Heavy Higgs Sector of the Two Higgs Doublet Model in the Natural Standard Model Alignment Limit

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with A. Pilaftsis, JHEP **1412**, 024 (2014) [arXiv:1408.3405 [hep-ph]]

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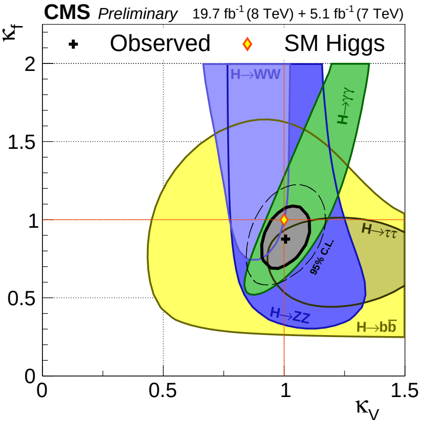
Outline

- Introduction
- Natural Alignment
- Maximally Symmetric 2HDM
- Collider Phenomenology
- Conclusion

A Higgs or *the* Higgs?



Unique opportunity for probing New Physics through the Higgs portal:

- Precision Higgs Study (Higgcision).
- Search for additional Higgses.

Why more Higgses?

Several theoretical motivations to go beyond the SM Higgs sector.

- Electroweak Baryogenesis
  - Additional sources of CP violation
  - Strong first order phase transition
- Dark Matter
- Supersymmetry
- Why not?

Two Higgs Doublets

- Any scalar sector in a local $SU(2) \times U(1)$ gauge theory must be consistent with $\rho_{\rm exp} = 1.0004^{+0.0003}_{-0.0004}$. [PDG '14]

- With n Higgs multiplets Φ_i (with $i = 1, 2, ..., n$):

$$\rho_{\rm tree} = \frac{\sum_{i=1}^{n} \left[T_i(T_i + 1) - Y_i^2 \right] v_i}{2 \sum_{i=1}^{n} Y_i^2 v_i} .$$

- Simplest choice: Add multiplets with $T(T + 1) = 3Y^2$ so that $\rho_{\rm tree} = 1$.
- SM: One $SU(2)_L$ doublet $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$ with $Y = \frac{1}{2}$.
- A simple extension: Two $SU(2)_L$ doublets $\Phi_i = \begin{pmatrix} \phi_i^+ \\ \phi_i^0 \end{pmatrix}$ (with $i = 1, 2$).

General 2HDM Potential

- Most general 2HDM potential in doublet field space $\Phi_{1,2}$:

$$
\begin{aligned}
V = & - \mu_1^2(\Phi_1^\dagger \Phi_1) - \mu_2^2(\Phi_2^\dagger \Phi_2) - \left[m_{12}^2(\Phi_1^\dagger \Phi_2) + \mathrm{H.c.} \right] \\
& + \lambda_1(\Phi_1^\dagger \Phi_1)^2 + \lambda_2(\Phi_2^\dagger \Phi_2)^2 + \lambda_3(\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4(\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) \\
& + \left[\frac{1}{2}\lambda_5(\Phi_1^\dagger \Phi_2)^2 + \lambda_6(\Phi_1^\dagger \Phi_1)(\Phi_1^\dagger \Phi_2) + \lambda_7(\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_2) + \mathrm{H.c.} \right] .
\end{aligned}
$$

- Four real mass parameters $\mu_{1,2}^2$, Re(m_{12}^2), Im(m_{12}^2).
- 10 real quartic couplings $\lambda_{1,2,3,4}$, Re($\lambda_{5,6,7}$), Im($\lambda_{5,6,7}$).
- Rich vacuum structure. [Battye, Brawn, Pilaftsis '11; Branco *et al* '12]

Higgs Spectrum in a General 2HDM

- Consider normal vacua with real vevs $v_{1,2}$, where $\sqrt{v_1^2 + v_2^2} = v_{\rm SM}$ and $\tan\beta = v_2/v_1$.
- Eight real scalar fields: $\phi_j = \begin{pmatrix} \phi_j^+ \\ \frac{1}{\sqrt{2}}(v_j + \rho_j + i\eta_j) \end{pmatrix}$ (with $j = 1, 2$).
- After EWSB, 3 Goldstone bosons ($G^\pm, G^0$), eaten by $W^\pm$ and Z.
- Five physical scalar fields: two CP-even (h, H), one CP-odd (a) and two charged ($h^\pm$).
- In the charged sector, $\begin{pmatrix} G^\pm \\ h^\pm \end{pmatrix} = \begin{pmatrix} \cos\beta & \sin\beta \\ -\sin\beta & \cos\beta \end{pmatrix} \begin{pmatrix} \phi_1^\pm \\ \phi_2^\pm \end{pmatrix}$.

$$M_{h^\pm}^2 = \frac{1}{s_\beta c_\beta}\left[\mathrm{Re}(m_{12}^2) - \frac{1}{2}\left(\{\lambda_4 + \mathrm{Re}(\lambda_5)\} s_\beta c_\beta + \mathrm{Re}(\lambda_6) c_\beta^2 + \mathrm{Re}(\lambda_7) s_\beta^2\right)\right].$$

- In the CP-odd sector, $\begin{pmatrix} G^0 \\ a \end{pmatrix} = \begin{pmatrix} \cos\beta & \sin\beta \\ -\sin\beta & \cos\beta \end{pmatrix} \begin{pmatrix} \eta_1 \\ \eta_2 \end{pmatrix}$.

$$M_a^2 = \frac{1}{s_\beta c_\beta}\left[\mathrm{Re}(m_{12}^2) - v^2\left(\mathrm{Re}(\lambda_5) s_\beta c_\beta + \frac{1}{2}\left\{\mathrm{Re}(\lambda_6) c_\beta^2 + \mathrm{Re}(\lambda_7) s_\beta^2\right\}\right)\right]$$
$$= M_{h^\pm}^2 + \frac{1}{2}\left[\lambda_4 - \mathrm{Re}(\lambda_5)\right] v^2.$$

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Higgs Spectrum in a General 2HDM

- In the *CP*-even sector, $\begin{pmatrix} H \\ h \end{pmatrix} = \begin{pmatrix} \cos\alpha & \sin\alpha \\ -\sin\alpha & \cos\alpha \end{pmatrix} \begin{pmatrix} \rho_1 \\ \rho_2 \end{pmatrix}$

$$
\begin{aligned}
M_S^2 &\equiv \begin{pmatrix} A & C \\ C & B \end{pmatrix} \\
&= M_a^2 \begin{pmatrix} s_\beta^2 & -s_\beta c_\beta \\ -s_\beta c_\beta & c_\beta^2 \end{pmatrix} \\
&\quad + v^2 \begin{pmatrix} 2\lambda_1 c_\beta^2 + \mathrm{Re}(\lambda_5) s_\beta^2 + 2\mathrm{Re}(\lambda_6) s_\beta c_\beta & \lambda_{34} s_\beta c_\beta + \mathrm{Re}(\lambda_6) c_\beta^2 + \mathrm{Re}(\lambda_7) s_\beta^2 \\ \lambda_{34} s_\beta c_\beta + \mathrm{Re}(\lambda_6) c_\beta^2 + \mathrm{Re}(\lambda_7) s_\beta^2 & 2\lambda_2 s_\beta^2 + \mathrm{Re}(\lambda_5) c_\beta^2 + 2\mathrm{Re}(\lambda_7) s_\beta c_\beta \end{pmatrix}
\end{aligned}
$$

with $\tan 2\alpha = 2C/(A-B)$.

- The SM Higgs boson is given by

$$H_{\mathrm{SM}} = \rho_1 \cos\beta + \rho_2 \sin\beta = H\cos(\beta-\alpha) + h\sin(\beta-\alpha)\ .$$

- SM alignment limit: $\alpha \to \beta$ (or $\beta - \pi/2$).
- Usually attributed to either decoupling or accidental cancellations.

