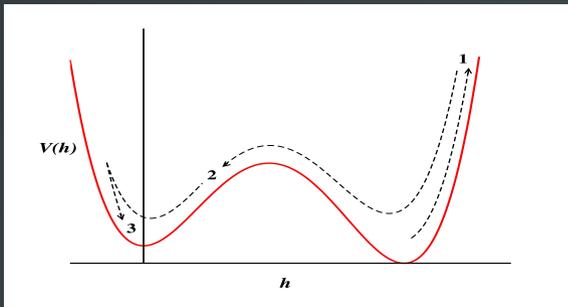
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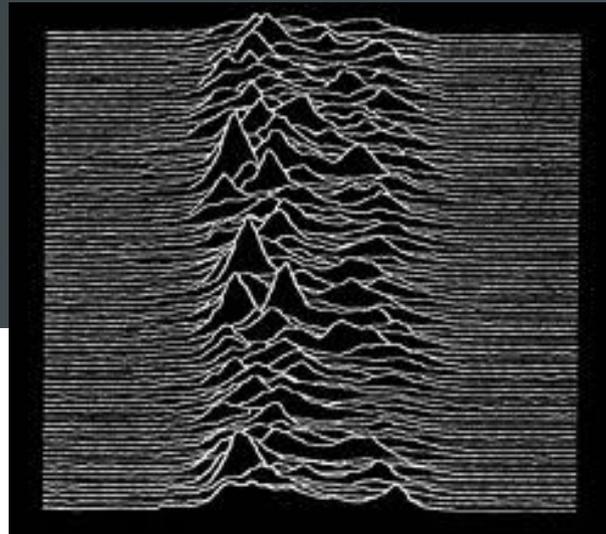
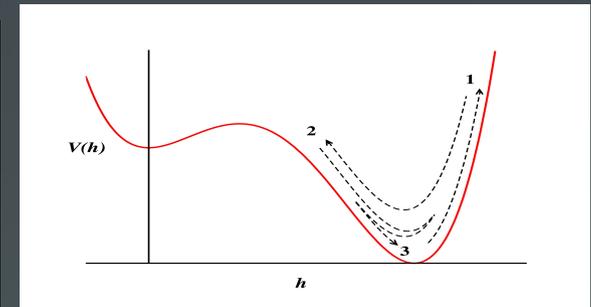
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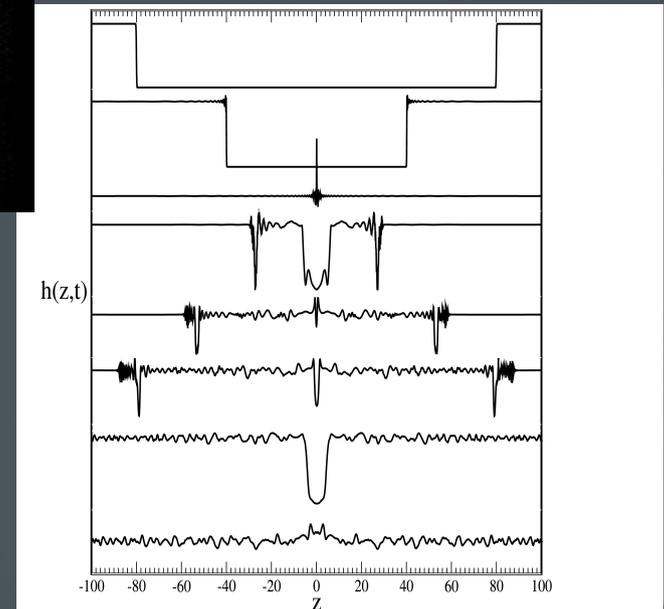
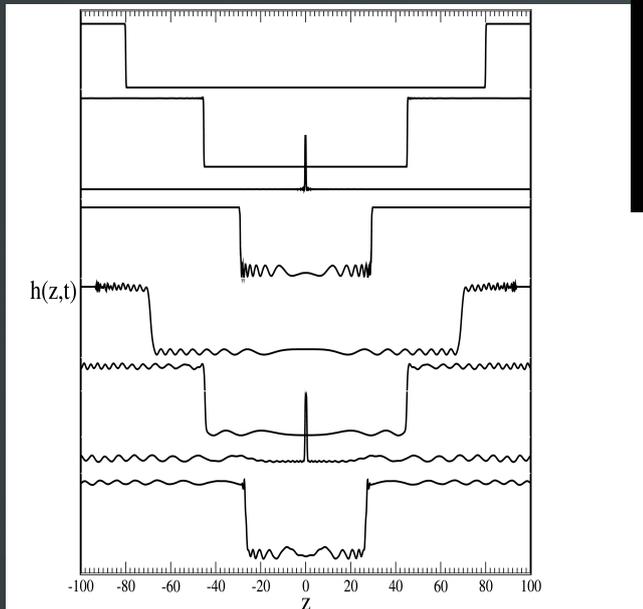
Elastic



Inelastic



Although this type of figure had appeared earlier!



① Modelling the Background: a Bubble Collision

Extreme Cases:

$$(\partial_t^2 - \partial_z^2) \phi(z, t) = -\frac{\partial V(\phi)}{\partial \phi}$$

Perfectly Elastic Collision

S. Hawking, I. Moss and A. Steward, Phys. Rev. D **26** (1982) 2681

$$\phi(z, t) = \frac{v}{2} \left[2 + \text{Tanh} \left(\gamma_w \frac{z - |t|}{l_w} \right) - \text{Tanh} \left(\gamma_w \frac{z + |t|}{l_w} \right) \right]$$

Totally Inelastic Collision

A. Falkowski and J. M. N. JHEP **1302** (2013) 034

$$V(\phi) \simeq \frac{m_h^2}{2} (\phi - v)^2 = \frac{m_h^2}{2} \delta\phi^2$$

$$\phi(z, t) = \begin{cases} 0 & \text{if } v_w t < z < -v_w t & t < 0, \\ 0 & \text{if } -v_w t < z < v_w t & t > 0, \\ v & \text{Otherwise,} \end{cases}$$

$$\phi(z, t > 0) = v \left[1 + \frac{l_w}{\gamma_w} \int_0^\infty dp_z \frac{p_z}{\sqrt{p_z^2 + m_h^2}} \frac{\text{Cos}(p_z z)}{\text{Sinh} \left(\frac{\pi l_w p_z}{2 \gamma_w} \right)} \text{Sin} \left(\sqrt{p_z^2 + m_h^2} t \right) \right]$$



① Modelling the Background: a Bubble Collision

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Of course, particle production will affect $\phi(\mathbf{r}, t)$  Backreaction



② Particle Production

Number of Produced Particles:

$$\mathcal{N} = 2 \text{Im}(\Gamma[\phi])$$

Effective action:

Generator of 1PI Green functions

$$\Gamma[\phi] = \frac{1}{2} \int d^4x_1 d^4x_2 \phi(x_1) \phi(x_2) \Gamma^{(2)}(x_1 - x_2) + \dots$$

$$\mathcal{N} = \int \frac{d^4p}{(2\pi)^4} \text{Im} \left(\tilde{\Gamma}^{(2)}(p^2) \right) \int d^4x_1 d^4x_2 \phi(x_1) \phi(x_2) e^{ip(x_1 - x_2)}$$

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$$\mathcal{N} = \int \frac{d^4p}{(2\pi)^4} \text{Im} \left(\tilde{\Gamma}^{(2)}(p^2) \right) \left| \tilde{h}(p) \right|^2$$

Planar
Bubble Walls

$$\frac{\mathcal{N}}{A} = 2 \int \frac{dp_z d\omega}{(2\pi)^2} \left| \tilde{h}(p_z, \omega) \right|^2 \text{Im} \left(\tilde{\Gamma}^{(2)}(\omega^2 - p_z^2) \right)$$

R. Watkins and L. M. Widrow, Nucl. Phys B374 (1992) 446

② Particle Production

$$\frac{\mathcal{N}}{A} = \frac{1}{2\pi^2} \int_0^\infty d\chi f(\chi) \text{Im} \left(\tilde{\Gamma}^{(2)}(\chi) \right)$$