  [Gunion, Haber '03; Ginzburg, Krawczyk '05; Carena, Low, Shah, Wagner '13]
- **Explore symmetries of the 2HDM potential to *naturally* justify the alignment limit.**

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Natural Alignment Condition

- Rewrite *CP*-even mass matrix as

$$M_S^2 = \begin{pmatrix} c_\beta & -s_\beta \\ s_\beta & c_\beta \end{pmatrix} \begin{pmatrix} \widehat{A}v^2 & \widehat{C}v^2 \\ \widehat{C}v^2 & M_a^2 + \widehat{B}v^2 \end{pmatrix} \begin{pmatrix} c_\beta & s_\beta \\ -s_\beta & c_\beta \end{pmatrix} \equiv O\widehat{M}_S^2 O^\mathsf{T} .$$

$$\widehat{A} = 2\Big[c_\beta^4\lambda_1 + s_\beta^2 c_\beta^2 \lambda_{345} + s_\beta^4 \lambda_2 \;+\; 2 s_\beta c_\beta \Big(c_\beta^2 \lambda_6 + s_\beta^2 \lambda_7\Big)\Big] ,$$

$$\widehat{B} = \lambda_5 \;+\; 2\Big[s_\beta^2 c_\beta^2 \Big(\lambda_1 + \lambda_2 - \lambda_{345}\Big) \;-\; s_\beta c_\beta \Big(c_\beta^2 - s_\beta^2\Big)\Big(\lambda_6 - \lambda_7\Big)\Big] ,$$

$$\widehat{C} = s_\beta^3 c_\beta \Big(2\lambda_2 - \lambda_{345}\Big) - c_\beta^3 s_\beta \Big(2\lambda_1 - \lambda_{345}\Big) + c_\beta^2 \Big(1 - 4 s_\beta^2\Big)\lambda_6 + s_\beta^2 \Big(4 c_\beta^2 - 1\Big)\lambda_7 .$$

- Exact alignment ($\alpha = \beta$) iff $\widehat{C} = 0$, i.e.

$$\lambda_7 t_\beta^4 - (2\lambda_2 - \lambda_{345}) t_\beta^3 + 3(\lambda_6 - \lambda_7) t_\beta^2 + (2\lambda_1 - \lambda_{345}) t_\beta - \lambda_6 \;=\; 0 .$$

- Natural alignment if happens for *any* value of $\tan\beta$, independent of non-SM Higgs spectra:

$$\boxed{\lambda_1 = \lambda_2 = \lambda_{345}/2 , \quad \lambda_6 = \lambda_7 = 0}$$

- CP-even Higgs masses are given by

$$M_H^2 = 2\lambda_2 v^2 \;\equiv\; \lambda_{\rm SM} v^2 ,$$
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Higgs Couplings in a General 2HDM

- With respect to the SM Higgs couplings $H_{\rm SM}VV$ ($V = W^{\pm}, Z$),

$$g_{hVV} = \sin(\beta - \alpha)\,, \qquad g_{HVV} = \cos(\beta - \alpha)\,.$$

$$g_{haZ} = \cos(\beta - \alpha)\,, \qquad g_{HaZ} = \sin(\beta - \alpha)\,,$$
$$g_{h^+hW^-} = \cos(\beta - \alpha)\,, \qquad g_{h^+HW^-} = \sin(\beta - \alpha)\,.$$

- Similar behavior for CP-even Higgs self-couplings:

$$g_{hHH} \propto \sin(\beta - \alpha)\,, \qquad g_{Hhh} \propto \cos(\beta - \alpha)\,.$$

- In the alignment limit ($\alpha \to \beta$), one of the neutral Higgses (h) is gaugephobic.
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Z_2-symmetric 2HDM

- General Yukawa Lagrangian

$$-\mathcal{L}_Y = \bar{Q}_L(h_1^u\Phi_1 + h_2^u\Phi_2)u_R + \bar{Q}_L(h_1^d\widetilde{\Phi}_1 + h_2^d\widetilde{\Phi}_2)d_R + \bar{L}_L(h_1^e\widetilde{\Phi}_1 + h_2^e\widetilde{\Phi}_2)e_R \, .$$

- Dangerous FCNC processes at tree-level.
- Can be naturally avoided by imposing a Z_2-symmetry. [Glashow, Weinberg '58]

	Z_2 charge							Coupling		
	Φ_1	Φ_2	Q_L	L_L	u_R	d_R	e_R	u_R	d_R	e_R
Type-I	$+$	$-$	$+$	$+$	$-$	$-$	$-$	Φ_2	Φ_2	Φ_2
Type-II (MSSM-type)	$+$	$-$	$+$	$+$	$-$	$+$	$+$	Φ_2	Φ_1	Φ_1
Type-X (Lepton-specific)	$+$	$-$	$+$	$+$	$-$	$-$	$+$	Φ_2	Φ_2	Φ_1
Type-Y (Flipped)	$+$	$-$	$+$	$+$	$-$	$+$	$-$	Φ_2	Φ_1	Φ_2

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Type-II (MSSM-type)	+	−	+	+	−	+	+	Φ_2	Φ_1	Φ_1
Type-X (Lepton-specific)	+	−	+	+	−	−	+	Φ_2	Φ_2	Φ_1
Type-Y (Flipped)	+	−	+	+	−	+	−	Φ_2	Φ_1	Φ_2

Symmetry Classifications of the 2HDM Potential

- Three classes of accidental symmetries of the 2HDM potential:
  - Higgs Family (HF) Symmetries involving transformations of $\Phi_{1,2}$ only (but not $\Phi^*_{1,2}$), e.g. Z_2 [Glashow, Weinberg '58], $U(1)_{\rm PQ}$ [Peccei, Quinn '77], $SO(3)_{\rm HF}$ [Deshpande, Ma '78; Ivanov '07; Ma, Maniatis '09; Ferreira, Haber, Maniatis, Nachtmann, Silva '10].
  - CP Symmetries relating $\Phi_{1,2}$ to $\Phi^*_{1,2}$, e.g. $\Phi_{1(2)} \to \Phi^*_{1(2)}$ (CP1) [Lee '73; Branco '80], $\Phi_{1(2)} \to (-)\Phi^*_{2(1)}$ (CP2) [Davidson, Haber '05], CP1 combined with $SO(2)_{\rm HF}/Z_2$ (CP3) [Ivanov '07; Ferreira, Haber, Silva '09; Ma, Maniatis '09; Ferreira, Haber, Maniatis, Nachtmann, Silva '10].
  - Additional mixed HF and CP symmetries that leave the gauge-kinetic terms of $\Phi_{1,2}$ invariant [Battye, Brawn, Pilaftsis '11].
- Includes *all* custodial symmetries of the 2HDM potential.
- Maximum of 13 *distinct* accidental symmetries of the general 2HDM potential.
- Each of them imposes specific relations among the scalar parameters.

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Bilinear Formalism

- Introduce an 8-dimensional complex multiplet: [Battye, Brawn, Pilaftsis '11; Nishi '11; Pilaftsis '12]

$$\mathbf{\Phi} = \begin{pmatrix} \Phi_1 \\ \Phi_2 \\ i\sigma^2\Phi_1^* \\ i\sigma^2\Phi_2^* \end{pmatrix} .$$

- $\mathbf{\Phi}$ satisfies the Majorana condition: $\mathbf{\Phi} = C\mathbf{\Phi}^*$, where $C = \sigma^2 \otimes \sigma^0 \otimes \sigma^2 = C^{-1} = C^*$.

- Define a null 6-dimensional Lorentz vector bilinear in Φ:

$$R^A = \Phi^\dagger \Sigma^A \Phi ,$$

(with $A = 0, 1, 2, 3, 4, 5$), where

$$\Sigma^0 = \frac{1}{2}\sigma^0 \otimes \sigma^0 \otimes \sigma^0 \equiv \frac{1}{2}\mathbf{1}_8, \quad \Sigma^1 = \frac{1}{2}\sigma^0 \otimes \sigma^1 \otimes \sigma^0, \quad \Sigma^2 = \frac{1}{2}\sigma^3 \otimes \sigma^2 \otimes \sigma^0,$$

$$\Sigma^3 = \frac{1}{2}\sigma^0 \otimes \sigma^3 \otimes \sigma^0, \quad \Sigma^4 = -\frac{1}{2}\sigma^2 \otimes \sigma^2 \otimes \sigma^0, \quad \Sigma^5 = -\frac{1}{2}\sigma^1 \otimes \sigma^2 \otimes \sigma^0.$$

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(with $A = 0, 1, 2, 3, 4, 5$), where

$$\Sigma^0 = \frac{1}{2}\sigma^0 \otimes \sigma^0 \otimes \sigma^0 \equiv \frac{1}{2}\mathbf{1}_8, \quad \Sigma^1 = \frac{1}{2}\sigma^0 \otimes \sigma^1 \otimes \sigma^0, \quad \Sigma^2 = \frac{1}{2}\sigma^3 \otimes \sigma^2 \otimes \sigma^0,$$
$$\Sigma^3 = \frac{1}{2}\sigma^0 \otimes \sigma^3 \otimes \sigma^0, \quad \Sigma^4 = -\frac{1}{2}\sigma^2 \otimes \sigma^2 \otimes \sigma^0, \quad \Sigma^5 = -\frac{1}{2}\sigma^1 \otimes \sigma^2 \otimes \sigma^0.$$

2HDM Potential in Bilinear Field Space

- The general 2HDM potential takes a simple form:

$$V = -\frac{1}{2}M_A R^A + \frac{1}{4}L_{AB}R^A R^B .$$

$$M = \Big(\mu_1^2+\mu_2^2,\quad 2\mathrm{Re}(m_{12}^2),\quad -2\mathrm{Im}(m_{12}^2),\quad \mu_1^2-\mu_2^2,\quad 0,\quad 0\Big),$$

$$R = \begin{pmatrix} \Phi_1^\dagger\Phi_1+\Phi_2^\dagger\Phi_2 \\ \Phi_1^\dagger\Phi_2+\Phi_2^\dagger\Phi_1 \\ -i(\Phi_1^\dagger\Phi_2-\Phi_2^\dagger\Phi_1) \\ \Phi_1^\dagger\Phi_1-\Phi_2^\dagger\Phi_2 \\ \Phi_1^{\mathsf T}i\sigma^2\Phi_2-\Phi_2^\dagger i\sigma^2\Phi_1^* \\ -i(\Phi_1^{\mathsf T}i\sigma^2\Phi_2+\Phi_2^\dagger i\sigma^2\Phi_1^*) \end{pmatrix},$$

$$L = \begin{pmatrix} \lambda_1+\lambda_2+\lambda_3 & \mathrm{Re}(\lambda_6+\lambda_7) & -\mathrm{Im}(\lambda_6+\lambda_7) & \lambda_1-\lambda_2 & 0 & 0 \\ \mathrm{Re}(\lambda_6+\lambda_7) & \lambda_4+\mathrm{Re}(\lambda_5) & -\mathrm{Im}(\lambda_5) & \mathrm{Re}(\lambda_6-\lambda_7) & 0 & 0 \\ -\mathrm{Im}(\lambda_6+\lambda_7) & -\mathrm{Im}(\lambda_5) & \lambda_4-\mathrm{Re}(\lambda_5) & -\mathrm{Im}(\lambda_6-\lambda_7) & 0 & 0 \\ \lambda_1-\lambda_2 & \mathrm{Re}(\lambda_6-\lambda_7) & -\mathrm{Im}(\lambda_6-\lambda_7) & \lambda_1+\lambda_2-\lambda_3 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

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$$L = \begin{pmatrix} \lambda_1 + \lambda_2 + \lambda_3 & \mathrm{Re}(\lambda_6 + \lambda_7) & -\mathrm{Im}(\lambda_6 + \lambda_7) & \lambda_1 - \lambda_2 & 0 & 0 \\ \mathrm{Re}(\lambda_6 + \lambda_7) & \lambda_4 + \mathrm{Re}(\lambda_5) & -\mathrm{Im}(\lambda_5) & \mathrm{Re}(\lambda_6 - \lambda_7) & 0 & 0 \\ -\mathrm{Im}(\lambda_6 + \lambda_7) & -\mathrm{Im}(\lambda_5) & \lambda_4 - \mathrm{Re}(\lambda_5) & -\mathrm{Im}(\lambda_6 - \lambda_7) & 0 & 0 \\ \lambda_1 - \lambda_2 & \mathrm{Re}(\lambda_6 - \lambda_7) & -\mathrm{Im}(\lambda_6 - \lambda_7) & \lambda_1 + \lambda_2 - \lambda_3 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

13 Symmetries of the 2HDM Potential

[Pilaftsis '12]

Table 1
Parameter relations for the 13 accidental symmetries [1] related to the $U(1)_Y$-invariant 2HDM potential in the diagonally reduced basis, where $\mathrm{Im}\,\lambda_5 = 0$ and $\lambda_6 = \lambda_7$. A dash signifies the absence of a constraint.

No.	Symmetry	μ_1^2	μ_2^2	m_{12}^2	λ_1	λ_2	λ_3	λ_4	$\mathrm{Re}\,\lambda_5$	$\lambda_6 = \lambda_7$
1	$Z_2 \times O(2)$	–	–	Real	–	–	–	–	–	Real
2	$(Z_2)^2 \times SO(2)$	–	–	0	–	–	–	–	–	0
3	$(Z_2)^3 \times O(2)$	–	μ_1^2	0	–	λ_1	–	–	–	0
4	$O(2) \times O(2)$	–	–	0	–	–	–	–	0	0
5	$Z_2 \times [O(2)]^2$	–	μ_1^2	0	–	λ_1	–	–	$2\lambda_1 - \lambda_{34}$	0
6	$O(3) \times O(2)$	–	μ_1^2	0	–	λ_1	–	$2\lambda_1 - \lambda_3$	0	0
7	$SO(3)$	–	–	Real	–	–	–	–	λ_4	Real
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10	$O(2) \times O(3)$	–	μ_1^2	0	–	λ_1	$2\lambda_1$	–	0	0
11	$SO(4)$	–	–	0	–	–	–	0	0	0
12	$Z_2 \times O(4)$	–	μ_1^2	0	–	λ_1	–	0	0	0
13	$SO(5)$	–	μ_1^2	0	–	λ_1	$2\lambda_1$	0	0	0

- Recall natural alignment condition $\lambda_1 = \lambda_2 = \lambda_{345}/2$, $\lambda_6 = \lambda_7 = 0$.
- *Only* three symmetries satisfy this.

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Maximal Symmetry Group

- *Maximal* symmetry group in the bilinear field space: $G^R_{\rm 2HDM} = {\rm SO}(5)$.
- In the original Φ-field space, $G^\Phi_{\rm 2HDM} = ({\rm Sp}(4)/Z_2) \otimes {\rm SU}(2)_L$.
- Conjecture: In a general nHDM, $G^\Phi_{n{\rm HDM}} = ({\rm Sp}(2n)/Z_2) \otimes {\rm SU}(2)_L$.
- For the SM, reproduces the well-known custodial symmetry $G^\Phi_{\rm SM} = ({\rm SU}(2)_C/Z_2) \otimes {\rm SU}(2)_L$. [Sikivie, Susskind, Voloshin, Zakharov '80].