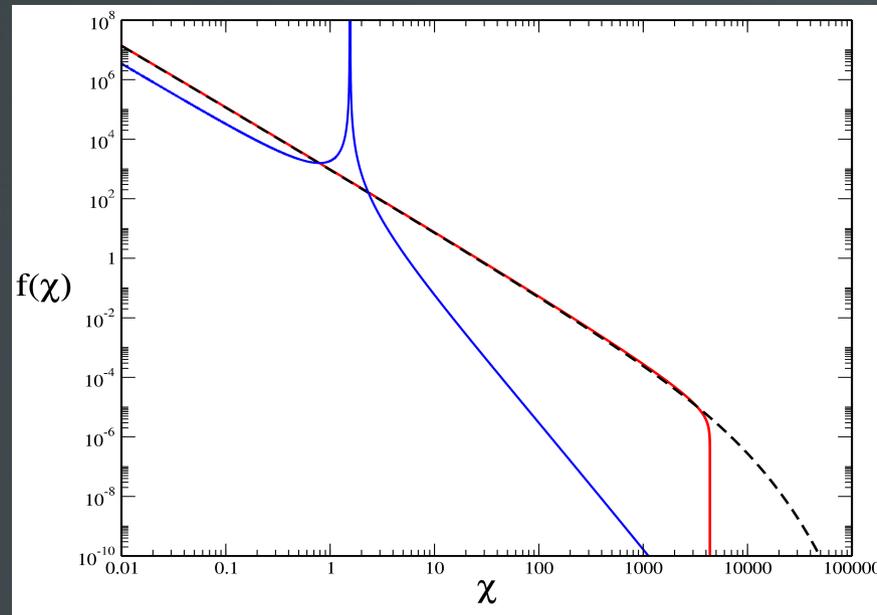
Particle Production Efficiency

Perfectly Elastic

$$f_{\text{PE}}(\chi) = \frac{16 v^2 \text{Log} \left[\frac{2 \left(\frac{\gamma w}{l_w} \right)^2 - \chi + 2 \frac{\gamma w}{l_w} \sqrt{\left(\frac{\gamma w}{l_w} \right)^2 - \chi}}{\chi} \right]}{\chi^2} \Theta \left[\left(\frac{\gamma w}{l_w} \right)^2 - \chi \right]$$

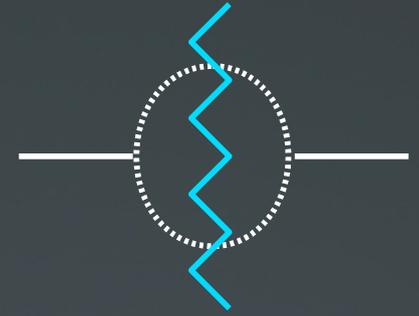
Totally Inelastic

$$f_{\text{TI}}(\chi) = 4 v^2 m_h^4 \frac{\text{Log} \left[\frac{2 \left(\frac{\gamma w}{l_w} \right)^2 + \chi + 2 \frac{\gamma w}{l_w} \sqrt{\left(\frac{\gamma w}{l_w} \right)^2 + \chi}}{\chi} \right]}{\chi^2 \left[(\chi - m_h^2)^2 + m_h^6 \frac{l_w^2}{\gamma^2} \right]}$$



② Particle Production

$$\frac{\mathcal{N}}{A} = \frac{1}{2\pi^2} \int_0^\infty d\chi f(\chi) \text{Im} \left(\tilde{\Gamma}^{(2)}(\chi) \right)$$

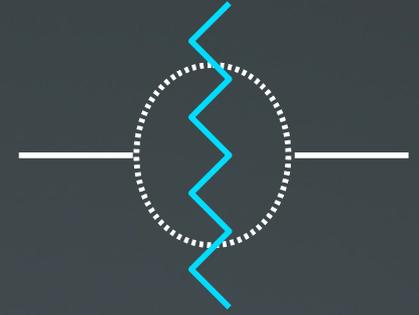


$$\text{Im} \left(\tilde{\Gamma}^{(2)}(\chi) \right) = \frac{1}{2} \sum_{\alpha} \int d\Pi_{\alpha} |\overline{\mathcal{M}}(h \rightarrow \alpha)|^2 \Theta[\chi - \chi_{\min}]$$