- In 2HDM, 3 different realizations of custodial symmetry with [Pilaftsis '12; BD, Pilaftsis '14]

$$
\begin{array}{llll}
\text{(i)} & h_1^u = e^{i\theta} h_1^d & \text{and} & h_2^u = e^{i\theta} h_2^d\,, \\
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Maximally Symmetric 2HDM

- In the SO(5) limit, the 2HDM potential is very simple:

$$V = -\mu^2 \left(|\Phi_1|^2 + |\Phi_2|^2\right) + \lambda \left(|\Phi_1|^2 + |\Phi_2|^2\right)^2 = -\frac{\mu^2}{2}\, \mathbf{\Phi}^\dagger \mathbf{\Phi} + \frac{\lambda}{4} (\mathbf{\Phi}^\dagger \mathbf{\Phi})^2 .$$

- More minimal than the MSSM scalar potential, which in the custodial limit $g' \to 0$, has a smaller symmetry: $\mathrm{O}(2) \otimes \mathrm{O}(3) \subset \mathrm{SO}(5)$.
- After EWSB in the MS-2HDM, one massive Higgs boson H with $M_H^2 = 2\lambda_2 v^2$, whilst remaining four (h, a and $h^\pm$) are massless [Goldstone theorem].
- **Natural SM alignment limit with $\alpha = \beta$.** [Recall $H_{\rm SM} = H\cos(\beta-\alpha) + h\sin(\beta-\alpha)$]
- (Pseudo)-Goldstones in MS-2HDM acquire mass due to custodial symmetry-breaking g' and Yukawa coupling effects.

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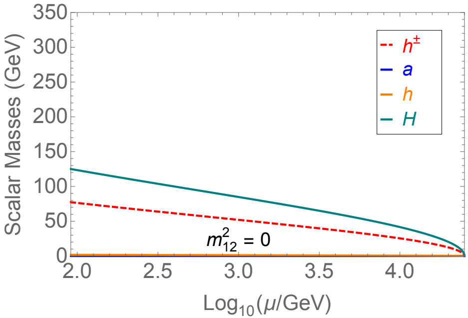
g' and Yukawa Coupling Effects

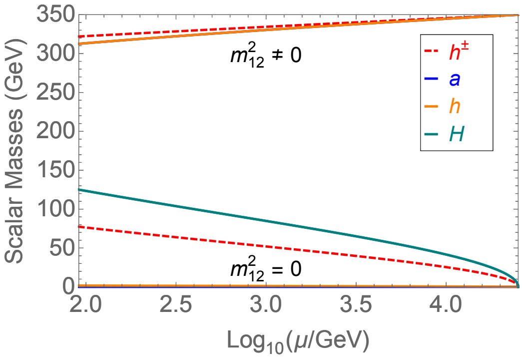
- Custodial symmetry broken by non-zero g' and Yukawa couplings.

$$\begin{aligned} \mathrm{SO}(5)\otimes\mathrm{SU}(2)_L \quad &\xrightarrow{g'\neq 0} \quad \mathrm{O}(3)\otimes\mathrm{O}(2)\otimes\mathrm{SU}(2)_L \;\sim\; \mathrm{O}(3)\otimes\mathrm{U}(1)_Y\otimes\mathrm{SU}(2)_L \\ &\xrightarrow{\text{Yukawa}} \quad \mathrm{O}(2)\otimes\mathrm{U}(1)_Y\otimes\mathrm{SU}(2)_L \;\sim\; \mathrm{U}(1)_{\mathrm{PQ}}\otimes\mathrm{U}(1)_Y\otimes\mathrm{SU}(2)_L \\ &\xrightarrow{\langle\Phi_{1,2}\rangle\neq 0} \quad \mathrm{U}(1)_{\mathrm{em}}\ . \end{aligned}$$

- Assume SO(5)-symmetry scale $\mu_X \gg v$.
- Use two-loop RGEs to find the mass spectrum at weak scale.



Soft Breaking Effects

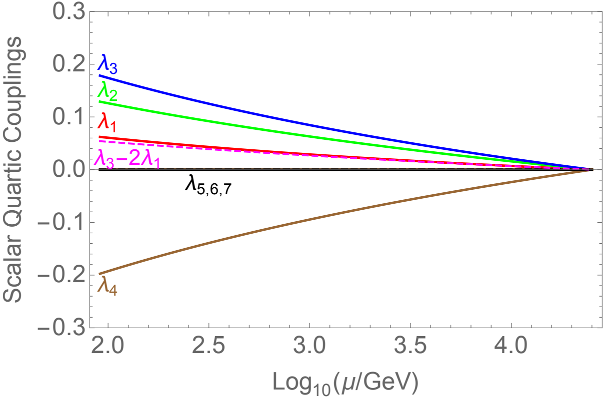


- In the $SO(5)$ limit for quartic couplings,

$$M_H^2 = 2\lambda_2 v^2 , \qquad M_h^2 = M_a^2 = M_{h^\pm}^2 = \frac{\mathrm{Re}(m_{12}^2)}{s_\beta c_\beta} .$$

- Still preserves natural alignment, irrespective of other 2HDM parameters.
- Predicts a quasi-degenerate heavy Higgs sector.

Quartic Coupling Unification



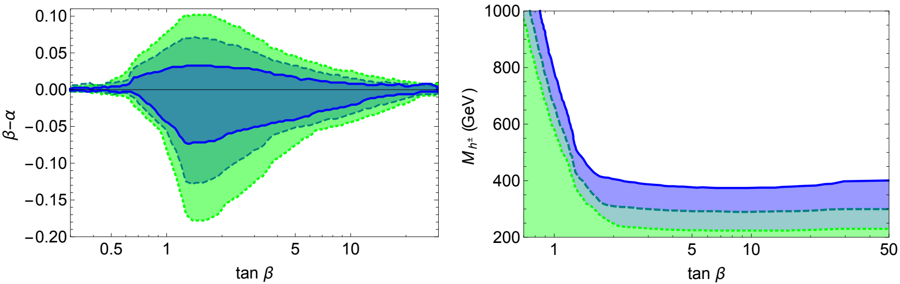
Constraints from Global Fit

- Electroweak precision observables.
- LHC signal strengths of the light *CP*-even Higgs boson.
- Limits on heavy *CP*-even scalar from $H \to WW, ZZ, \tau\tau$ searches.
- Flavor observables such as B_s mixing and $B \to X_s\gamma$.

- Stability of the potential:

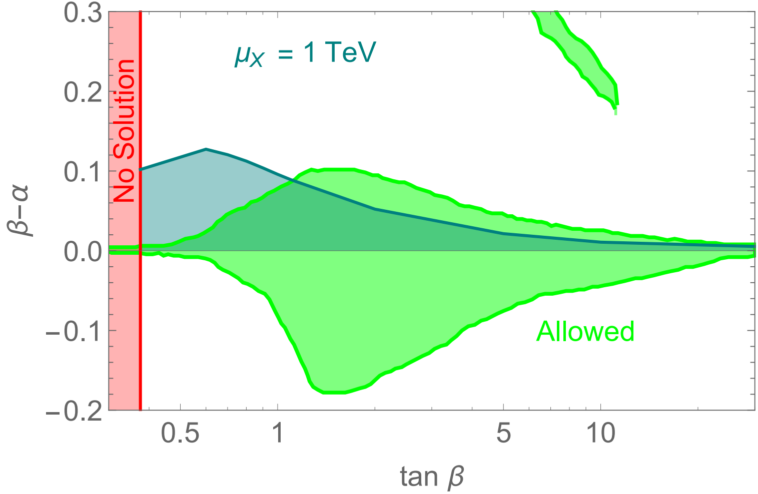
$$\lambda_{1,2} > 0, \quad \lambda_3 + \sqrt{\lambda_1\lambda_2} > 0, \quad \lambda_3 + \lambda_4 + \sqrt{\lambda_1\lambda_2} - \mathrm{Re}(\lambda_5) > 0.$$

- Perturbativity of the Higgs self-couplings: $\|S_{\Phi\Phi\to\Phi\Phi}\| < \frac{1}{8}$.



[reinterpreted from Baglio, Eberhardt, Nierste, Wiebusch ’13]

Misalignment Predictions in MS-2HDM



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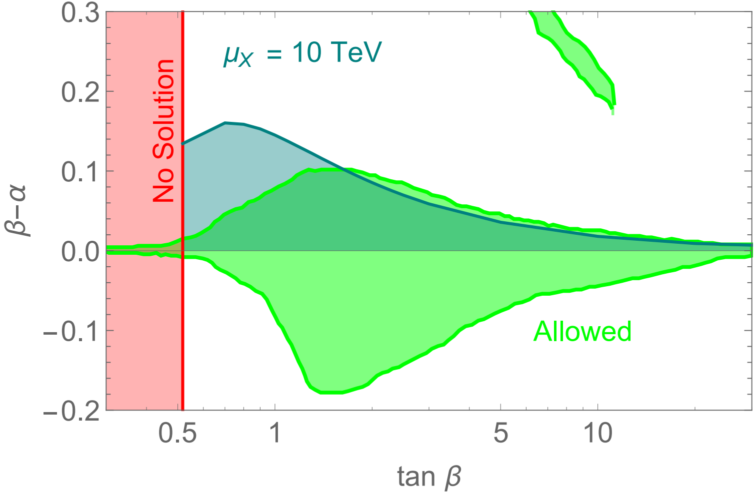

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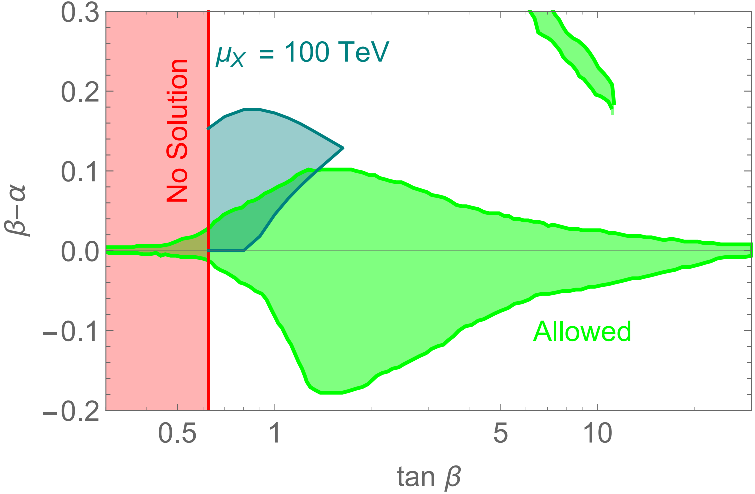

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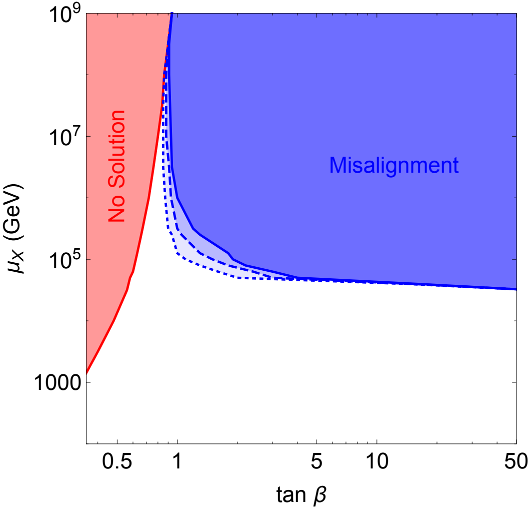
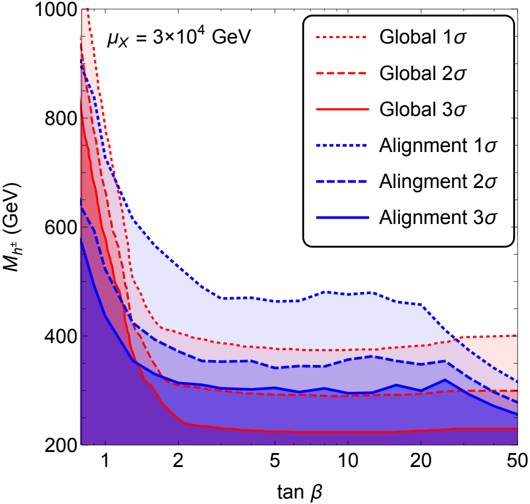
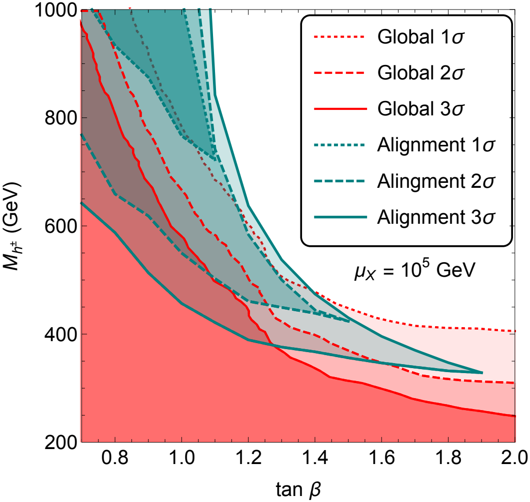


Lower Limit on Charged Higgs Mass

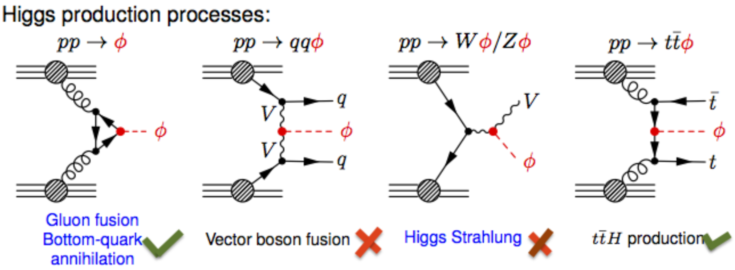


Lower *and* Upper Limits on Charged Higgs Mass

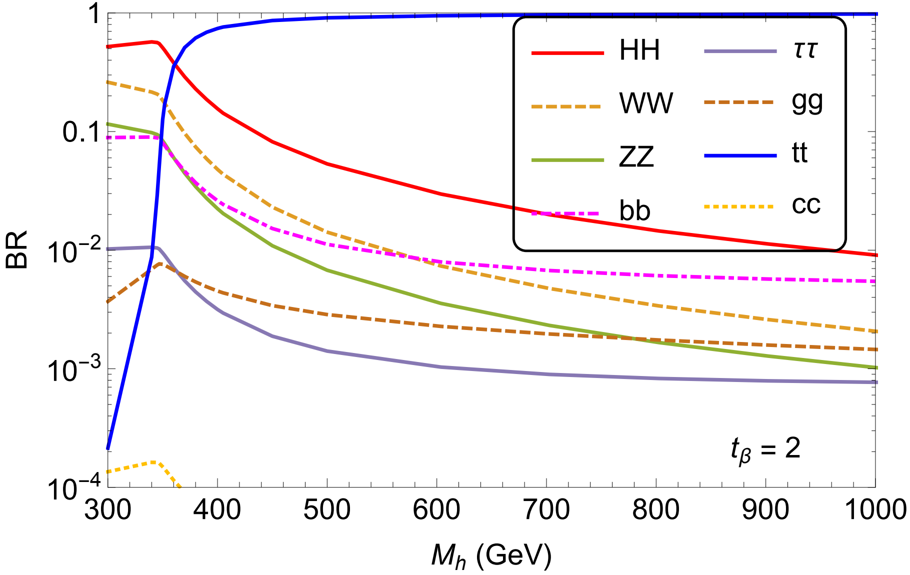


Implications for the LHC Searches

- Recall that $g_{hVV} = \sin(\beta - \alpha)$, $g_{HVV} = \cos(\beta - \alpha)$.
- In the alignment limit $\alpha \to \beta$, H is SM-like and the heavy Higgs h is gaugephobic.
- Dominant production modes at the LHC: ggF and associated production with $t\bar{t}$.