② Particle Production

$$\frac{\mathcal{N}}{A} = \frac{1}{2\pi^2} \int_0^\infty d\chi f(\chi) \text{Im} \left(\tilde{\Gamma}^{(2)}(\chi) \right)$$



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⇒ Production of Scalars, Fermions, Vector Bosons.

$$\text{Im} \left[\tilde{\Gamma}^{(2)}(\chi) \right]_S = \frac{\lambda_s^2 v^2}{8\pi} \sqrt{1 - 4\frac{m_s^2}{\chi}} \Theta(\chi - 4m_s^2)$$

$$\text{Im} \left[\tilde{\Gamma}^{(2)}(\chi) \right]_f = \frac{m_f^2}{4\pi v^2} \chi \left(1 - \frac{4m_f^2}{\chi} \right)^{\frac{3}{2}} \Theta(\chi - 4m_f^2)$$

$$\text{Im} \left[\tilde{\Gamma}^{(2)}(\chi) \right]_V = \frac{\lambda_V^2 m_V^2}{8\pi} \left(3 - \frac{\chi}{m_V^2} + \frac{\chi^2}{4m_V^4} \right) \sqrt{1 - 4\frac{m_V^2}{\chi}} \Theta(\chi - 4m_V^2)$$

→ If effective coupling to the Higgs, production suppressed.