Branching Fractions



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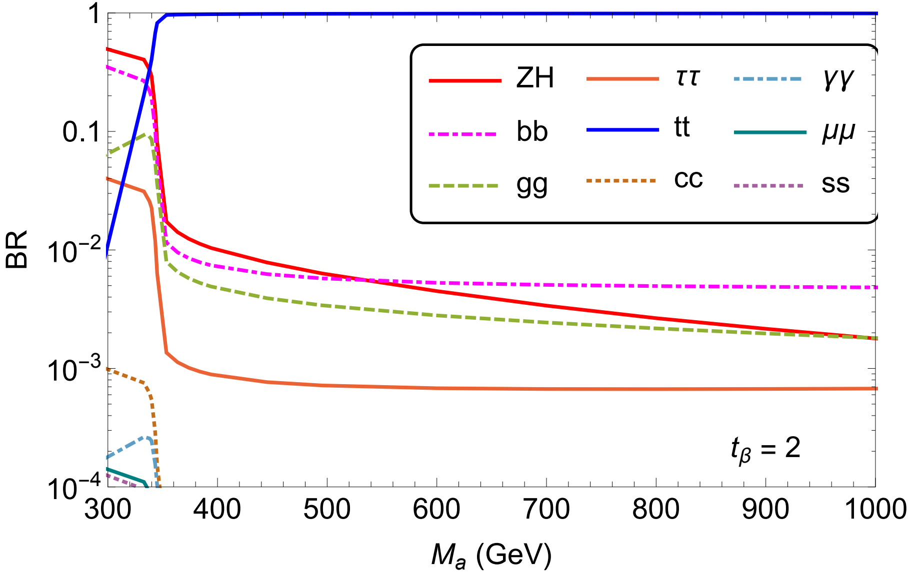

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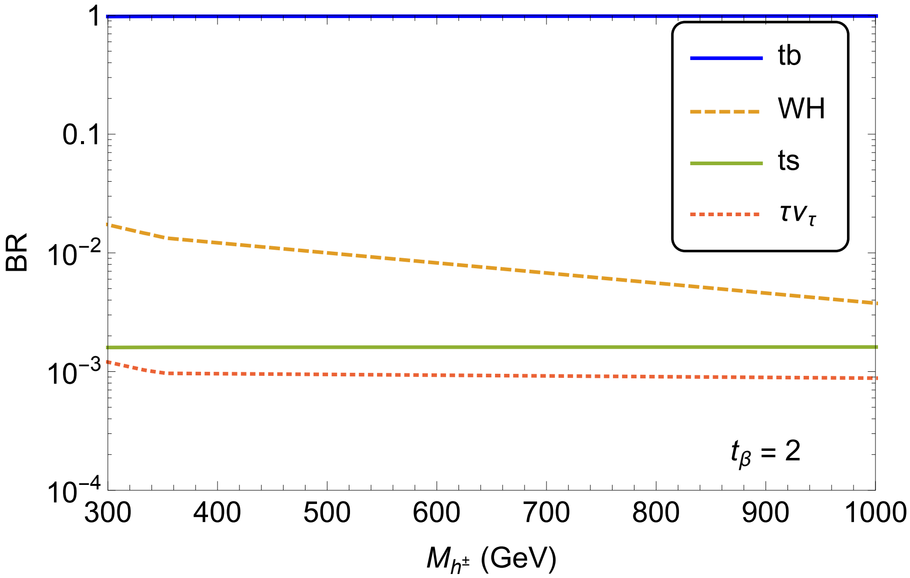
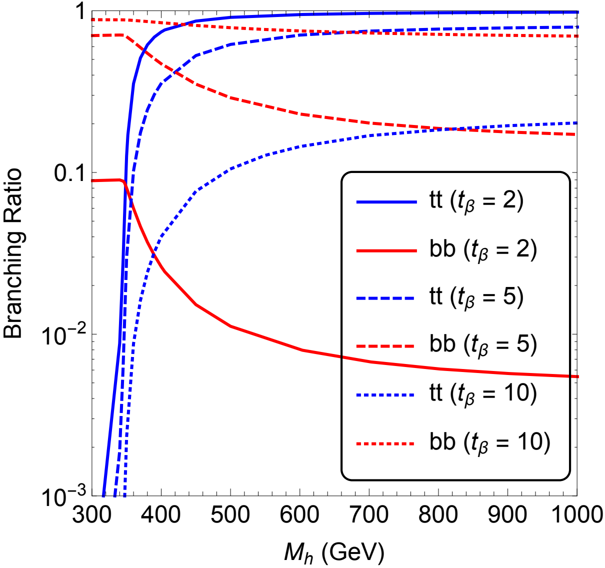


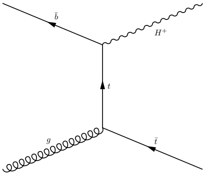
tan β Dependance

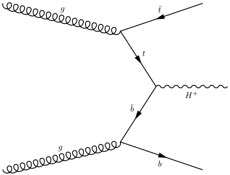


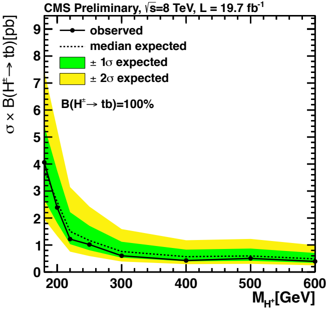
LHC Searches so far

- Existing collider limits on the heavy Higgs sector derived from *WW* and *ZZ* modes are not applicable in the alignment limit.
- Limits from $gg \to h \to \tau^+\tau^-$ and $gg \to b\bar{b}h \to b\bar{b}\tau^+\tau^-$ are easily satisfied.
- Similarly for $h \to HH \to \gamma\gamma bb$.
- In the charged-Higgs sector, most of the searches focus on the low-mass regime ($M_{h^\pm} < M_t$): $pp \to tt \to Wbbh^+,\ h^+ \to cs$.
- Recently, the search was extended beyond the top-threshold: [CMS-PAS-HIG-13-026]

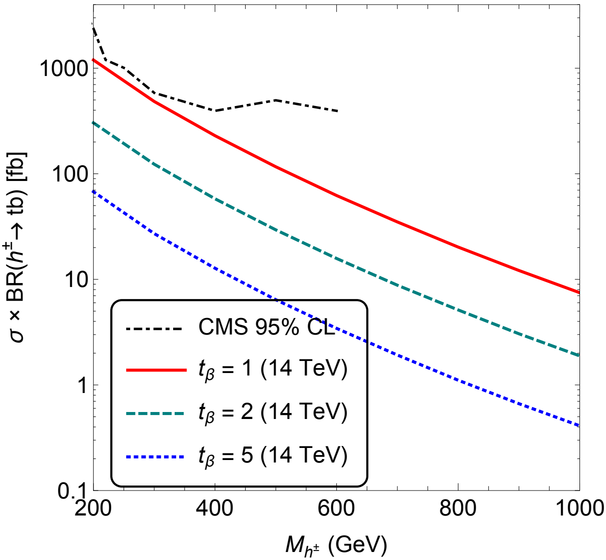
$$gg \to h^+ tb \to (\ell\nu bb)(\ell'\nu b)b$$







Predictions in the MS-2HDM



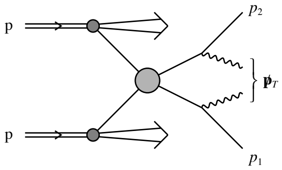
Simulations for $\sqrt{s} = 14$ TeV LHC

- Used `MadGraph5_aMC@NLO`.
- Event reconstruction using the CMS cuts:

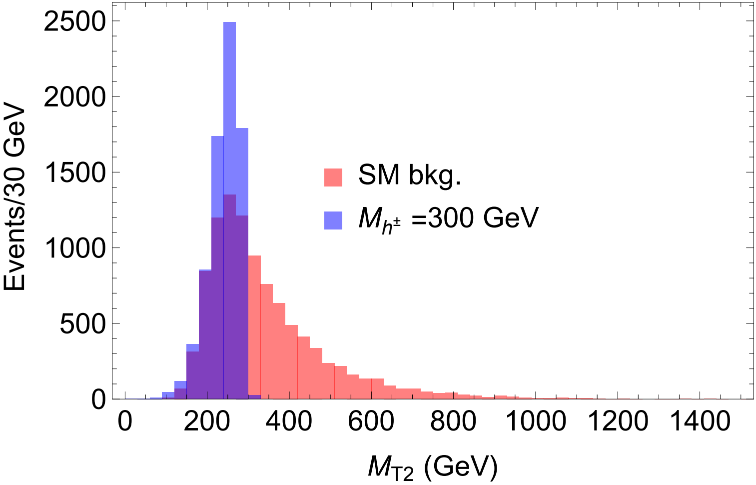
$$\begin{aligned} p_T^\ell &> 20\ \mathrm{GeV}, \quad |\eta^\ell| < 2.5, \quad \Delta R^{\ell\ell} > 0.4, \\ M_{\ell\ell} &> 12\ \mathrm{GeV}, \quad |M_{\ell\ell} - M_Z| > 10\ \mathrm{GeV}, \\ p_T^j &> 30\ \mathrm{GeV}, \quad |\eta^j| < 2.4, \quad \not{E}_T > 40\ \mathrm{GeV}. \end{aligned}$$

- For charged Higgs mass reconstruction, used 'stransverse mass' variable [Lester, Summers '99]

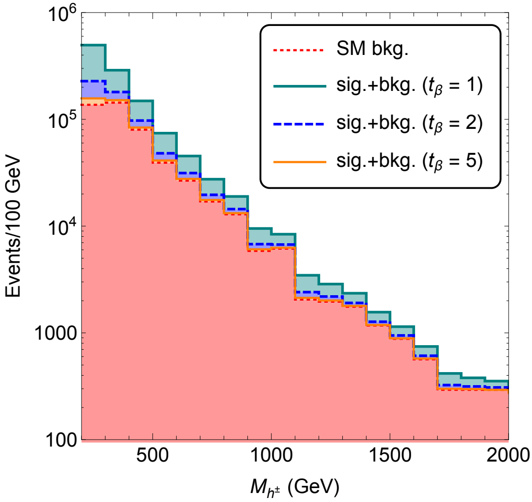
$$M_{T2} = \min_{\left\{\not{p}_{T_1} + \not{p}_{T_2} = \not{p}_T\right\}} \left[\max\left\{m_{T_1}, m_{T_2}\right\}\right] .$$



Mass Reconstruction using M_{T2}



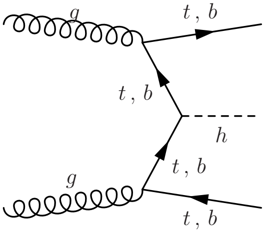
Reach at 14 TeV LHC

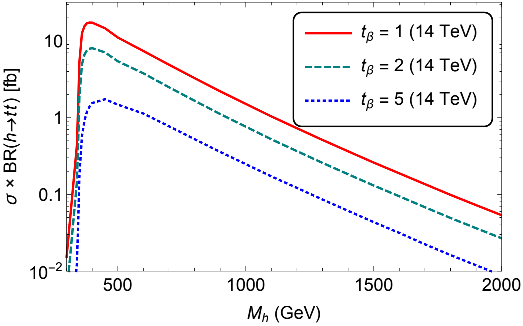


New Signal in the Neutral Higgs Sector

$$gg \rightarrow t\bar{t}h \rightarrow t\bar{t}t\bar{t}$$

- Existing 95% CL experimental upper limit on $\sigma_{t\bar{t}t\bar{t}}$ is 32 fb (CMS).
- SM prediction for $\sigma(pp \rightarrow t\bar{t}t\bar{t} + X) \simeq$ 10–15 fb at NLO. [Bevilacqua, Worek '12]
- Still lot of room for BSM contribution.





Towards a Full Analysis of the $t\bar{t}t\bar{t}$ Signal

35 final states, grouped into five channels:

- Fully hadronic: 12 jets, with 4 b-jets.
- Mostly hadronic: 6 light jets, 4 b-jets, one charged lepton and $\not{E}_T$.
- Semi-leptonic/hadronic: 4 light jets, 4 b-jets, 2 charged leptons and $\not{E}_T$.
- Mostly leptonic: 2 light jets, 4 b-jets, 3 charged leptons and $\not{E}_T$.
- Fully leptonic: 4 b-jets, 4 charged leptons and $\not{E}_T$.

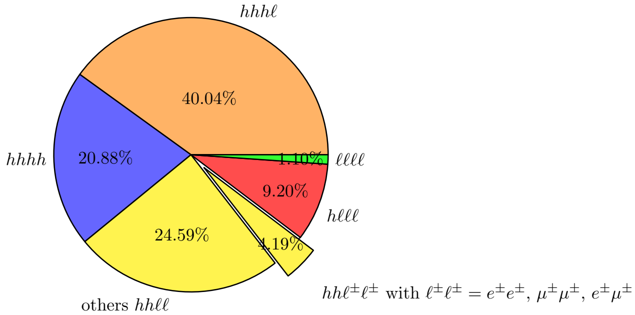


Figure 1.4: Branching fractions for the different decays of the four top quarks, depending on whether the W boson decays hadronically (h) or leptonically (ℓ).

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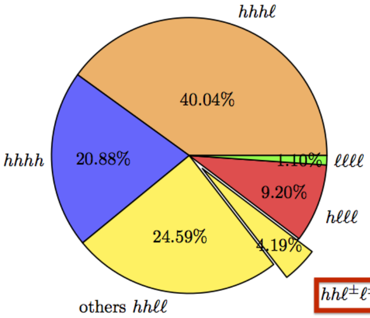
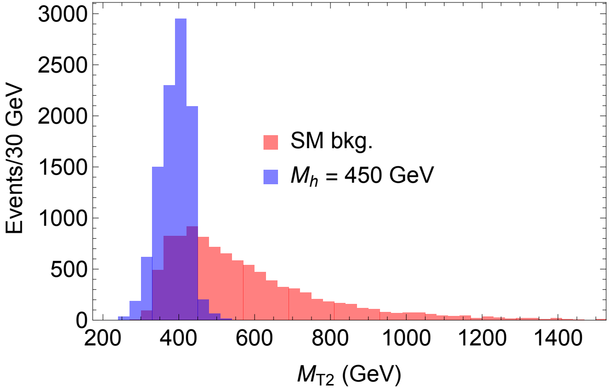


$hh\ell^{\pm}\ell^{\pm}$ with $\ell^{\pm}\ell^{\pm} = e^{\pm}e^{\pm},\ \mu^{\pm}\mu^{\pm},\ e^{\pm}\mu^{\pm}$

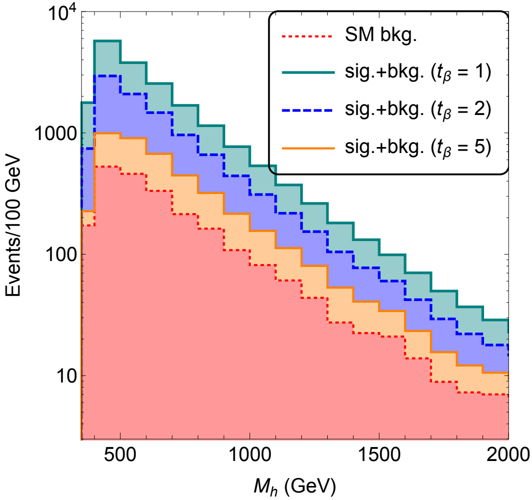
experimentally favored channel

Figure 1.4: Branching fractions for the different decays of the four top quarks, depending on whether the W boson decays hadronically (h) or leptonically (ℓ).