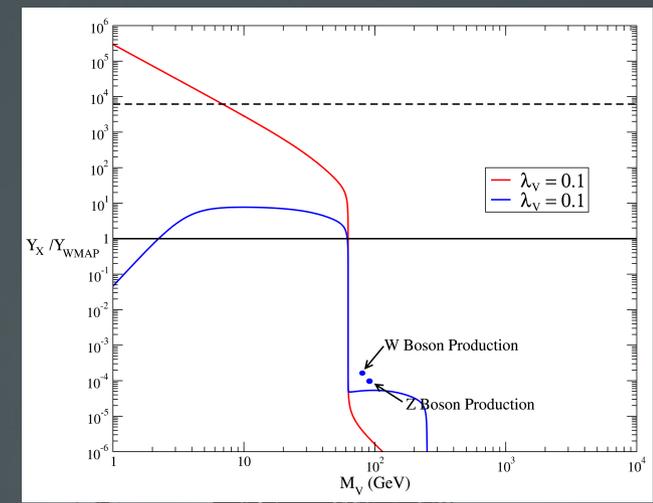
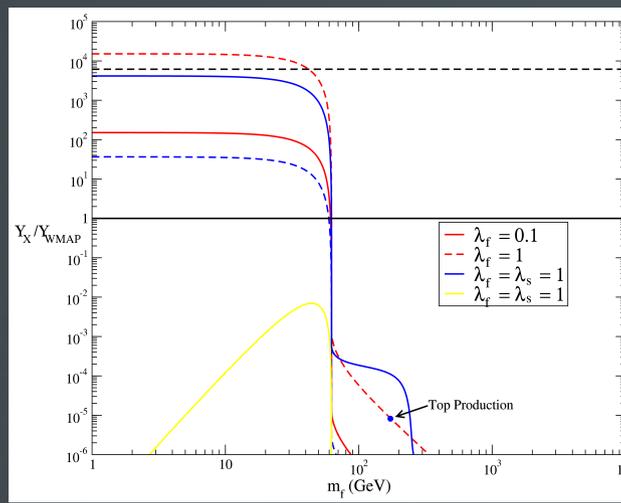
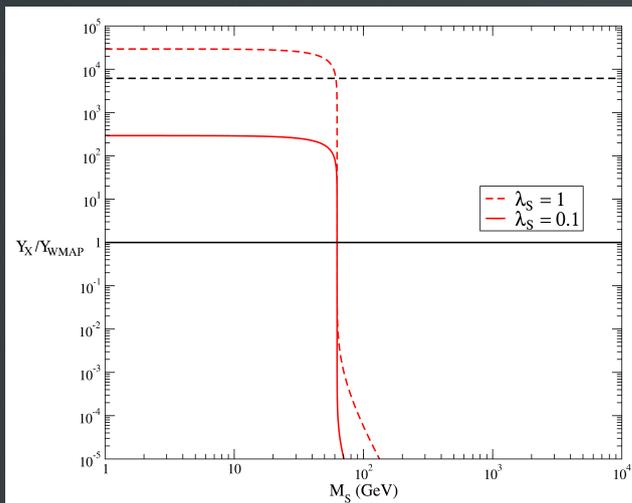
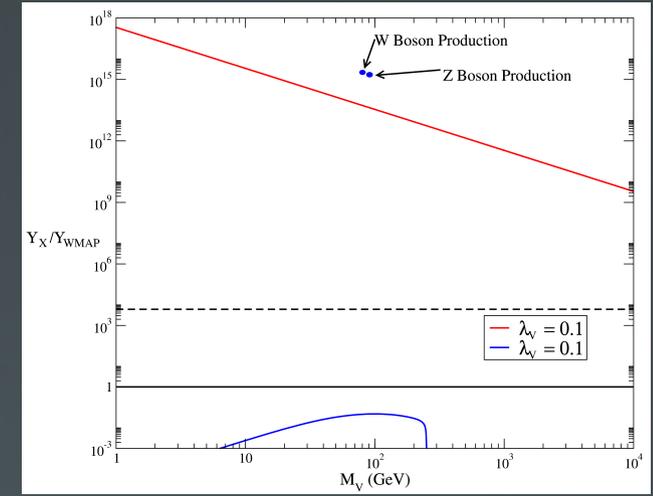
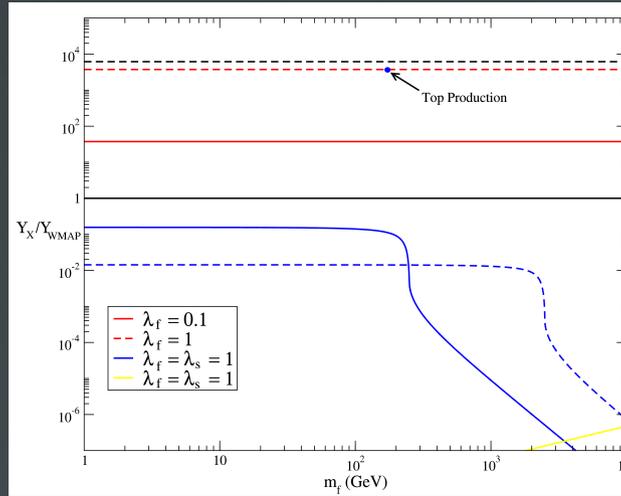
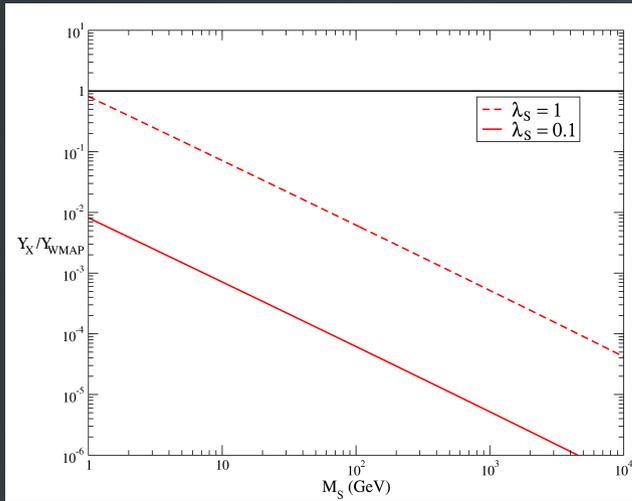
Bubble collisions excite very massive field modes $p^2 \gg \Lambda^2$ (eff. description breaks down)

→ Possible Issues with Backreaction.