Mass Reconstruction using M_{T2}



$pp \to t\bar{t}h \to t\bar{t}t\bar{t}$ Signal



Conclusions

- Examined the SM alignment limit of the 2HDM potential.
- Listed the symmetries leading to natural alignment.
- Analyzed the simplest one, namely, the Maximally Symmetric 2HDM potential with SO(5) symmetry.
- Deviations from alignment are induced naturally by RG effects due to g' and Yukawa couplings, and due to soft-breaking mass parameter.
- Predicts a quasi-degenerate and ‘gaugephobic’ heavy Higgs sector.
- Using the alignment constraints, we predict lower limits on the heavy Higgs spectrum, which prevail the present limits in a wide range of parameter space.
- Depending on the SO(5)-breaking scale, we also obtain an upper limit on the heavy Higgs masses, which could be completely probed during LHC run-II.
- Initiated study on a new collider signal with four top quarks in the final state, which can become a valuable observational tool to directly probe the heavy Higgs sector in the alignment limit.

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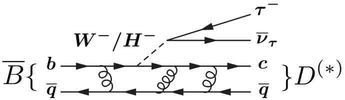
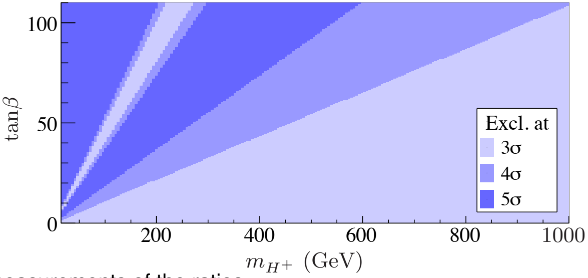
ps: BABAR Result





- Improved measurements of the ratios

$$\mathcal{R}(D) = \frac{\mathrm{BR}(\overline{B} \to D\tau^-\overline{\nu}_\tau)}{\mathrm{BR}(\overline{B} \to D\ell^-\overline{\nu}_\ell)}, \qquad \mathcal{R}(D^*) = \frac{\mathrm{BR}(\overline{B} \to D^*\tau^-\overline{\nu}_\tau)}{\mathrm{BR}(\overline{B} \to D^*\ell^-\overline{\nu}_\ell)},$$

- Excess w.r.t. the SM predictions at 3.4σ. Charged Higgs interpretation?
- Disfavored in Type-II 2HDM, but by the same token, SM is also excluded !!

Symmetry Generators

Table 2

Symmetry generators [cf. (10), (14)] and discrete group elements [cf. (17)] for the 13 accidental symmetries of the $U(1)_Y$-invariant 2HDM potential. For each symmetry, the maximally broken SO(5) generators compatible with a neutral vacuum are displayed, along with the pseudo-Goldstone bosons (given in parentheses) that result from the Goldstone theorem.

No.	Symmetry	Generators $T^a \leftrightarrow K^a$	Discrete group elements	Maximally broken SO(5) generators	Number of pseudo-Goldstone bosons
1	$Z_2 \times O(2)$	T^0	D_{CP1}	–	0
2	$(Z_2)^2 \times SO(2)$	T^0	D_{Z_2}	–	0
3	$(Z_2)^3 \times O(2)$	T^0	D_{CP2}	–	0
4	$O(2) \times O(2)$	T^3, T^0	–	T^3	1 (a)
5	$Z_2 \times [O(2)]^2$	T^2, T^0	D_{CP1}	T^2	1 (h)
6	$O(3) \times O(2)$	$T^{1,2,3}, T^0$	–	$T^{1,2}$	2 (h, a)
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13	$SO(5)$	$T^{0,1,2,\dots,9}$	–	$T^{1,2,8,9}$	4 $(h, a, h^\pm)$

[Pilaftsis '12]

- T^a and K^a are the generators of $SO(5)$ and $Sp(4)$ respectively ($a = 0, ..., 9$).
- T^0 is the hypercharge generator in R-space, which is equivalent to the electromagnetic generator $Q_{\rm em} = \frac{1}{2}\sigma^0 \otimes \sigma^0 \otimes \sigma^3 + K^0$ in Φ-space.
- $Sp(4)$ contains the custodial symmetry group $SU(2)_C$.
- Three *independent* realizations of custodial symmetry induced by (i) $K^{0,4,6}$, (ii) $K^{0,5,7}$, (iii) $K^{0,8,9}$.

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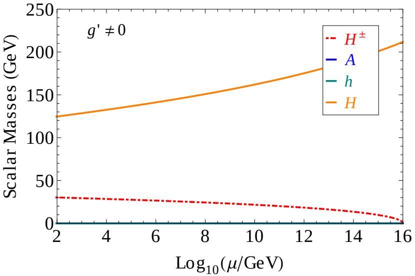
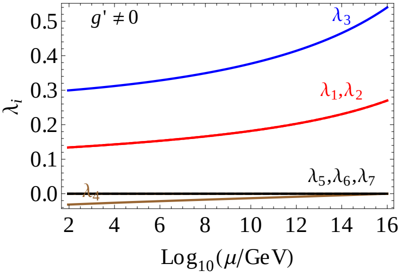
Quark Yukawa Couplings

- By convention, choose $h_1^u = 0$. For Type-I (Type-II) 2HDM, $h_1^d (h_2^d) = 0$.
- Quark yukawa couplings w.r.t. the SM are given by

Coupling	Type-I	Type-II
$g_{ht\bar{t}}$	$\cos\alpha / \sin\beta$	$\cos\alpha / \sin\beta$
$g_{hb\bar{b}}$	$\cos\alpha / \sin\beta$	$-\sin\alpha / \cos\beta$
$g_{Ht\bar{t}}$	$\sin\alpha / \sin\beta$	$\sin\alpha / \sin\beta$
$g_{Hb\bar{b}}$	$\sin\alpha / \sin\beta$	$\cos\alpha / \cos\beta$
$g_{at\bar{t}}$	$\cot\beta$	$\cot\beta$
$g_{ab\bar{b}}$	$-\cot\beta$	$\tan\beta$

g' Effect





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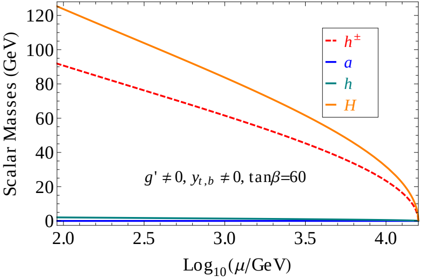
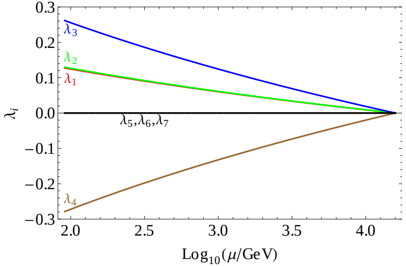
Yukawa Coupling Effects





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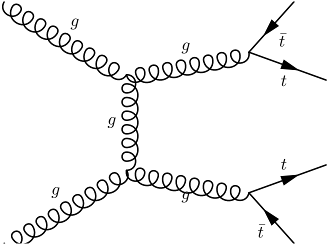
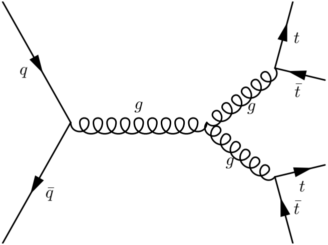
Production of 4 tops in the SM



(a)



(b)

Production of 4 tops in BSM