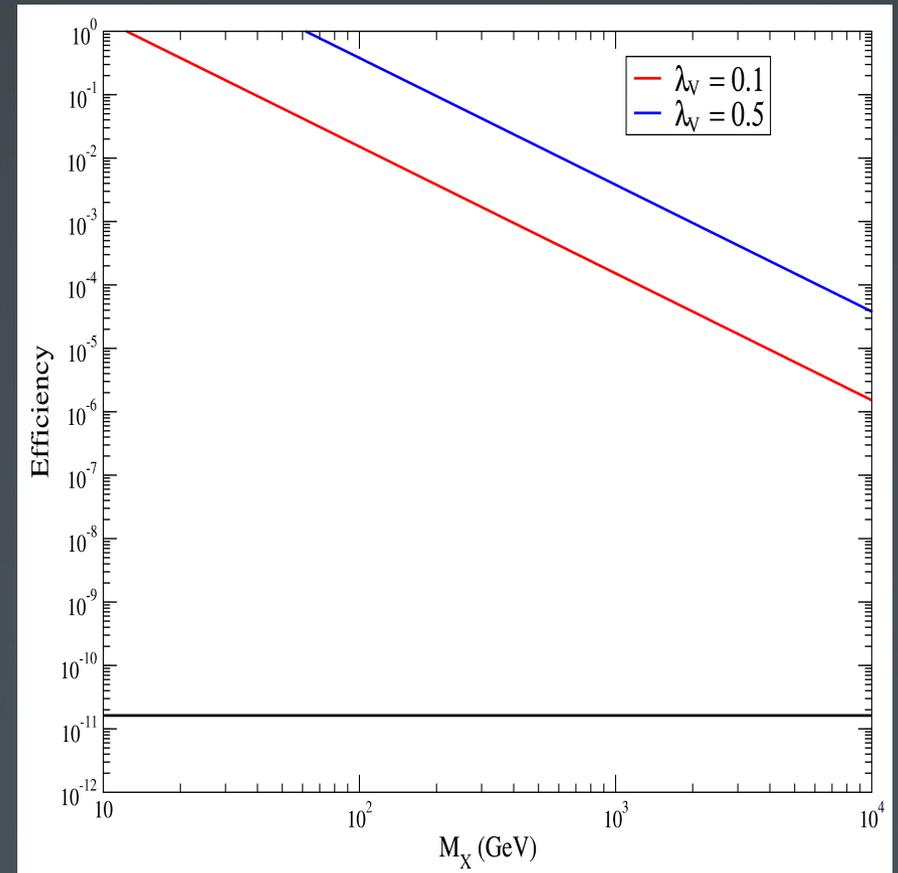
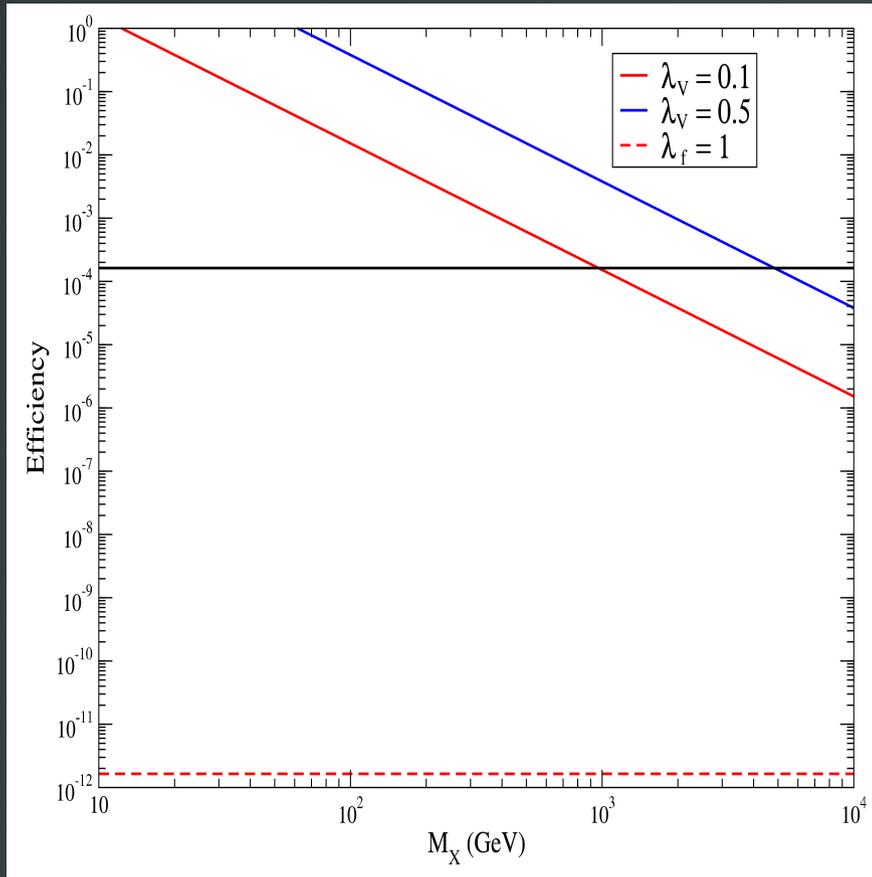
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⇒ Production of Scalars, Fermions, Vector Bosons.



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Non-Thermal Dark Matter

→ Relatively Elastic Collisions

→ Dark Matter is a Heavy Vector Boson V_μ (with direct coupling to Higgs)

① $m_V \sim \text{TeV} \Rightarrow V_\mu$ is a heavy WIMP

Possibilities: Kaluza-Klein DM, Little Higgs DM...

② $m_V \gg \text{TeV} \Rightarrow$ Baby-Zillas



Non-Thermal *muti-TeV* WIMPs

$$\mathcal{L}_V = \frac{1}{2}m_V^2 V_\mu V_\mu + \lambda_V v_T h V_\mu V_\mu$$

Non-thermal DM
Production

Thermal Abundance

$$\left[\frac{\lambda_V}{M_V(\text{TeV})} \right]_{\text{WMAP}} \approx 0.3$$

Direct Detection

$$\sigma_{VN \rightarrow VN} \approx 4.2 \cdot 10^{-44} \text{cm}^2 \left[\frac{\lambda_V}{M_V(\text{TeV})} \right]^2 < M_V \cdot 2.2 \cdot 10^{-44} \text{cm}^2$$

XENON bound



Non-Thermal *multi-TeV* WIMPs

$$\mathcal{L}_V = \frac{1}{2} m_V^2 V_\mu V_\mu + \lambda_V v_T h V_\mu V_\mu$$

Non-thermal DM
Production
 $\gamma_w \sim 10^8$

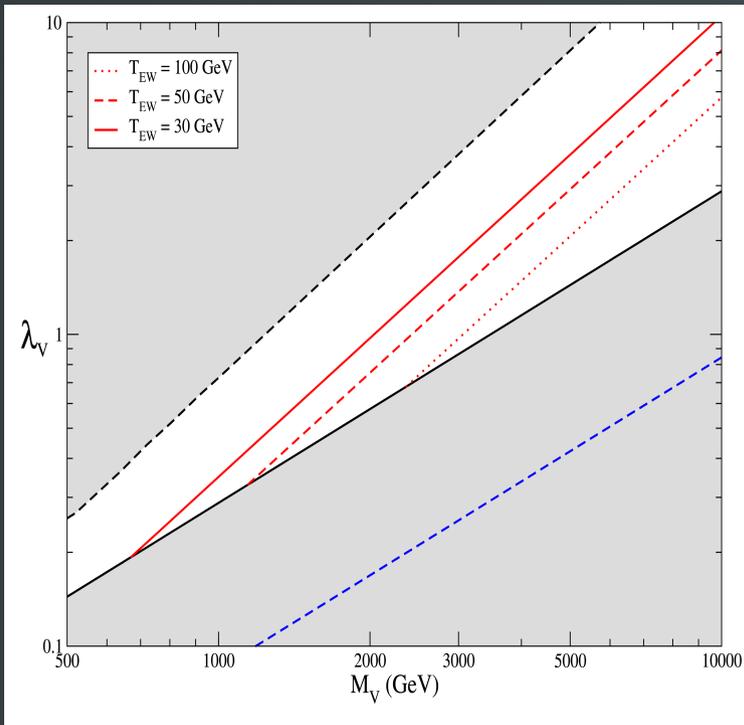
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XENON bound



$$\frac{dY}{dz} = -\alpha \frac{\langle \sigma v \rangle M_{\text{Pl}} M_V}{z^2} Y(z)^2 \rightarrow \frac{dy}{dz} = -\frac{1}{z^2} y^2(z) \rightarrow \frac{1}{y(z)} \frac{1}{y(z_{\text{EW}})} = \frac{1}{z_{\text{EW}}} \frac{1}{z}$$

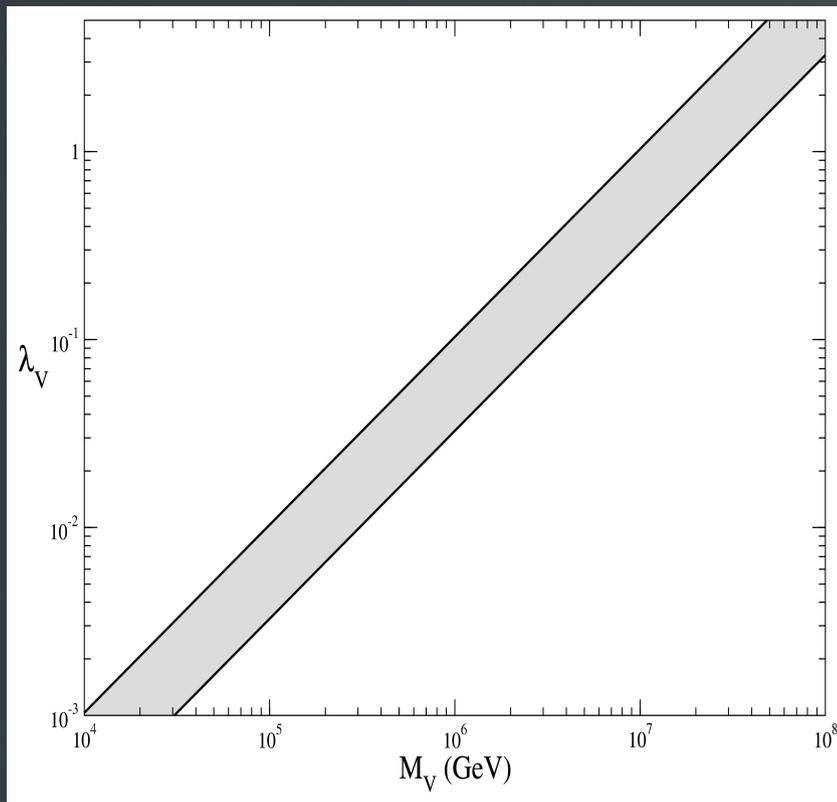
$$y(\infty) \simeq z_{\text{EW}}$$

$$\langle \sigma v \rangle = \langle \sigma v \rangle_{\text{WMAP}} \frac{T_{\text{fo}}}{T_{\text{EW}}}$$

Non-thermal
"WIMP miracle"

Non-Thermal *Baby-Zillas*

→ If $\gamma_w \sim \gamma_{\max} \sim 10^{15}$, possible to produce $m_\nu \gg \text{TeV}$



$$M_{\text{GUT}} \gg M_V \gg v$$

Baby-Zillas

$$T_{\text{RH}} < M_V/10$$

Low Reheating T after
Inflation Required

Conclusions

→ Efficient Particle Production in Strongly 1st Order EW Phase Transitions

→ Possible Non-thermal Production of Vector DM

⇒ Multi-TeV WIMP's Kaluza-Klein DM, Little Higgs DM...

⇒ Baby-Zillas

